

Dwarf Novae in the Shortest Orbital Period Regime

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DN in the shortest P_{orb} regime

- Two interesting points of view
 - Binary evolution
 - Final stage of the low-mass binary evolution
 - Dynamics of accretion disks
 - Extreme mass-ratio
 - Strong tidal effect

In the view point of

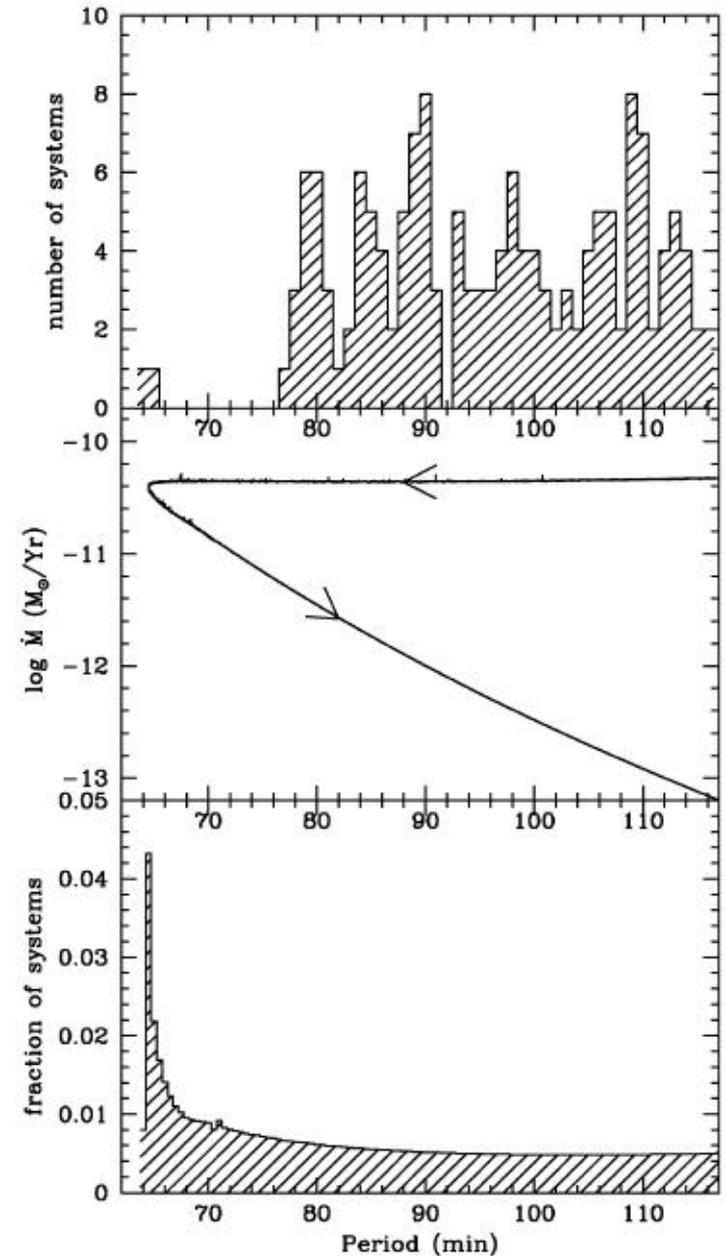
Binary evolution

- WZ Sge-type DNe as a “missing” population
 - Uemura+10 (PASJ, 62, 613)

Problems in the P_{orb} distribution

- Period spike problem
- Period minimum problem

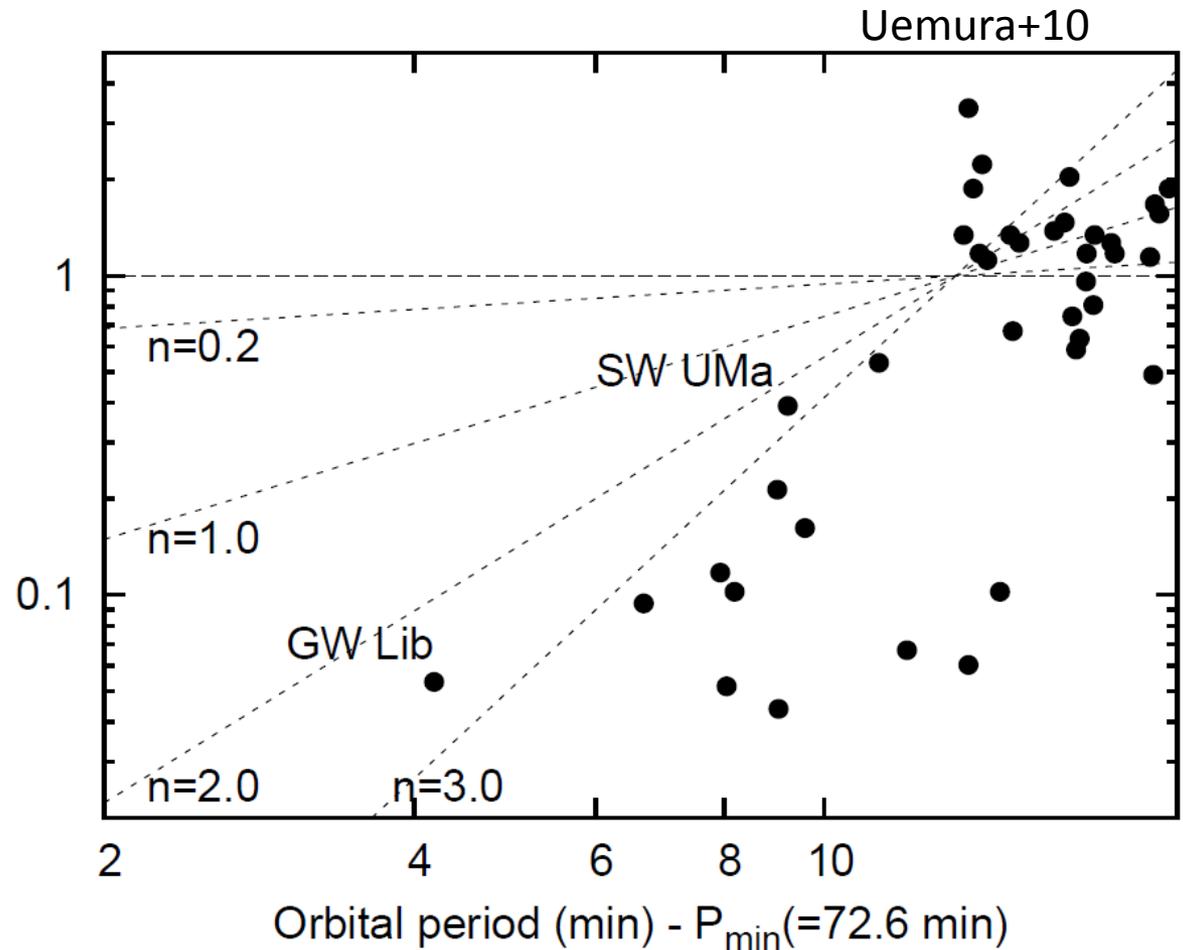
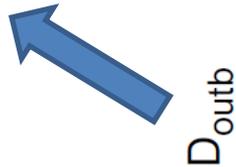
Porb distribution
(the figures from Barker & Kolb 2003)



P_{orb} dependence on recurrence times

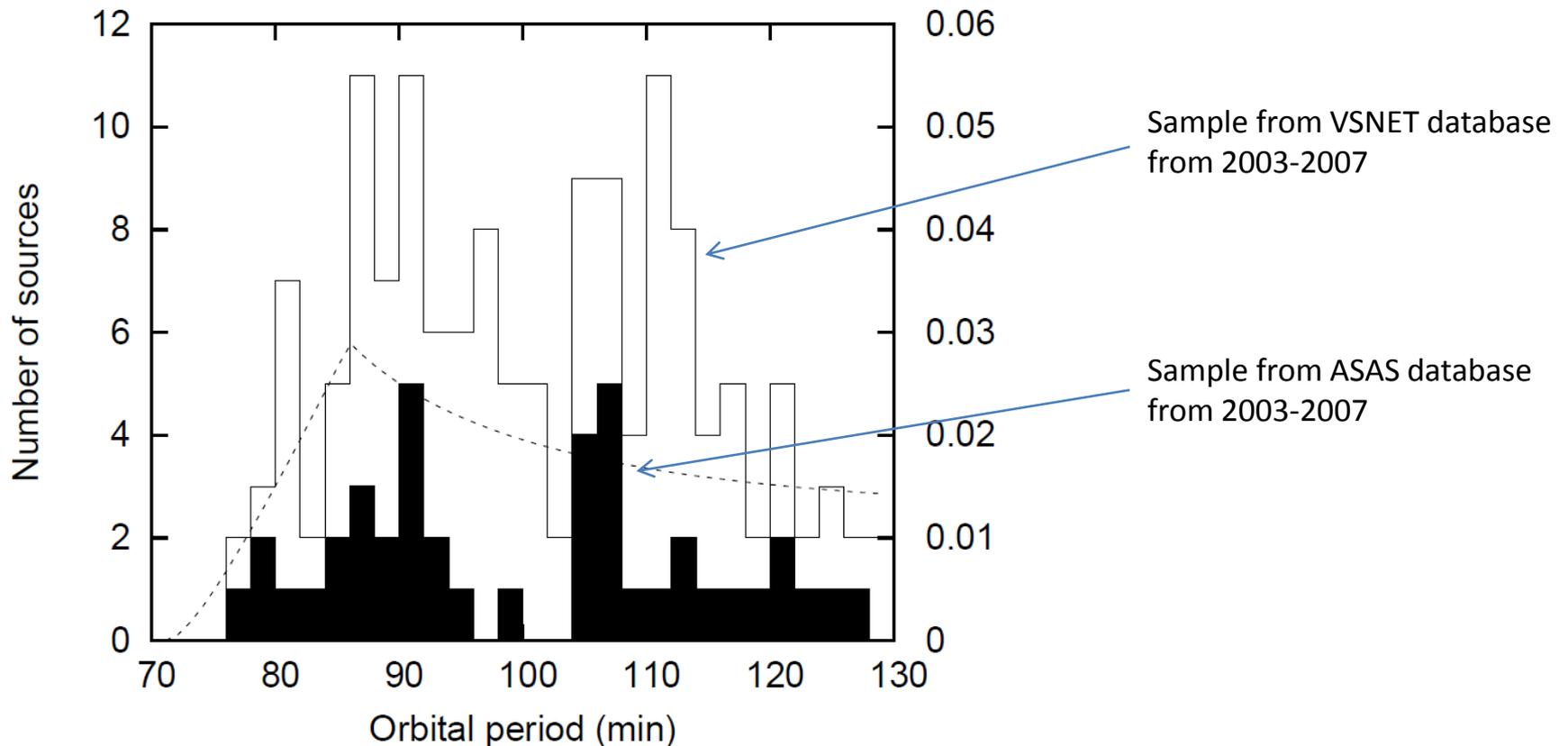
- Long recurrence time of superoutbursts (T_s) near P_{min}
- Intrinsic population of WZ Sge-type DNe (=very long T_s)???

Detectability of outbursts
= $1/T_s$
(normalized at 300 d)



Bayesian estimation of the intrinsic P_{orb} distribution of DNe

- Sample: DNe whose superoutbursts are detected in a certain period of time with a certain equipment.
 - “Observed sample” = “intrinsic distribution” x “outburst detectability”
- Estimation of the intrinsic distribution, using a Bayesian analysis



Models

- Observed distribution, Q , Outburst detectability (D_{outb}), Intrinsic distribution, I

$$Q(a, P_{\min}) = D_{\text{mag}} \cdot D_{\text{outb}} \cdot I(a, P_{\min}) / A_Q(a, P_{\min})$$

The intrinsic distribution.
(The form has no physical meaning)

$$I(p) = \begin{cases} p^{-\alpha} e^{-\alpha/p} / A_I & (p \geq 1) \\ 0 & (p < 1) \end{cases}$$
$$p = P_{\text{orb}} - P_{\min} + 1 \text{ (min)}$$

Detectability depending on
the absolute magnitude.

$$D_{\text{mag}}(P_{\text{orb}}) \propto 10^{-0.6M_V(P_{\text{orb}})}.$$

$$M_{V,\text{SU}}(P_{\text{orb}}) = C_1 - 0.383P_{\text{orb}} \quad (P_{\text{orb}} \geq P_{\text{crit}})$$
$$M_{V,\text{WZ}}(P_{\text{orb}}) = C_2 \quad (P_{\text{orb}} < P_{\text{crit}})$$

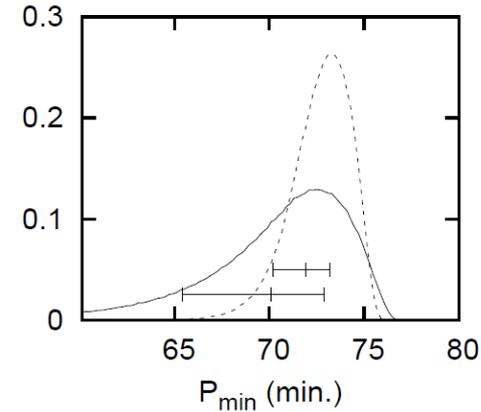
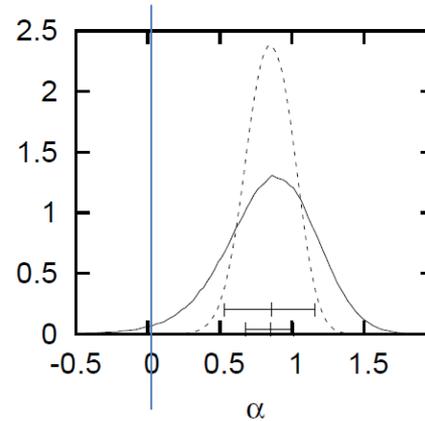
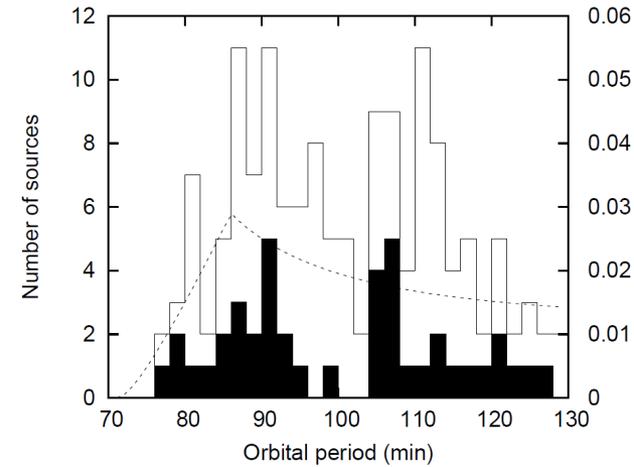
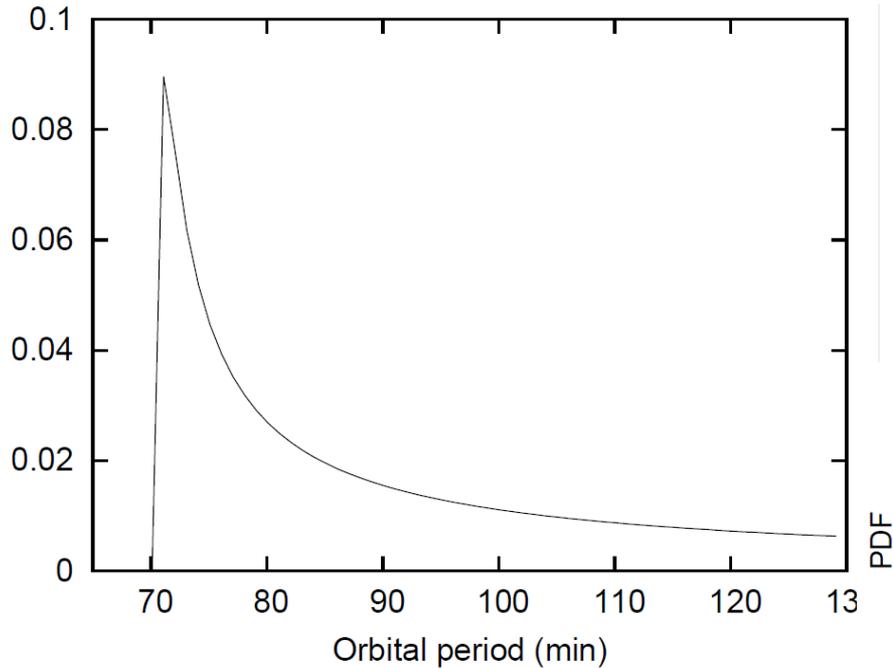
Detectability of outbursts depending on the recurrence time.

$$D_{\text{outb}}(P_{\text{orb}}) = \begin{cases} 1 & (P_{\text{orb}} \geq P_{\text{crit}}) \\ \left(\frac{P_{\text{orb}} - P_{\min}}{P_{\text{crit}} - P_{\min}} \right)^n & (P_{\text{orb}} < P_{\text{crit}}) \end{cases}$$

Estimating α and P_{\min} , with Markov-Chain Monte Carlo method

Results

- The intrinsic distribution has a “period spike” and shorter P_{\min} .



Among sources between 70—130 min,
59 % of sources are concentrated
between 70—86 min.

(Flat distribution)

Summary 1

- Our experimental approach showed that WZ Sge-type DNe could be a part of “missing population” near P_{\min} .
- The result depends on the assumed form of the intrinsic distribution.
 - Another form?
- The result should be tested by another sample, or another period of time.

In the view point of Dynamics of accretion disks

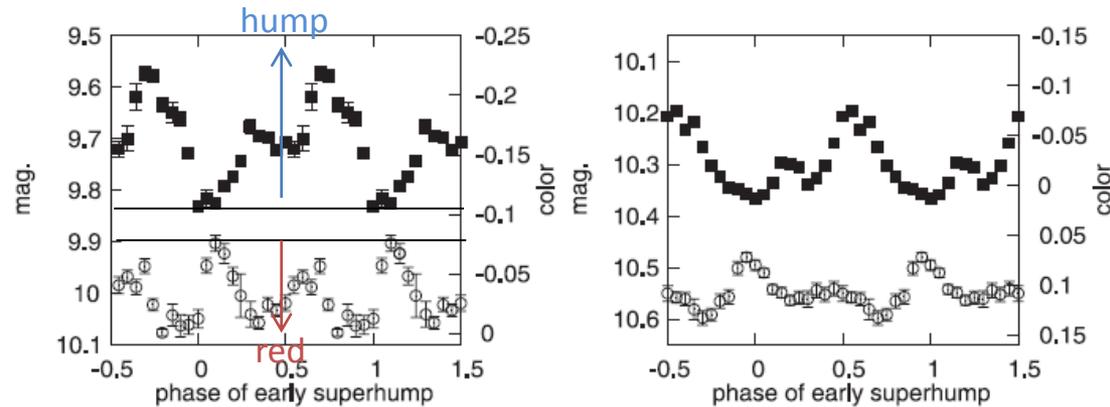
- New type of tomography:
 Early superhump mapping
(Uemura+11, in prep.)

Early superhump

- Only observed in WZ Sge-type DN
 - Period = P_{orb}
 - Doubly-peaked profile
- Rotation effect of non-axisymmetric structure of accretion disks

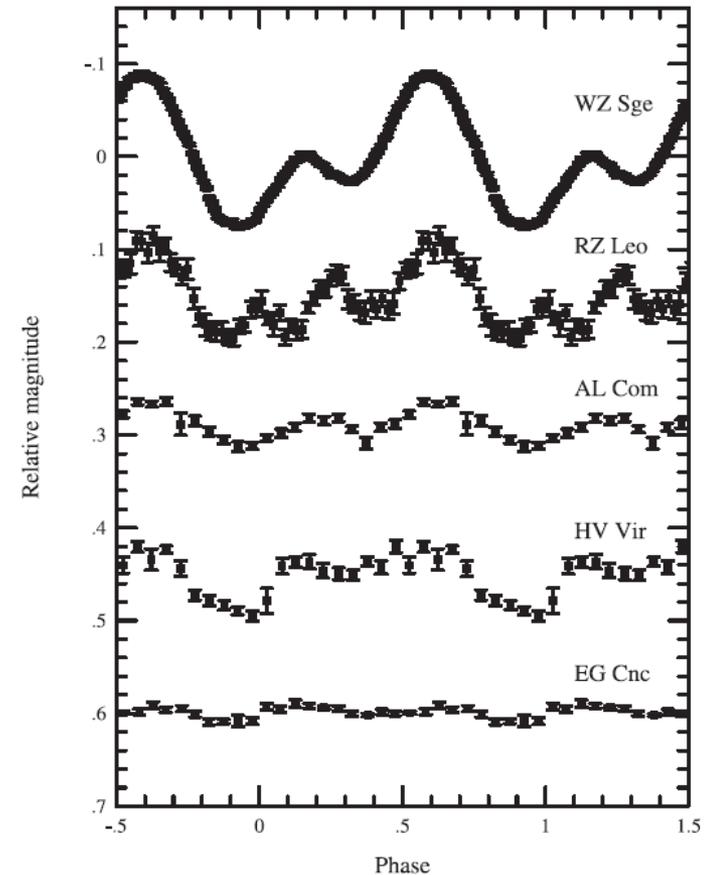
Color variation in early superhumps of V455 And (Matsui, et al. 2009)

Red color of hump component \rightarrow vertically elongated structure?



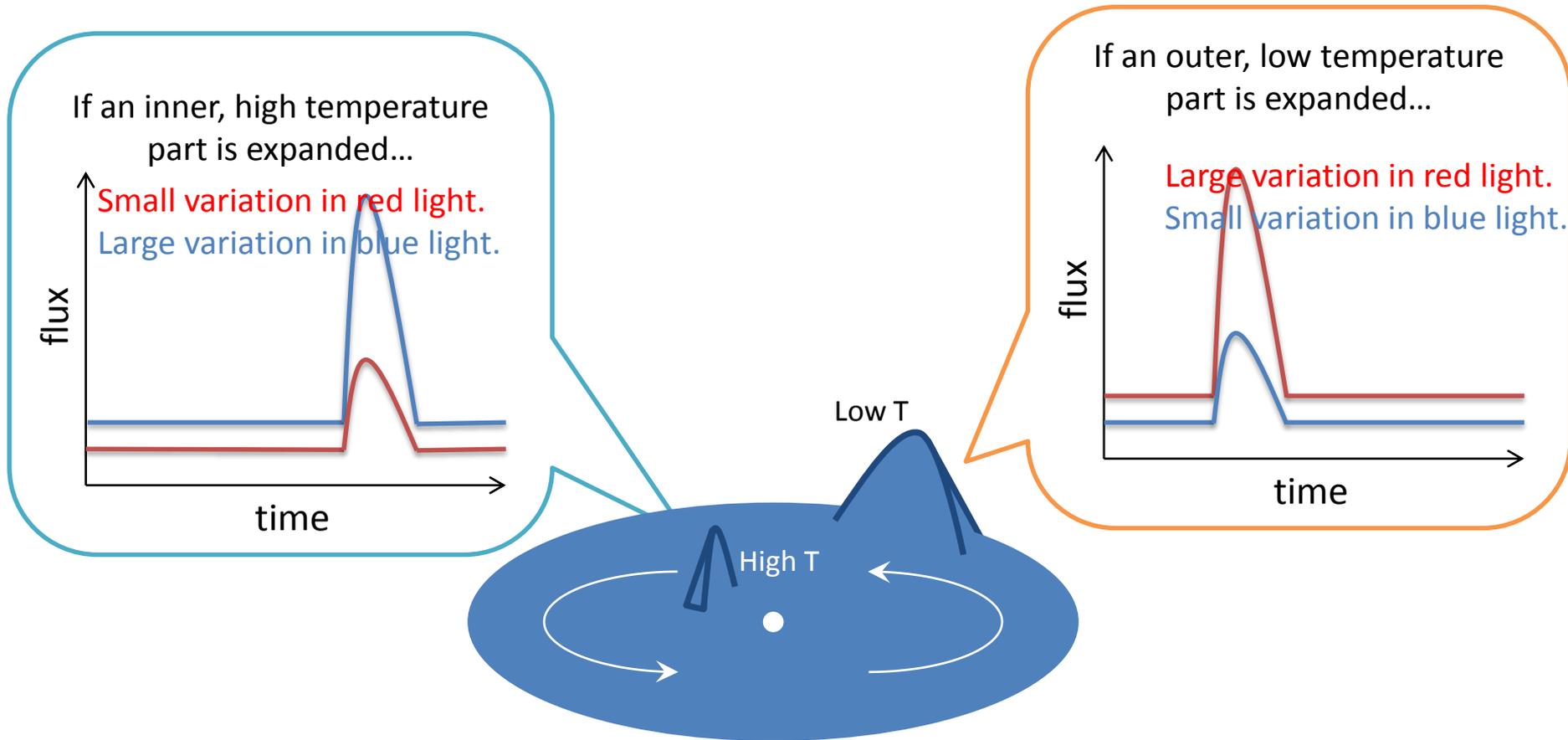
Zoo of early superhumps (Kato, et al. 2002).

Larger amplitude in edge-on systems.



Reconstruction of the accretion-disk structure from early superhumps

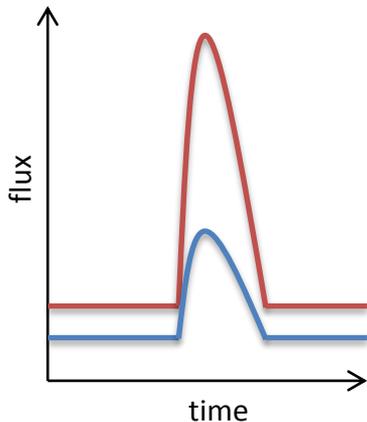
- Phase information → azimuthal structure in disks
- Color information → radial structure in disks



Details of our Bayesian model

Input

Multiband light curves



Model

Bayesian estimation of the height, $h(i,j)$

$$P(h) \propto L[f_{\nu,obs}(\phi), f_{\nu,model}(\phi)]\pi(h)$$

Posterior

likelihood

prior

* Likelihood function (defined by the observed and model LC)

$$L \propto \prod_{i,j} \exp - \frac{[f_{\nu_i,obs}(\phi_j) - f_{\nu_i,model}(\phi_j)]^2}{2\sigma^2}$$

* Prior distribution

(locally smoother)

$$\pi_{smooth}(h) \propto \prod_{l,m} [\exp - \frac{(h_{l,m} - 2h_{l-1,m} + h_{l-2,m})^2}{2w^2} \exp - \frac{(h_{l,m} - 2h_{l,m-1} + h_{l,m-2})^2}{2w^2}],$$

(default image to be $h=0.1r$)

$$\pi_{disk} \propto \begin{cases} \prod_{l,m} \exp - \frac{(h_{l,m} - h_{disk,l,m})^2}{2h_{disk,l,m}^2} & (h_{l,m} \geq 0) \\ 0 & (h_{l,m} < 0) \end{cases}$$

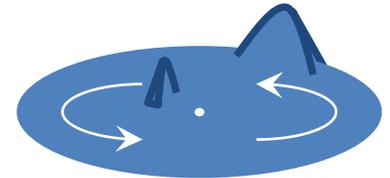
* Estimation of “h” is done with Markov-chained montecarlo (MCMC)

* The temperature distribution is like an standard disk model, as

$$T = T_{in} \left(\frac{r}{r_{in}} \right)^{-3/4}$$

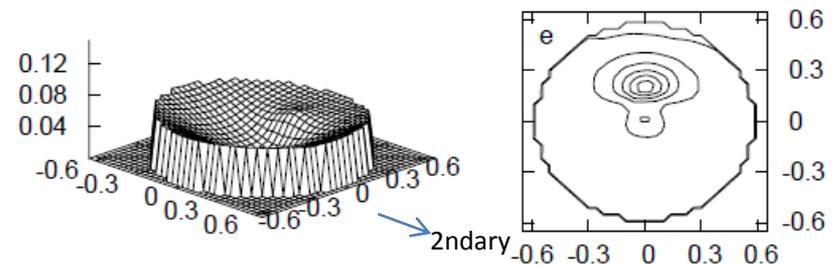
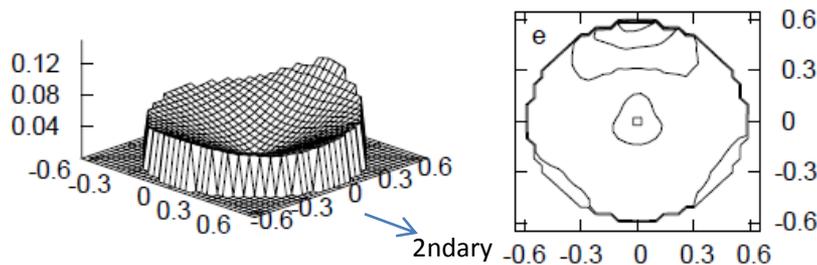
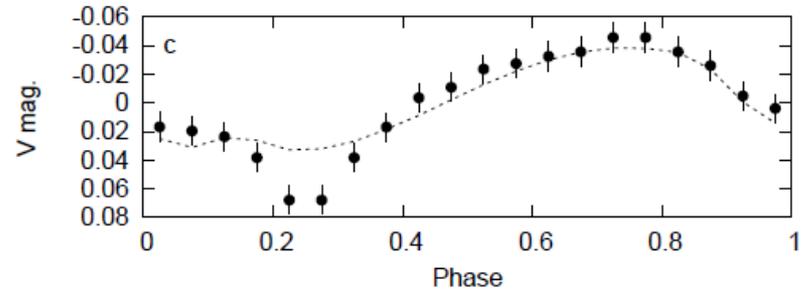
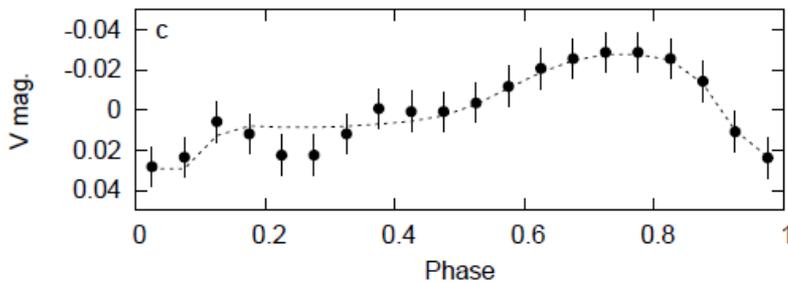
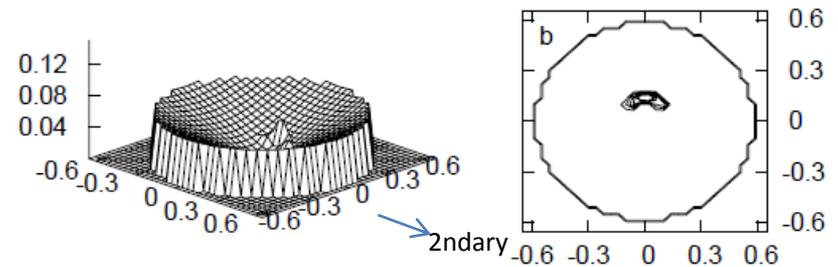
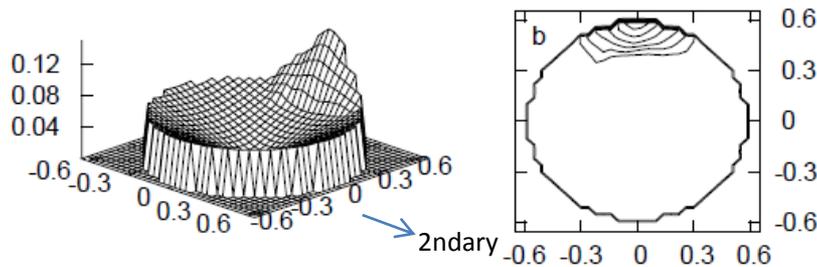
Output

Height map of disk



Demonstration with artificial data sets

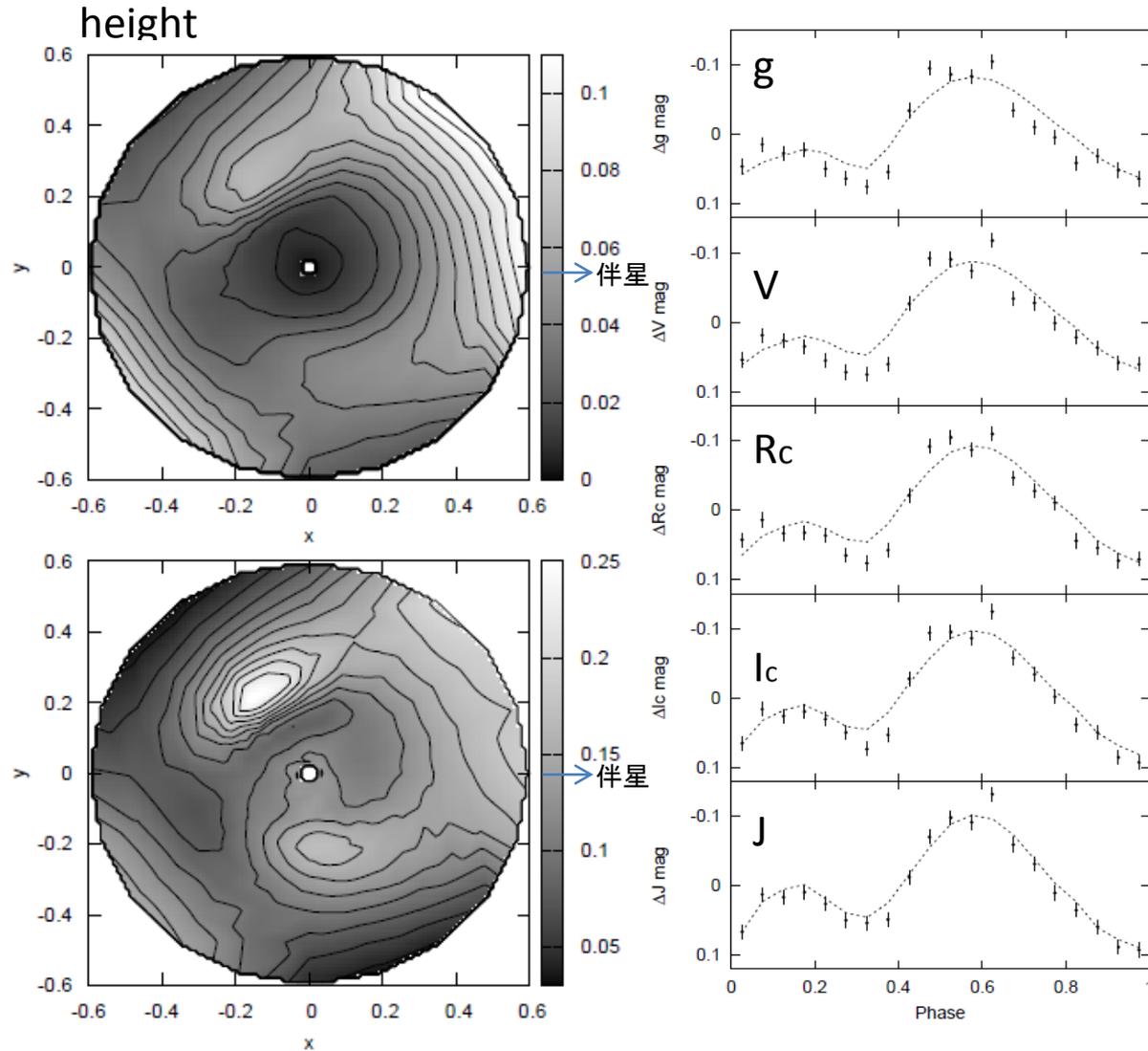
- Working as expected
 - Outer structures are reconstructed to be outer, and vice versa
 - Smoother structures than assumed, due to the prior distribution.



Reconstruction of the disk from early superhumps in V455 And: The 5th day of the superoutburst

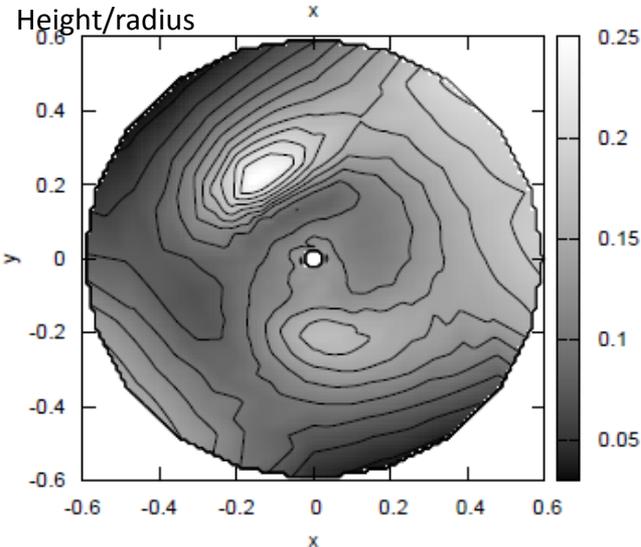
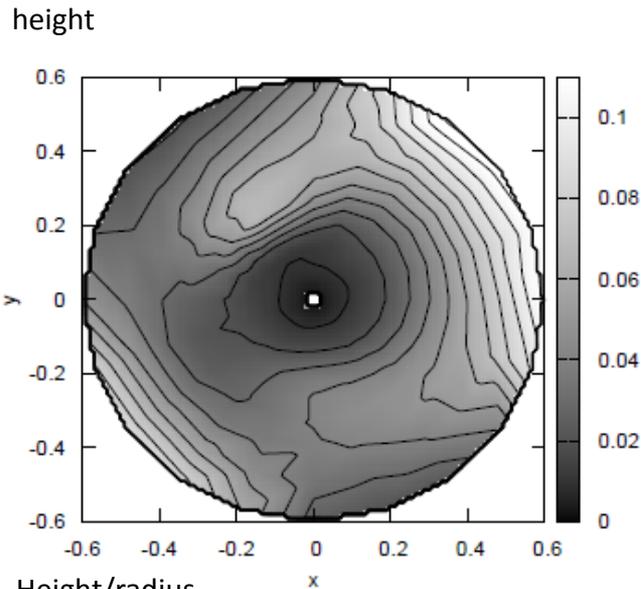
- 5-band data (g,V,R,I,J)
- Flaring outermost parts making primary and secondary maxima of the light curve
- “arm”-like structures

Height/radius

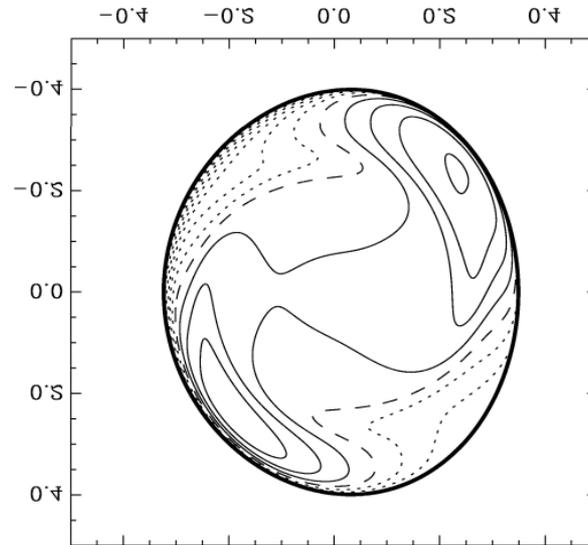


Comparison with theoretical models

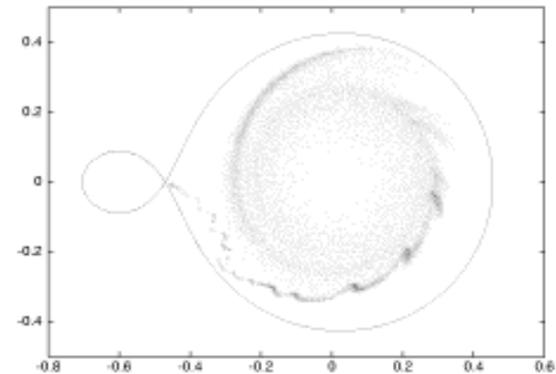
- Tidal distortion? 2:1 resonance?



Theoretical image of the height distortion by the tidal effect (Ogilvie 2002)



SPH simulation of the 2:1 (Kunze & Speith 2005)



Reconstructed disks are similar to the disk distorted by the tidal effect, but the upper-left flaring part cannot be explained.

Summary 2

- We have developed a Bayesian model which reconstruct the height map of disks from the multi-band light curves of early superhumps.
- The reconstructed disk has flaring outermost parts making the primary and secondary maxima of the light curve, and “arm”-like structures.
- The structure is similar to the disk structure distorted by the tidal effect, but the part for the secondary minimum cannot be explained.
- Future work:
 - Does the model make similar results for another objects?
 - How the disk evolves with time?
 - Can the reconstructed disk explain the emission-like profile in WZ Sge stars?