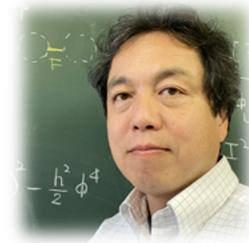


Modified Theories of Gravity and their Phenomenological Consequences

Tomohiro Inagaki
Hiroshima University



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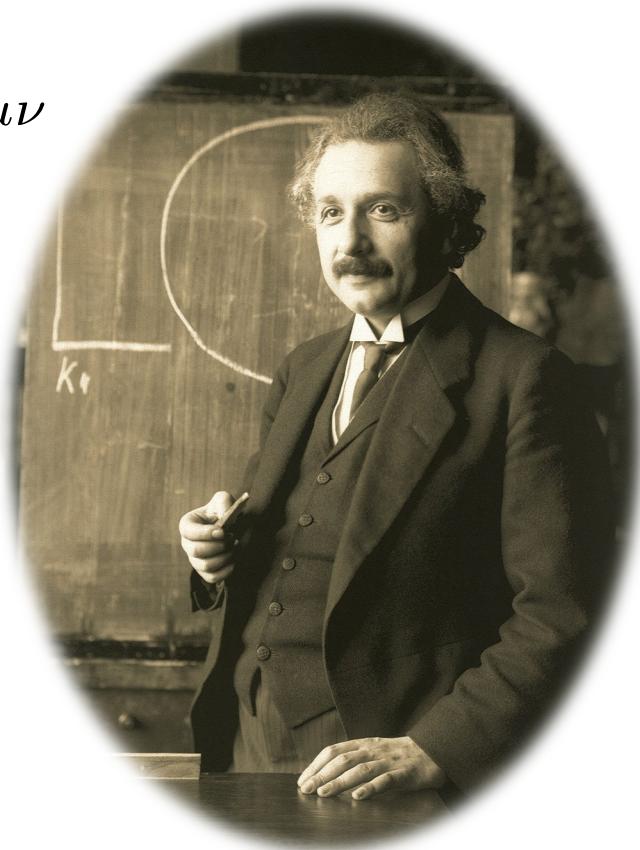
- Introduction
- Modified theories of Gravity
- Phenomenological consequences
- Conclusion

Introduction

Why?

Einstein's General Relativity

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$



Why?

Einstein's General Relativity

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

- Phenomenological consequences
 - Apsidal precession of the planet Mercury,
 - Dense stars, Black holes,
 - Gravitational lens,
 - Gravitational wave,
 - Expansion of the universe,

...

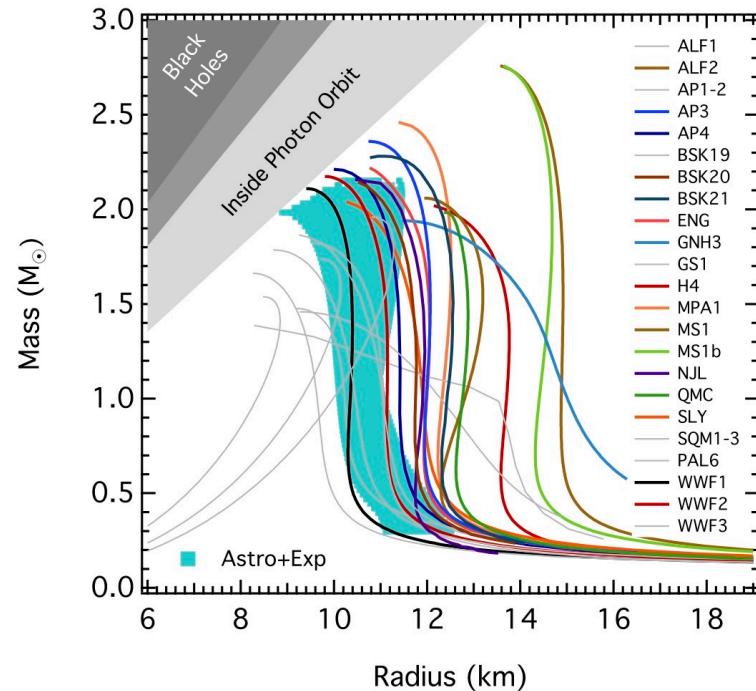
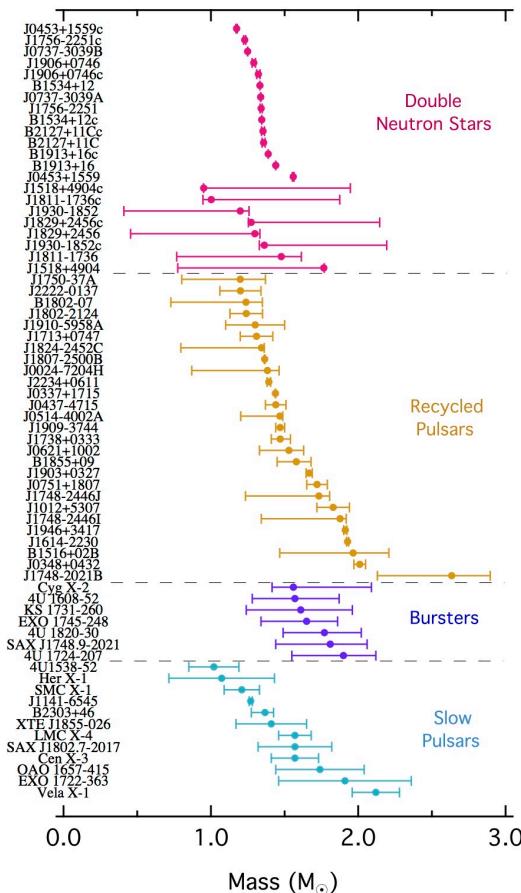
Why?

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- Phenomenological consequences
 - Apsidal precession of the planet Mercury,
 - Dense stars, Black holes, ←Heavy neutron stars
 - Gravitational lens, ←Dark matter
 - Gravitational wave, ←Just started
 - Expansion of the universe, ←Accelerated expansion
 - ... ←Small scale < 0.01mm

Heavy neutron stars



Ozel & Freire 2016

<http://xtreme.as.arizona.edu/NeutronStars/>

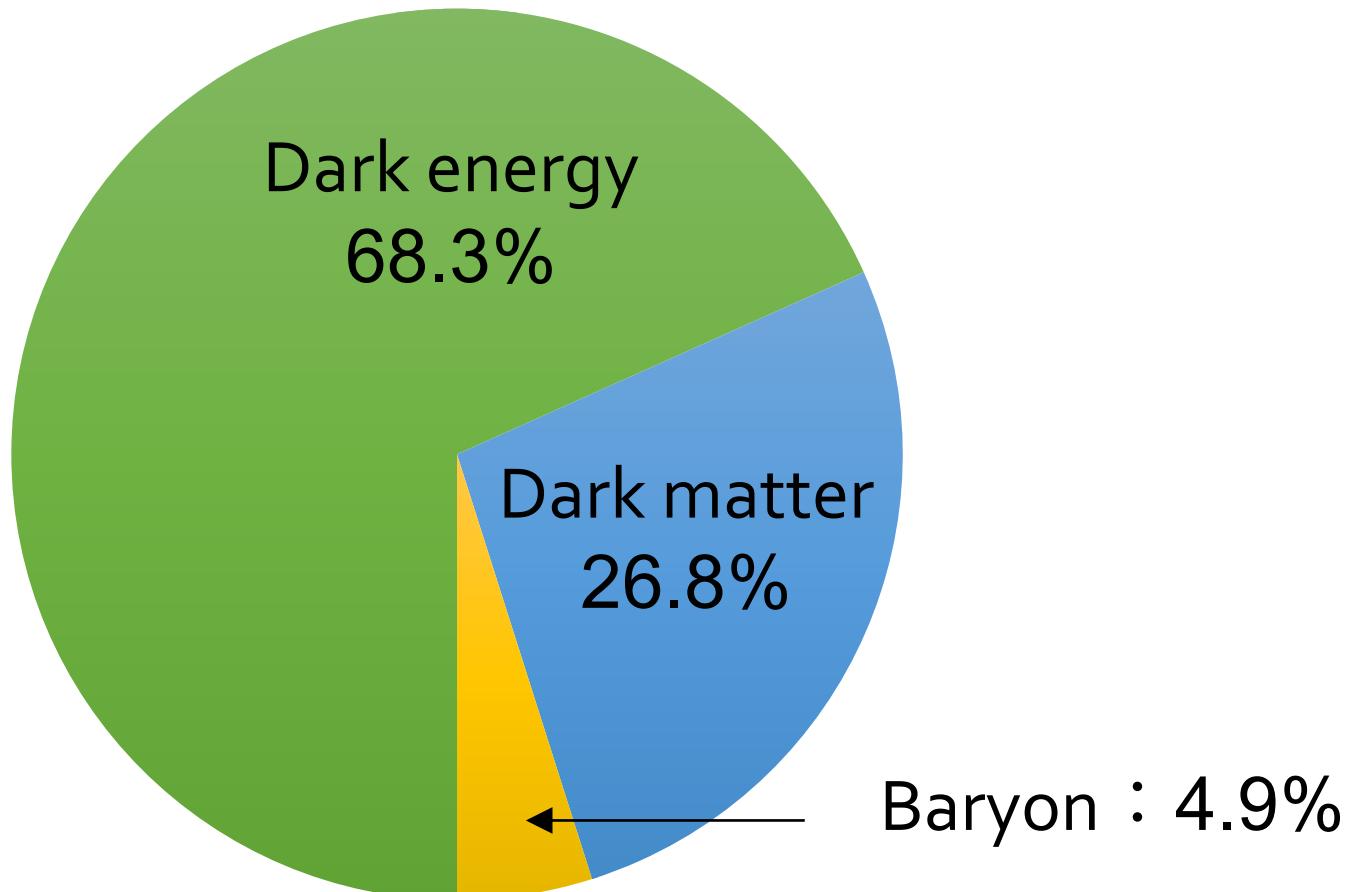
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Dark matter and dark energy



How?

Einstein equation

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Action

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$

How?

Einstein equation

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Action

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Modified Gravity

- Higher order
- Non-local
- Gauss-Bonnet
- Torsion
- ...

How?

Einstein equation

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

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Particle physics

- Neutrino
- Axion
- Superpartner
- ...

How?

Einstein equation

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Spacetime

- Extra dimensions
- D-brane
- ...

Action

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$

Modified Theories of Gravity

Modification of the Einstein eq.

Einstein equation

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Action

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$

Modified Gravity

- Higher order
- Non-local
- Gauss-Bonnet
- Torsion
- ...

$F(R)$ gravity

- Extension of the Einstein Hilbert action

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$



$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} F(R) + S_{\text{matter}}$$

Starobinsky model

A. A. Starobinsky, 1979

- Action

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R + \kappa^2 R^2) + S_{\text{matter}}$$

- After the Weyl transformation of the metric tensor

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-\tilde{g}} \tilde{R} + S_{\text{matter}}$$

$$+ \int d^4x \sqrt{-\tilde{g}} \left[-\frac{1}{2} \tilde{g}^{\mu\nu} (\partial_\mu \phi)(\partial_\nu \phi) - V(\phi) \right]$$

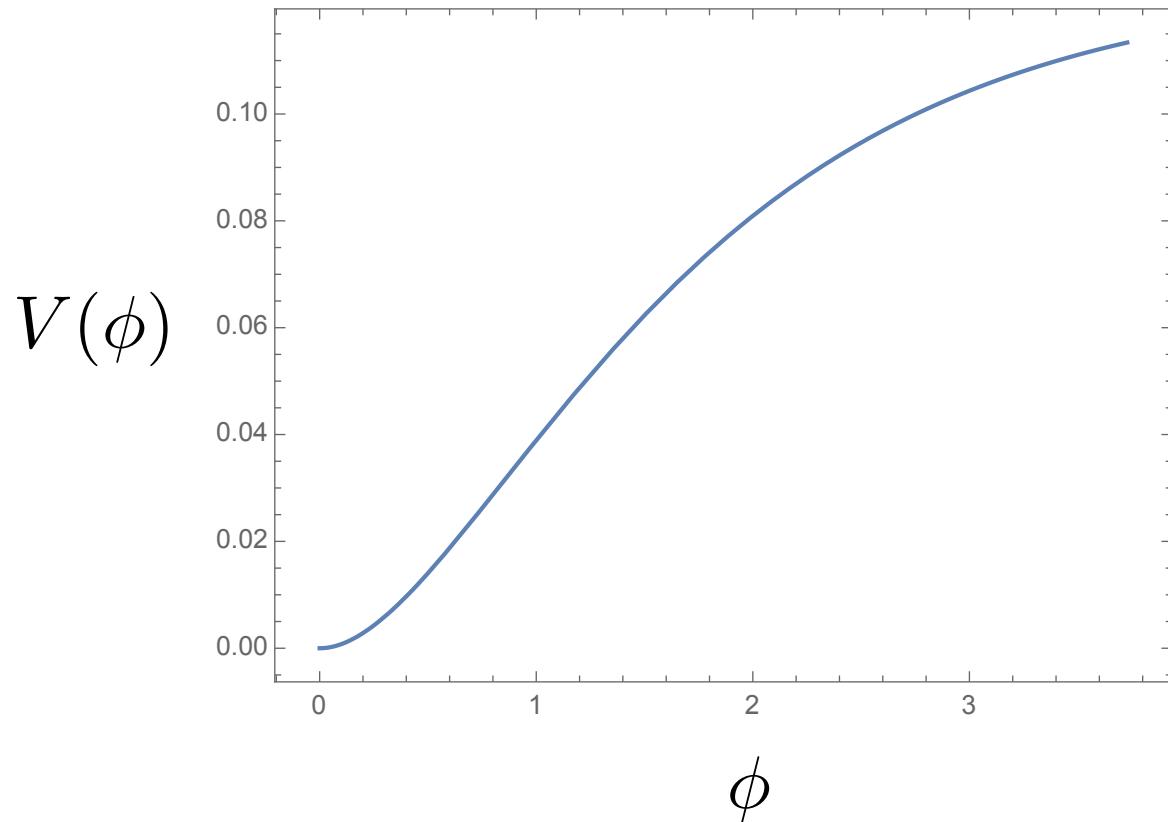
Scalarmon



Starobinsky model

A. A. Starobinsky, 1979

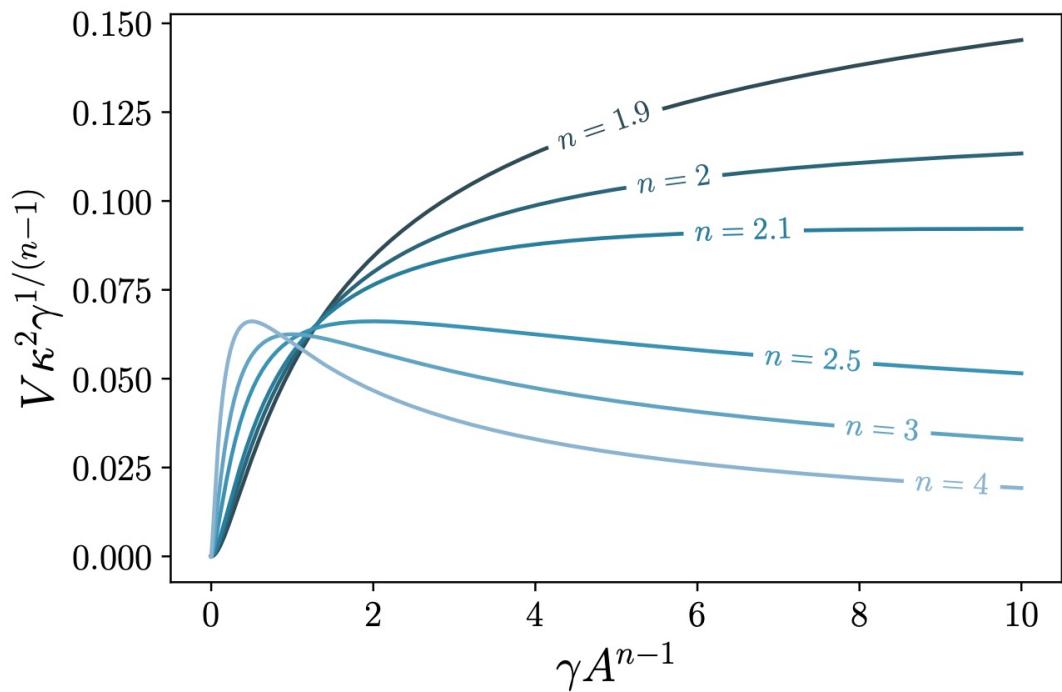
- Scalaron potential



More general cases

- Power law model

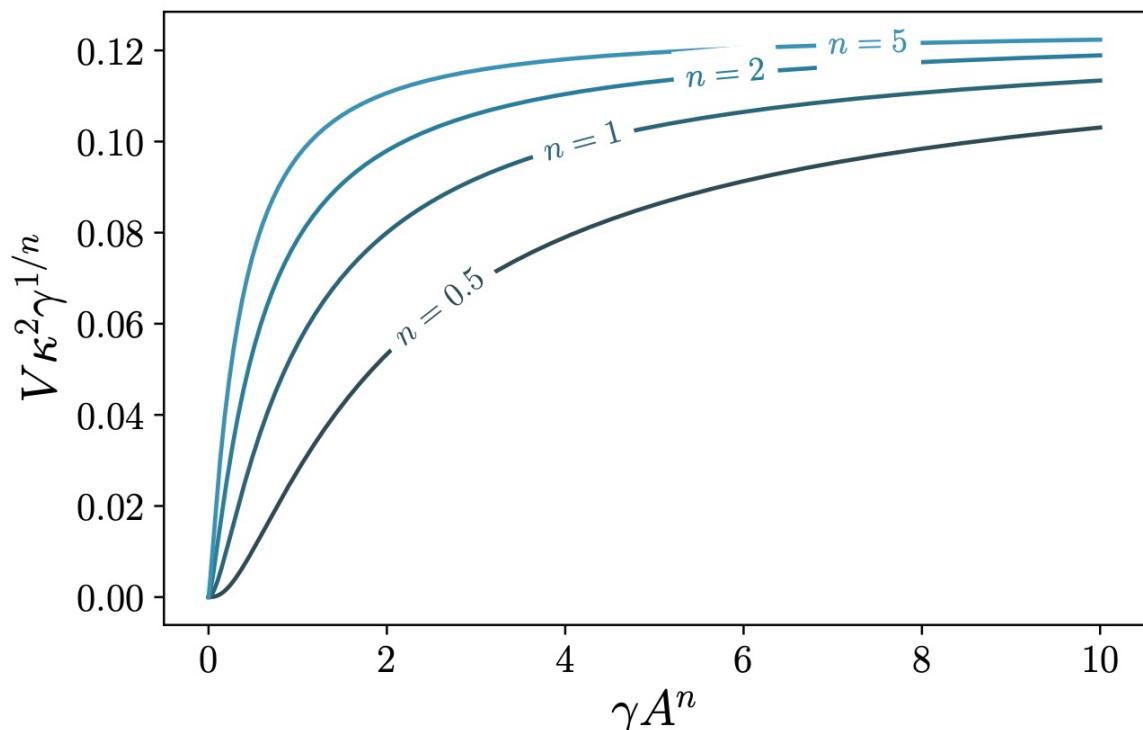
$$F(R) = R + \gamma R^n$$



More general cases

- Modified power law model

$$F(R) = R(1 + \gamma R^n)^{1/n}$$



$F(G)$ gravity

- Gauss-Bonnet invariant

$$G = R^2 - 4R^{\mu\nu}R_{\mu\nu} + R^{\mu\nu\rho\sigma}R_{\mu\nu\rho\sigma}$$

- Extension of the Einstein Hilbert action

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$



$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R + F(G)) + S_{\text{matter}}$$

S. Nojiri and S. D. Odintsov, Phys. Lett. B 631, (2005)

$F(R, G)$ gravity

- Extension of the Einstein Hilbert action

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$



$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} F(R, G) + S_{\text{matter}}$$

Cartan formarism

- Vierbein

$$g_{\mu\nu} = \eta_{ij} e_\mu^i e_\nu^j$$

- Fermionic fields are described in a local Lorentz frame.
- Torsion tensor is naturally introduced.

$$T^\lambda{}_{\mu\nu} = \Gamma^\lambda{}_{\mu\nu} - \Gamma^\lambda{}_{\nu\mu}$$

T. W. B. Kibble (1961), D. W. Sciama(1962)

Modification of the Einstein eq.

Einstein equation

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Particle physics

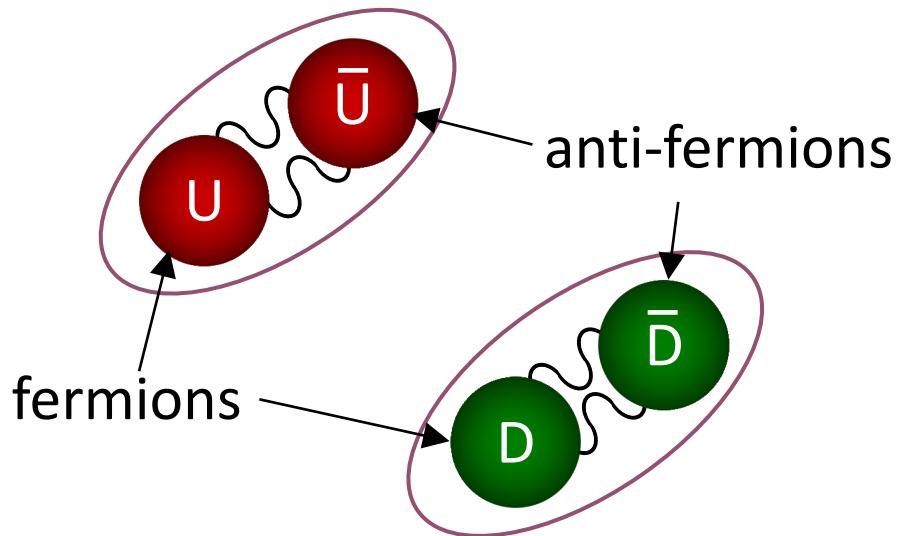
- Neutrino
- Axion
- Superpartner
- ...

Action

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$

Gauged NJL model

- Low energy effective theory of light pseudo-scalar mesons constructed by quarks and anti-quarks
- We scale up the model from the QCD scale to the inflation scale



T. I., S. D. Odintsov and H. Sakamoto, Astr. Space Sci. (2015),
T. I., S. D. Odintsov and H. Sakamoto, Nucl. Phys. B (2017),
T. I., S. D. Odintsov and H. Sakamoto, Europhys. Lett. (20170).

Modification of the Einstein eq.

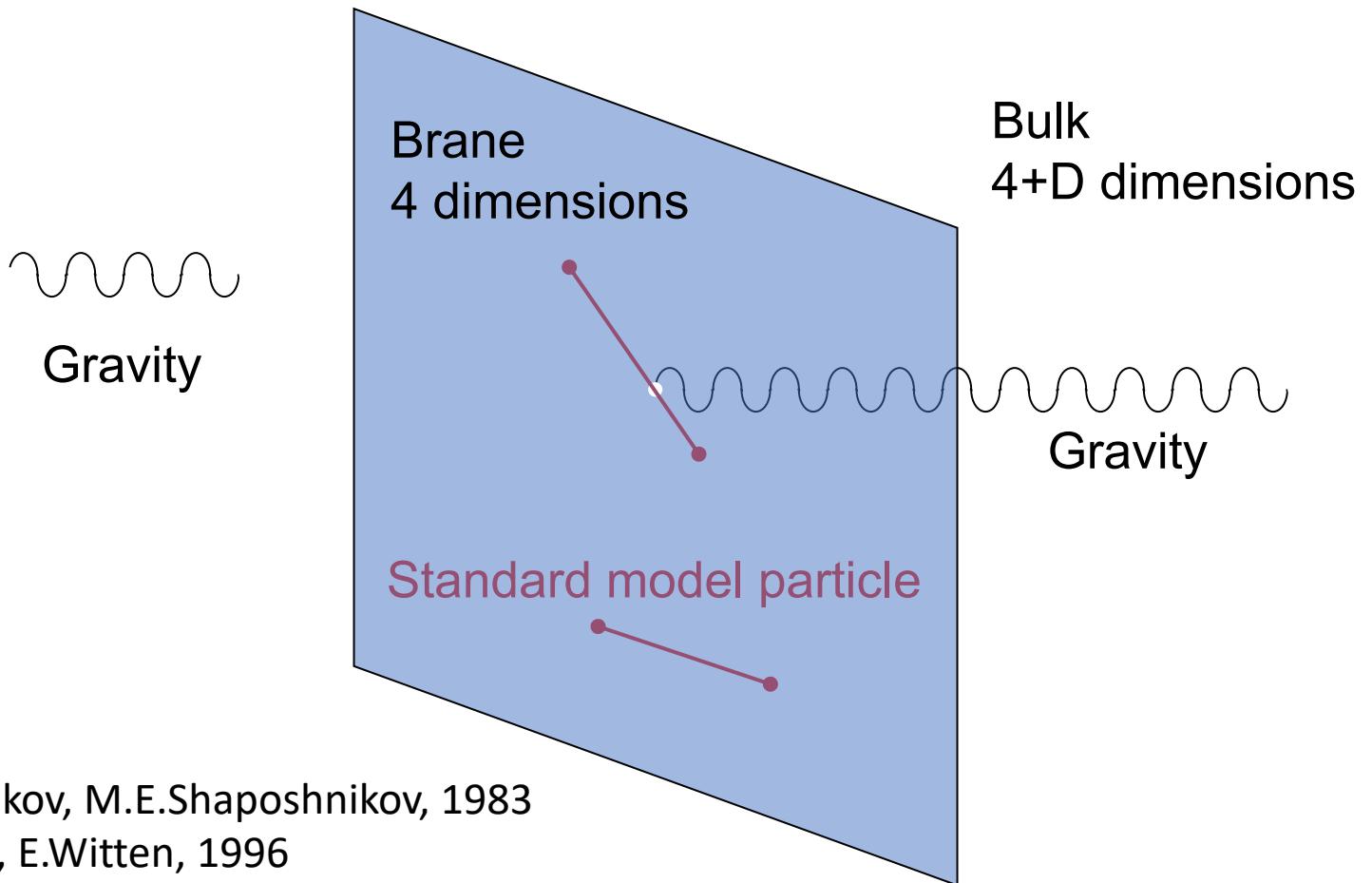
Einstein equation

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Action

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$

Brane world



V.A.Rubakov, M.E.Shaposhnikov, 1983

P.Horava, E.Witten, 1996

N.Arkani-Hamed, S.Dimopoulos, G.Dvali, 1998

L.Randall, R.Sundrum, 1999

Phenomenological consequences

Origin of accelerated expansion

- In a homogeneous and isotropic spacetime

$$ds^2 = c^2 dt^2 + a^2(t)(dx^2 + dy^2 + dz^2)$$

- Sources of energy density

Radiation	$a(t) \propto t^{1/2}$
Matter	$a(t) \propto t^{2/3}$
Potential energy	$a(t) \propto \exp(\alpha t)$
Cosmological constant	

Quasi de-Sitter expansion

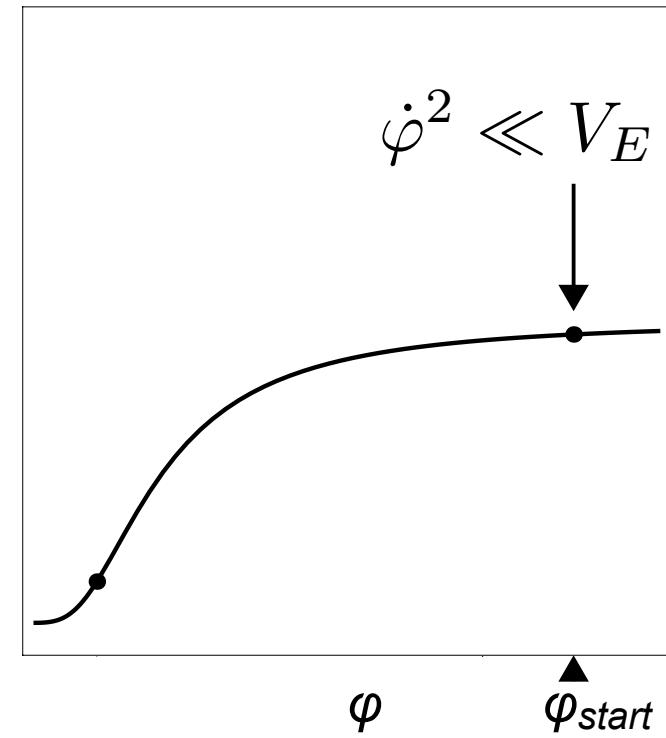
- Friedman equation

$$3 \left(\frac{\dot{a}}{a} \right)^2 = \frac{1}{2} \dot{\varphi}^2 + V_E$$

- Assumption $\dot{\varphi}^2 \ll V_E$



$$a(t + \Delta t) \sim a(t) e^{\sqrt{\frac{V_E}{3}} \Delta t}$$



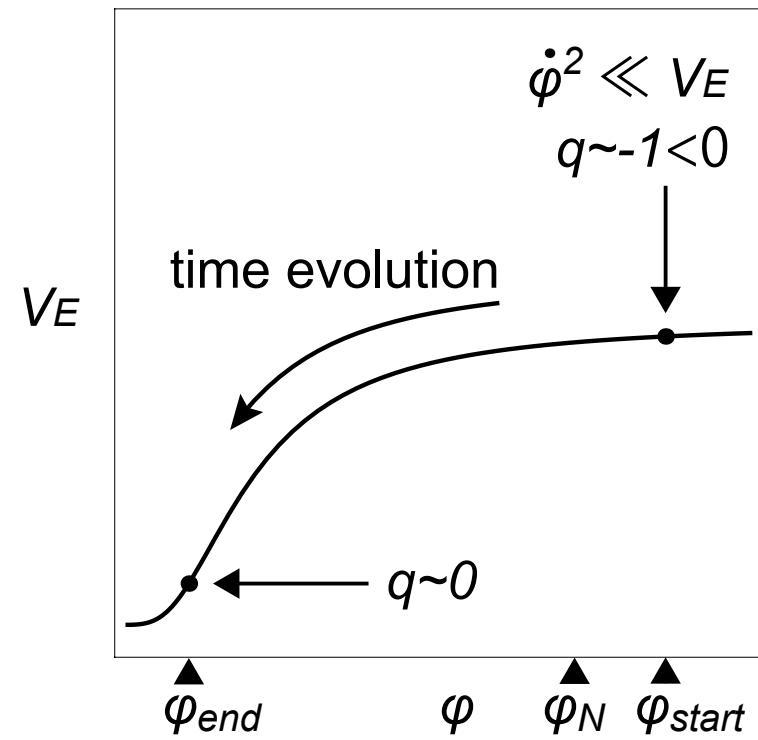
Exit from Inflation

- Equation of motion

$$\ddot{\varphi} + 3\frac{\dot{a}}{a}\dot{\varphi} = -\frac{\partial V_E}{\partial \varphi}$$

- Deceleration parameter

$$q \equiv -\frac{a\ddot{a}}{\dot{a}^2} \rightarrow 0$$



Exit from Inflation

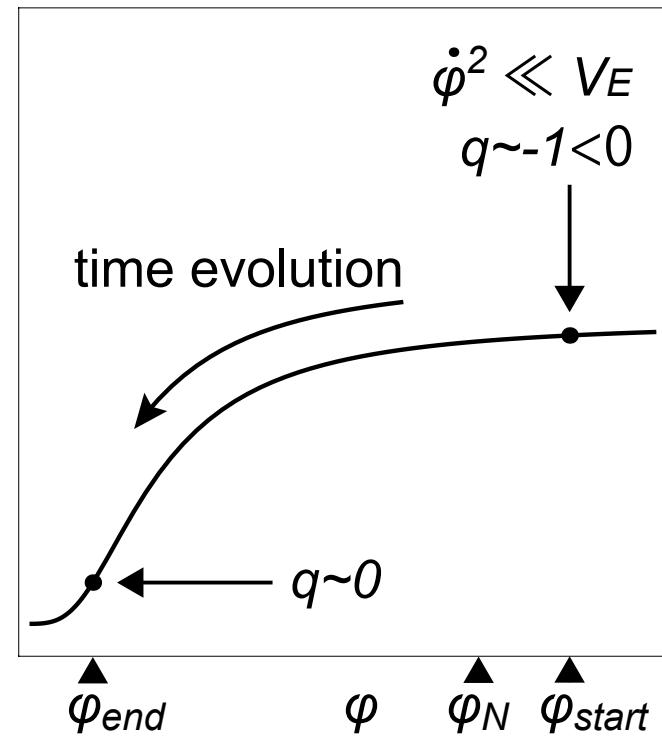
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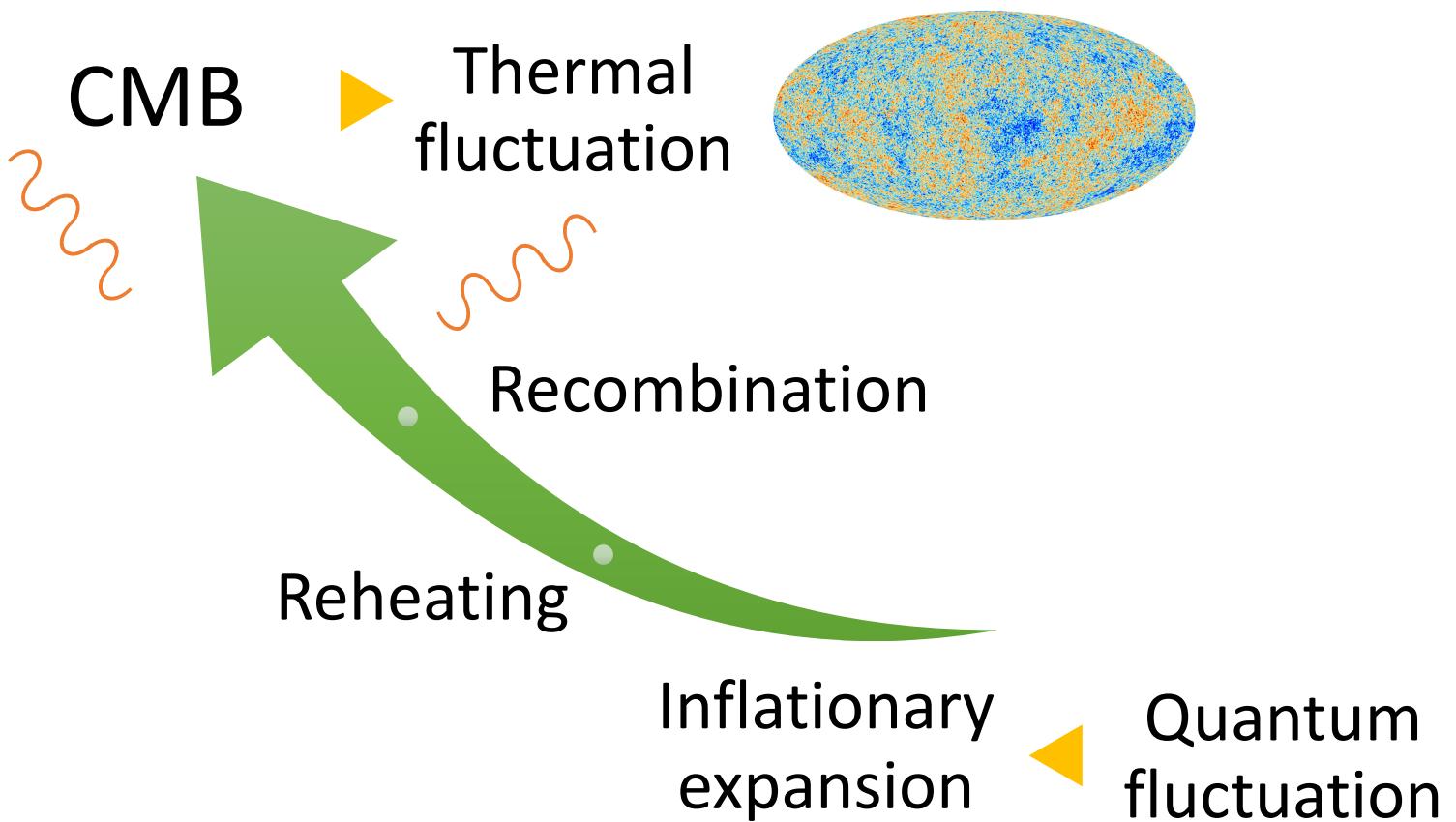
- Deceleration parameter

$$q \equiv -\frac{a\ddot{a}}{\dot{a}^2} \rightarrow 0$$

Modified gravity: Scalaron
gNJL: Composite scalar



Evidence for Inflation



Quantum fluctuations

$$\begin{aligned}\varphi + \delta\varphi \\ \rightarrow \mathcal{P}_s(k)\end{aligned}$$

Scalar type fluctuation
Origin: quantum
fluctuation of scalar field

Tensor type fluctuation
Origin: quantum
fluctuation of space-time

$$\begin{aligned}g^{\mu\nu} + \delta h^{\mu\nu} \\ \rightarrow \mathcal{P}_t(k)\end{aligned}$$

Observed CMB fluctuations

- Rescaled scalar type fluctuation
- Tensor to scalar ratio

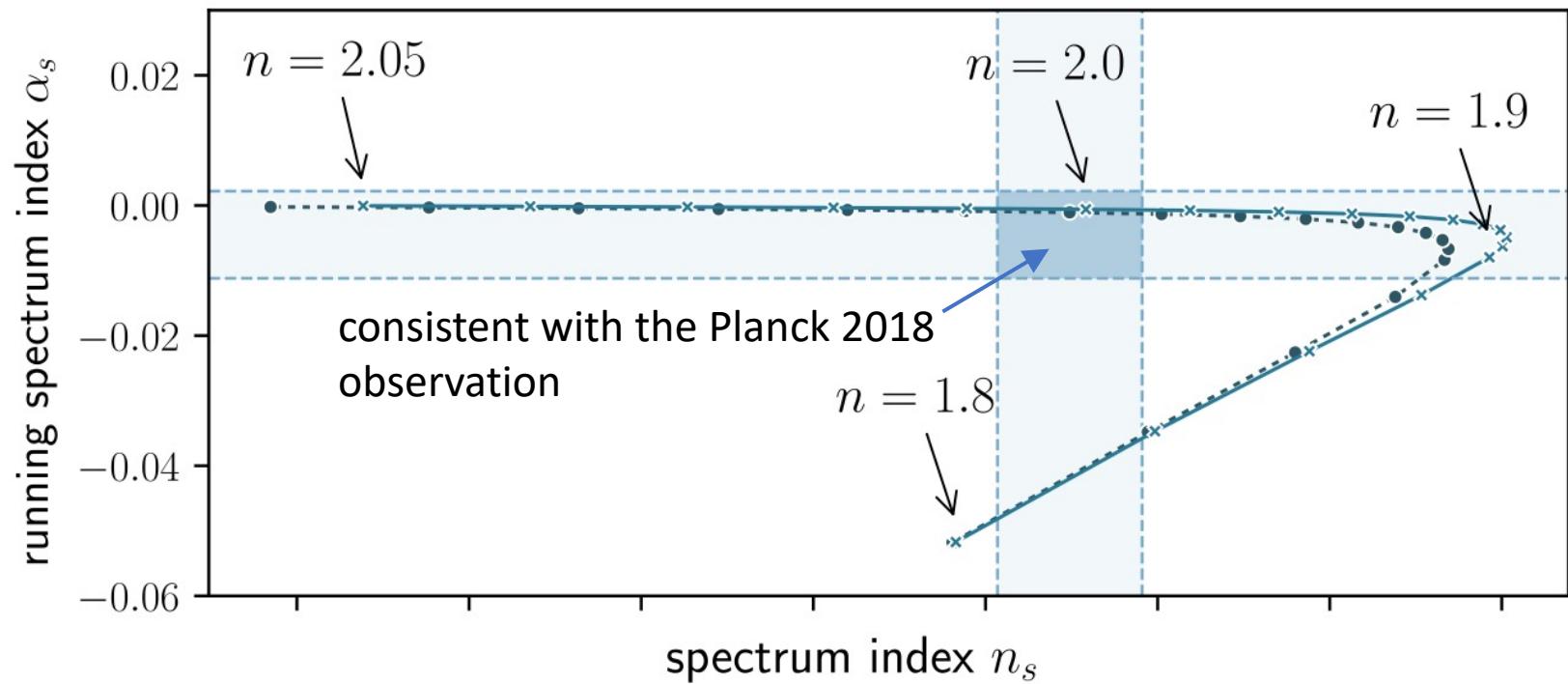
$$\mathcal{P}_s(k) \equiv A_s \left(\frac{k}{k_0} \right)^{n_s - 1}$$

$$r \equiv \frac{\mathcal{P}_t(k)}{\mathcal{P}_s(k)}$$

- Rescaled tensor type fluctuation

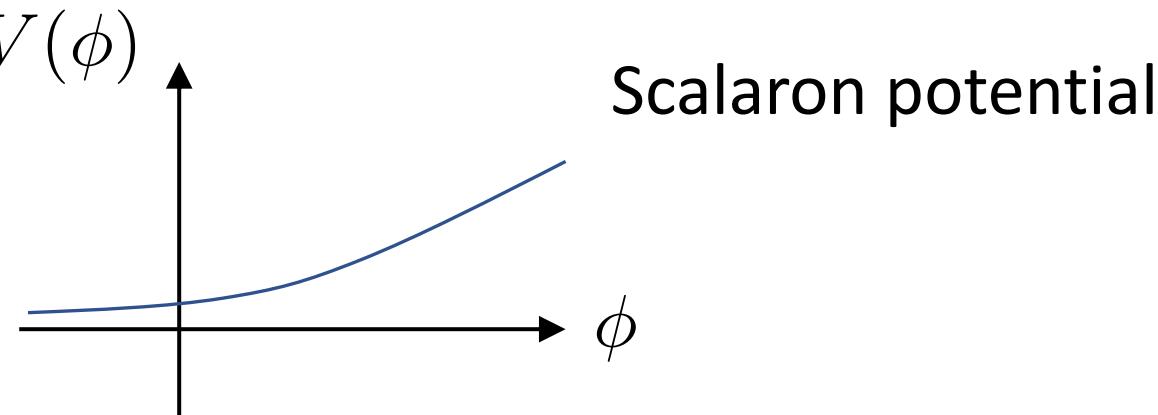
$$\mathcal{P}_t(k) \equiv A_t \left(\frac{k}{k_0} \right)^{n_t}$$

Power law model $F(R) = R + \gamma R^n$



Screening mechanism

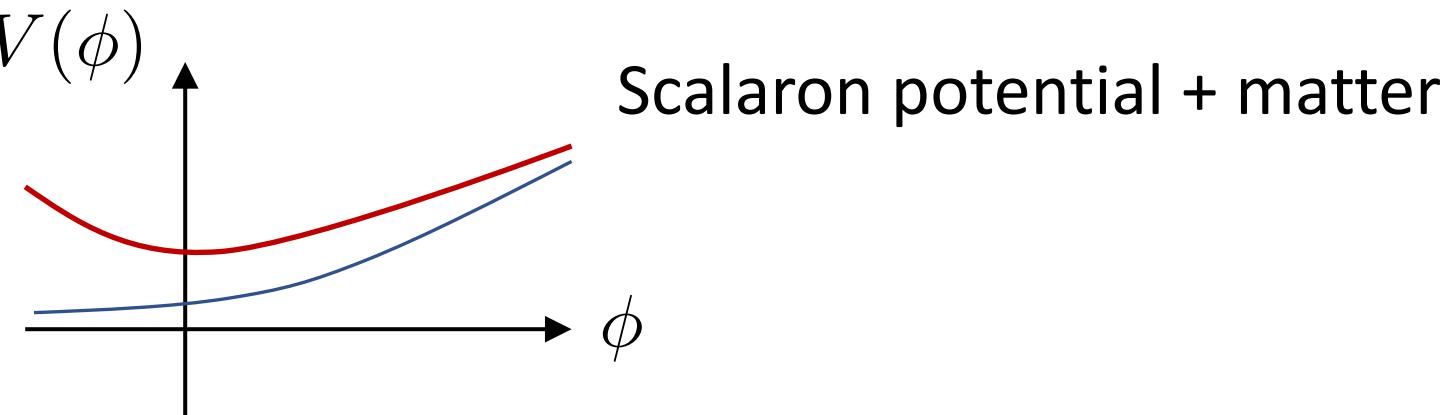
- Modified theories of Gravity should coincide with the general relativity around us. Some screening mechanism are proposed.
- Chameleon mechanism:



J. Khouri and A. Weltman, 2004,
P. Brax, C. van de Bruck, A. C. Davis and D. J. Shaw, 2008,
T. Katsuragawa and S. Matsuzaki, 2017

Screening mechanism

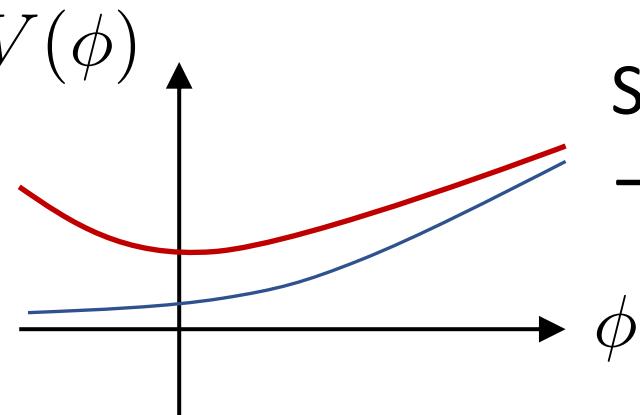
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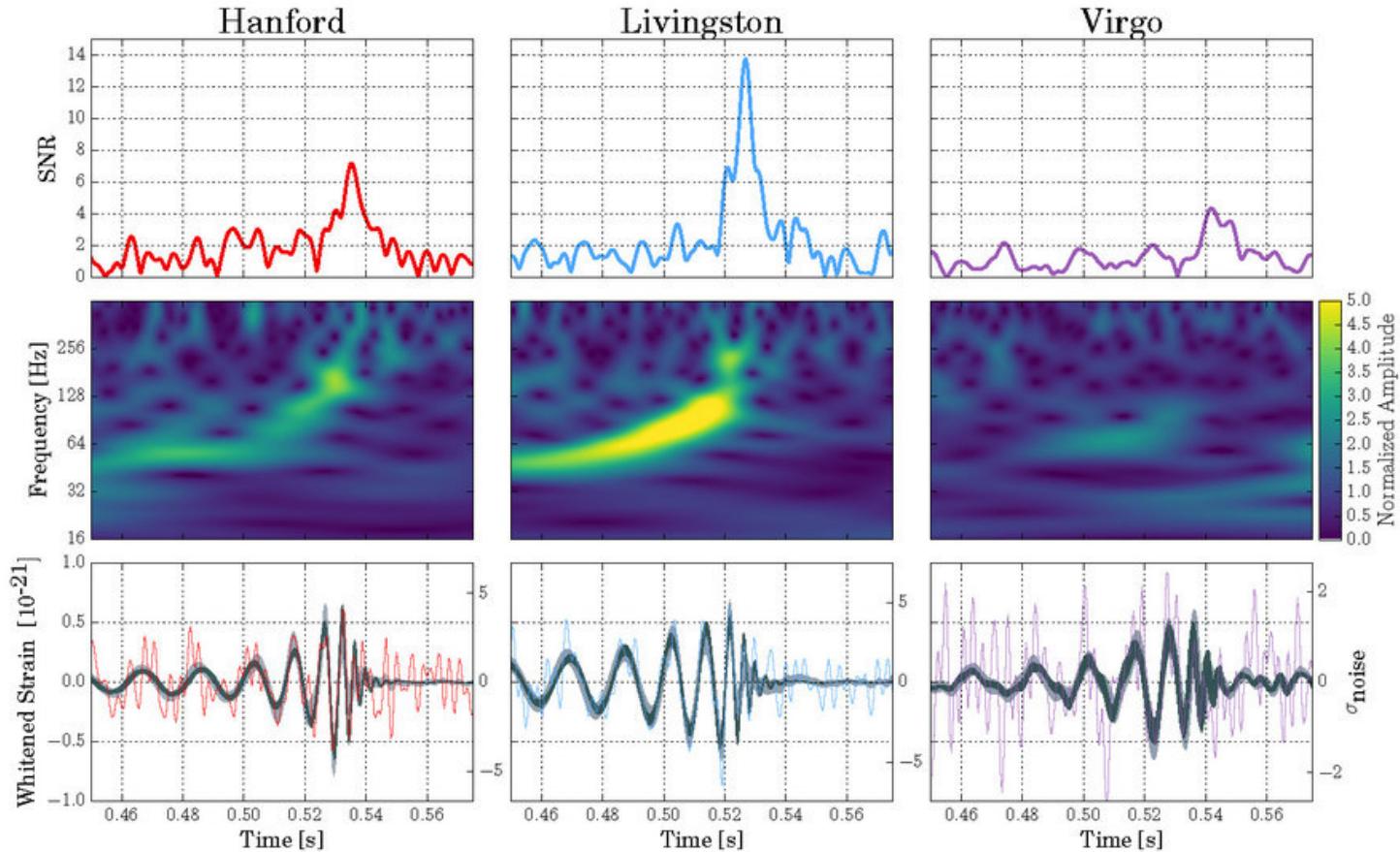
Screening mechanism

- Modified theories of Gravity should coincide with the general relativity around us. Some screening mechanism are proposed.
- Chameleon mechanism:



Scalarm potential + matter
→ Scalaron becomes heavy and screened out.
Dark matter candidate??

Gravitational waves



Gravitational wave signatures detected by each observatory
LIGO/Caltech/MIT/LSC 2017

Gravitational waves

- General relativity:
two massless tensor modes
- $F(R)$ gravity:
two massless tensor + one massive scalar mode

$$[\square - m_{F(R)}^2] \delta\Phi = 0, \quad m_{F(R)}^2 = \frac{1}{3} \left(\frac{F'(\tilde{R})}{F''(\tilde{R})} - \tilde{R} \right)$$

Y. S. Myung, 2016,

Y. Gong and S. Hou, 2018,

F. Moretti, F. Bombacigno and G. Montani, 2019,

T. Katsuragawa, T. Nakamura, T. Ikeda and S. Capozziello, 2019.

$F(G)$ gravity case

- Einstein equation in De Sitter background

$$g_{\mu\nu} = \tilde{g}_{\mu\nu} + h_{\mu\nu},$$

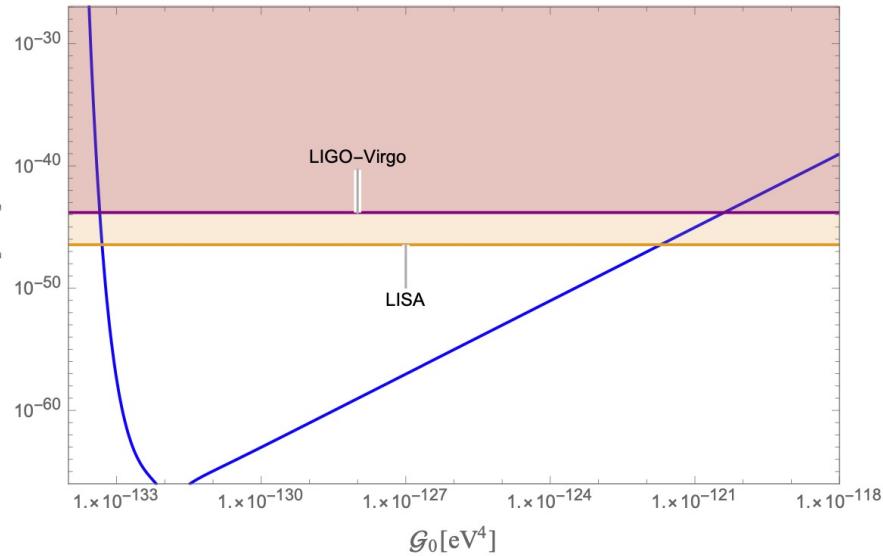
We obtain

$$\frac{M_{\text{pl}}^2}{4} \left[\square h_{\mu\nu}^T - \frac{\tilde{R}}{6} h_{\mu\nu}^T \right] - \frac{\tilde{R}}{3} \left[\nabla_\mu \nabla_\nu - \frac{m_{F(\mathcal{G})}^2}{4} \tilde{g}_{\mu\nu} \right] F''(\tilde{\mathcal{G}}) \delta \mathcal{G} = 0$$

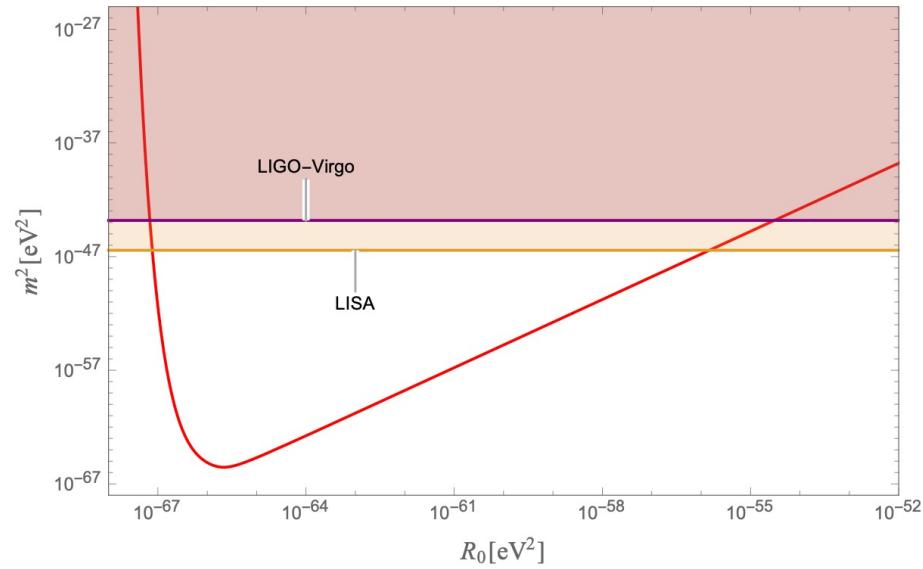
Massless tensor modes

Massive scalar mode

Massive scalar mode



$$F(R) = R - 2\Lambda(1 - e^{-R/R_0})$$



$$F(\mathcal{G}) = -\frac{M_{\text{pl}}^2}{2} 2\Lambda(1 - e^{-\mathcal{G}/\mathcal{G}_0})$$

Conclusion

Concluding remarks

- Alternative theory of the gravity is necessary.
- There are many possibilities.

Modification of the Einstein eq.

Einstein equation

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Action

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Concluding remarks

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Concluding remarks

- Alternative theory of the gravity is necessary.
- There are many possibilities.
- To restrict the possibilities
 - Phenomenological consequences
 - Theoretical constraints