

# Relativistic Hydrodynamic Model in High-Energy Heavy-Ion Collisions

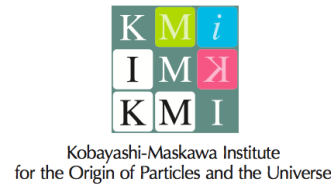
Graduate School of Advanced Science and Engineering, Hiroshima University  
International Institute for Sustainability with Knotted Chiral Meta Matter / SKCM<sup>2</sup>.

Hiroshima University

Kobayashi Maskawa Institute, Nagoya University

Department of Physics, Nagoya University

*Chiho NONAKA*

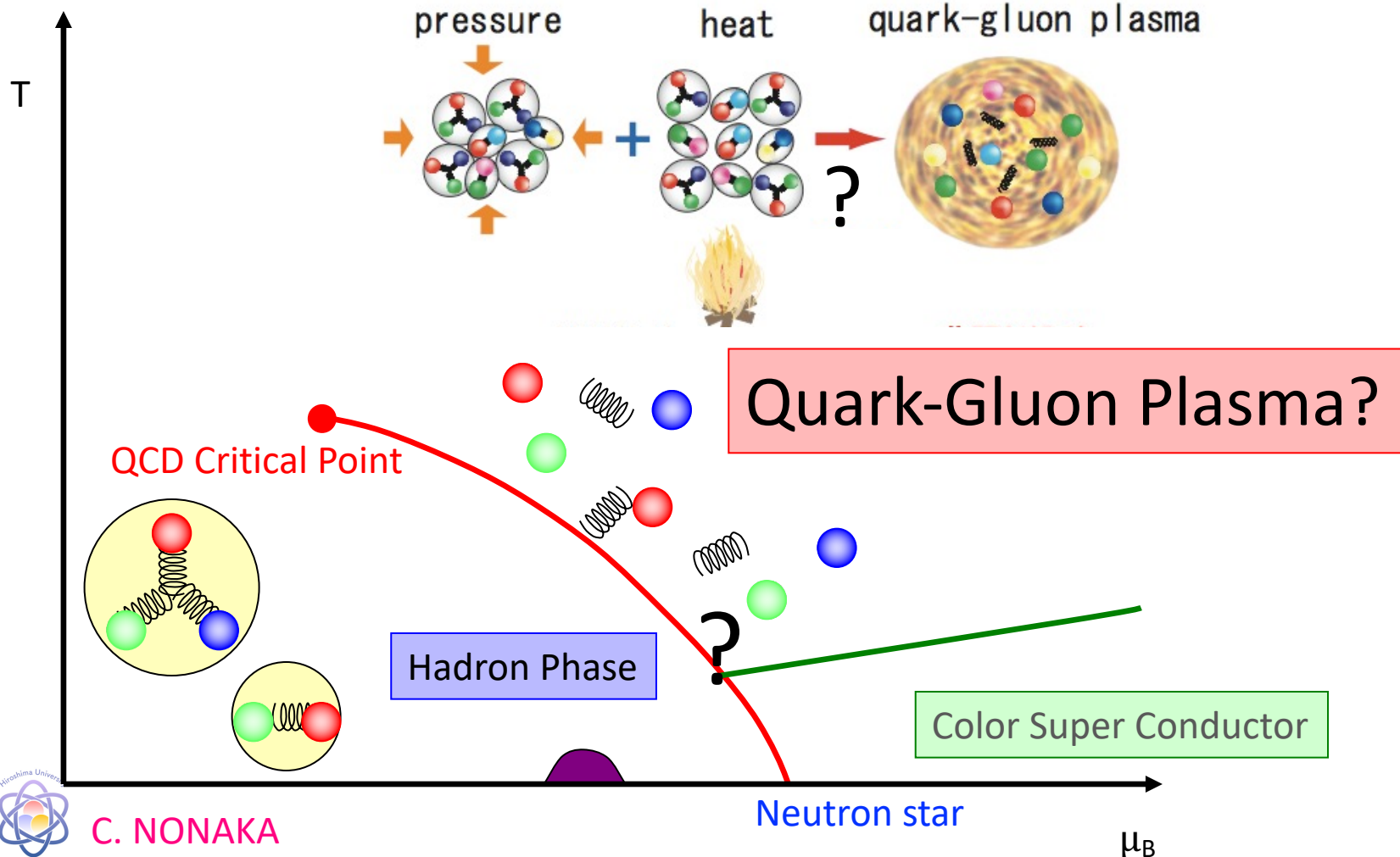


February 21, 2023@3rd IITB-Hiroshima workshop in HEP

# What is the QGP?

Quark-Gluon Plasma

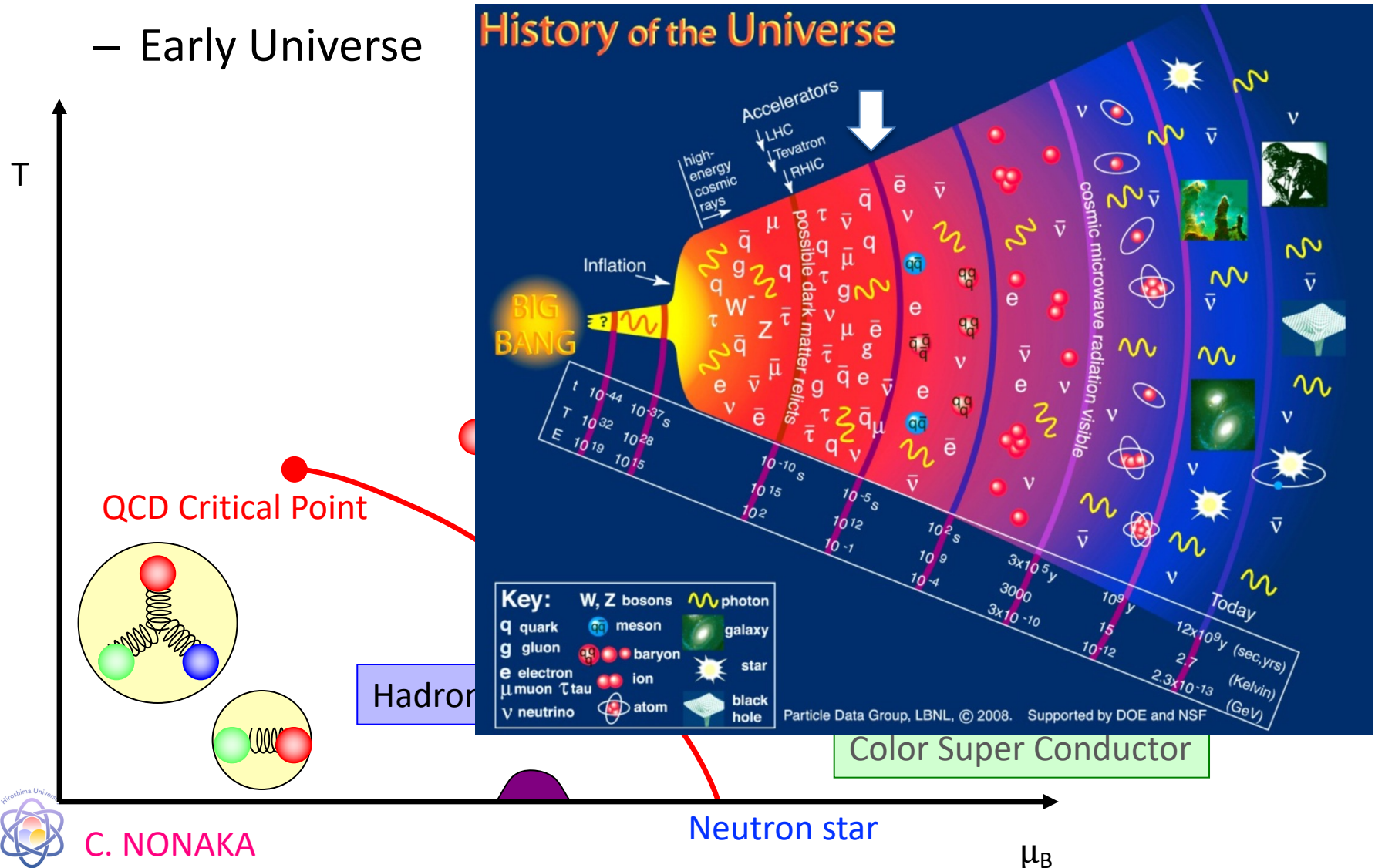
- Quarks and gluons at extreme conditions
  - High temperature and/or high density



# What is the QGP?

Quark-Gluon Plasma

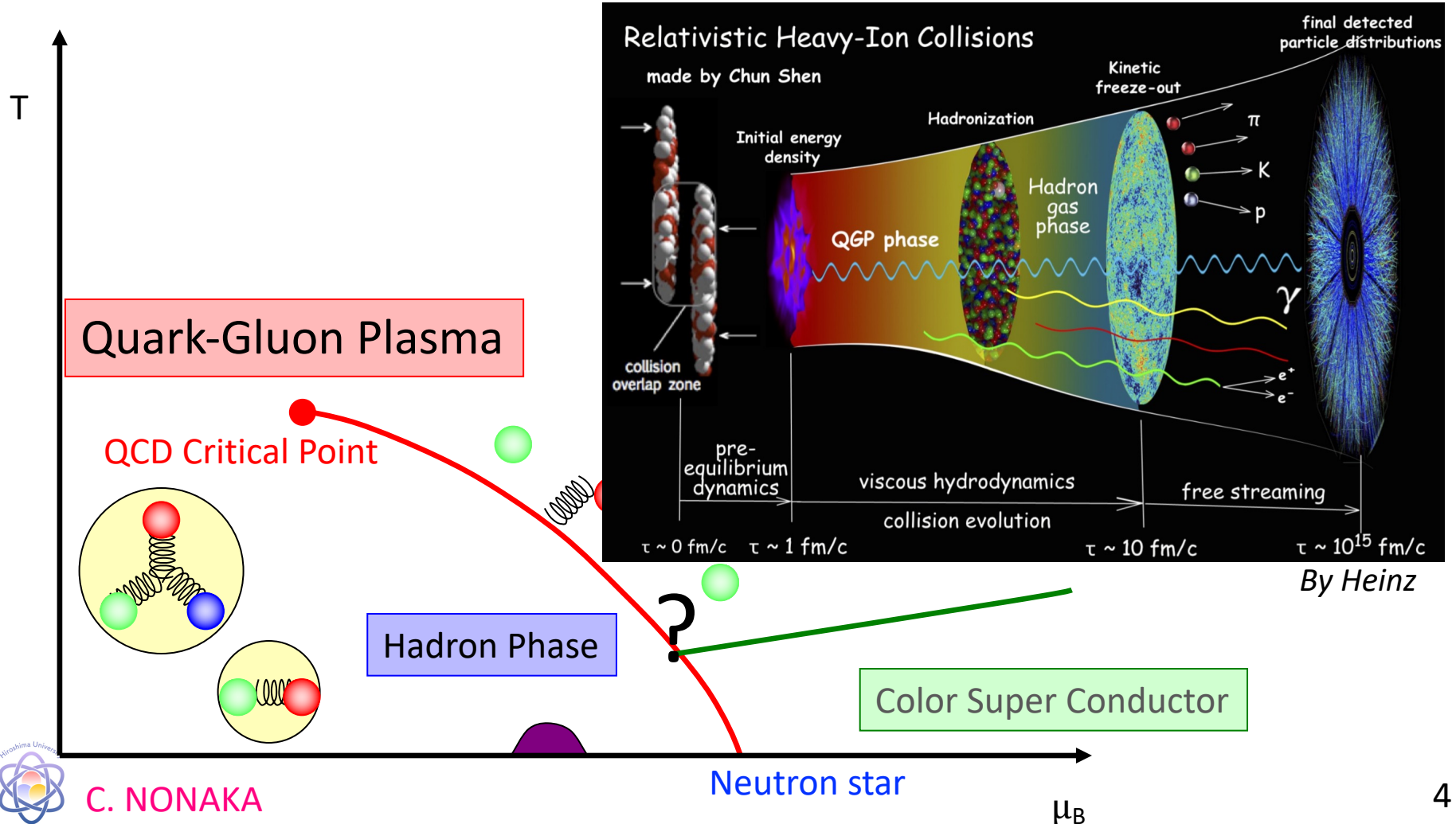
- Quarks and gluons at extreme conditions
  - Early Universe



# What is the QGP?

Quark-Gluon Plasma

- Quarks and gluons at extreme conditions
  - Relativistic Heavy Ion Collisions : Little Bang

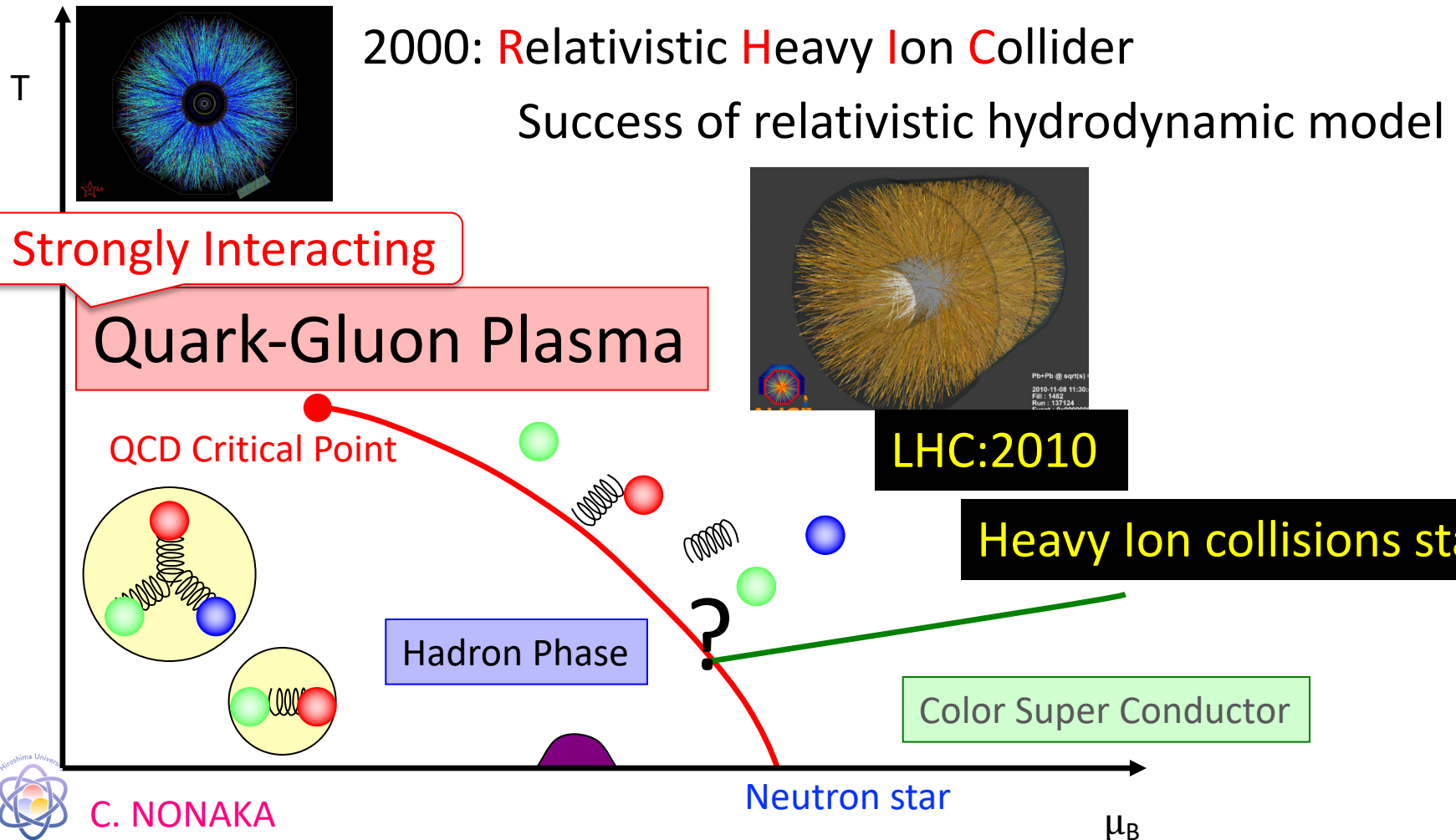




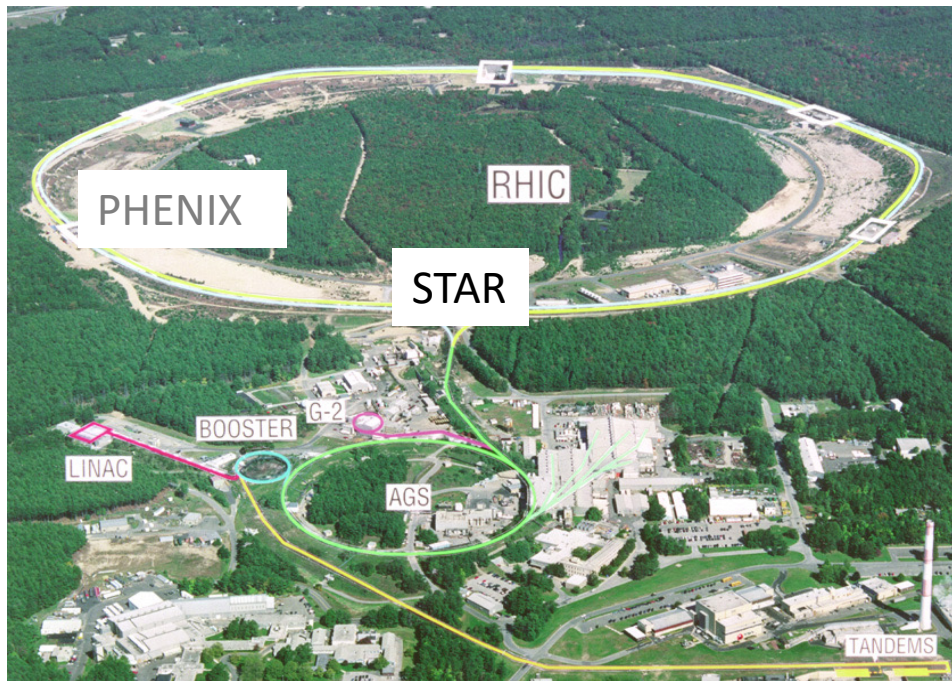
# What is the sQGP?

Quark-Gluon Plasma

- Quarks and gluons at extreme conditions
  - Relativistic Heavy Ion Collisions

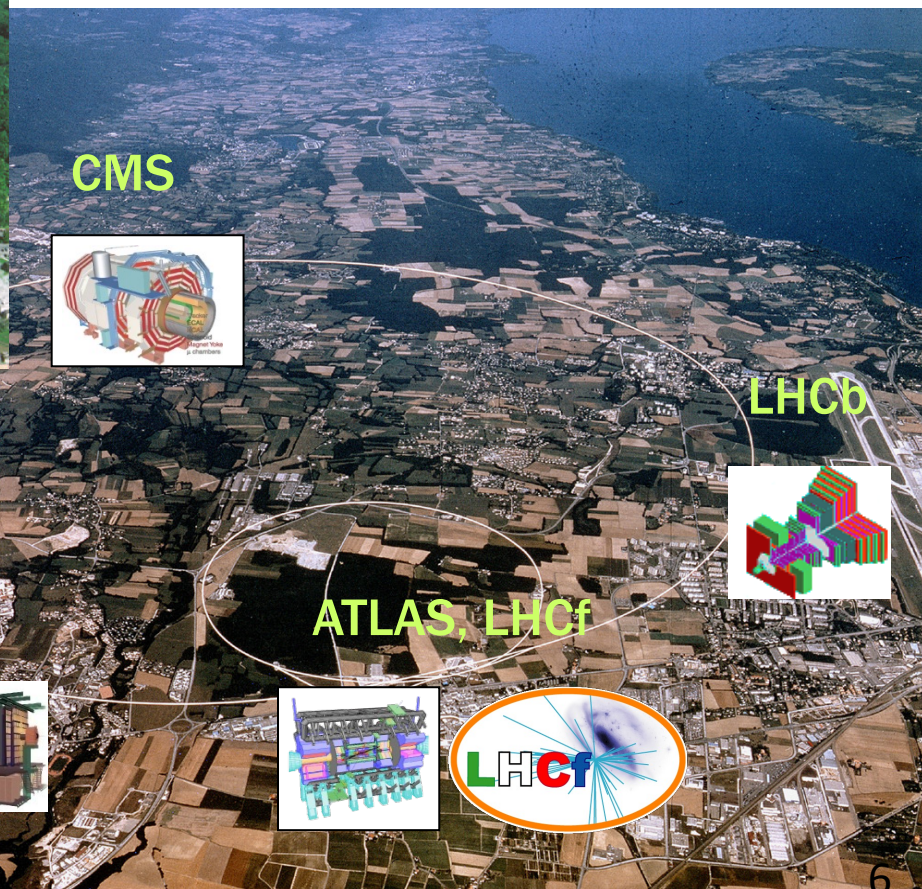


# Heavy Ion Collisions



RHIC@BNL

Large Hadron Collider@CERN

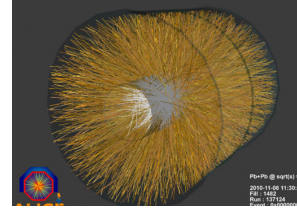
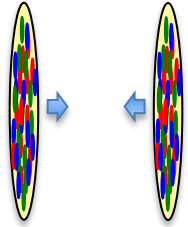




# Hydrodynamic Model

Experimental data

collisions



Hadrons  
Leptons  
Photons

Multiple particle production

➔ Hydrodynamic picture

Landau  
Bjorken

Success of hydrodynamic model at RHIC  
Relativistic viscous hydrodynamic model  
One of important phenomenological models

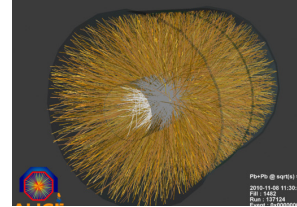
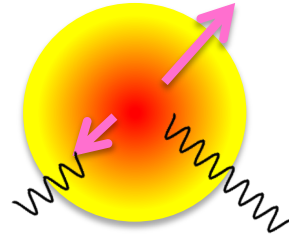
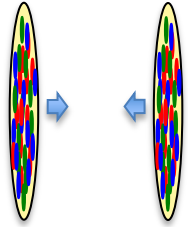


# Equation of State

Experimental data

collisions

hydrodynamics



Hydrodynamics

QCD phase diagram

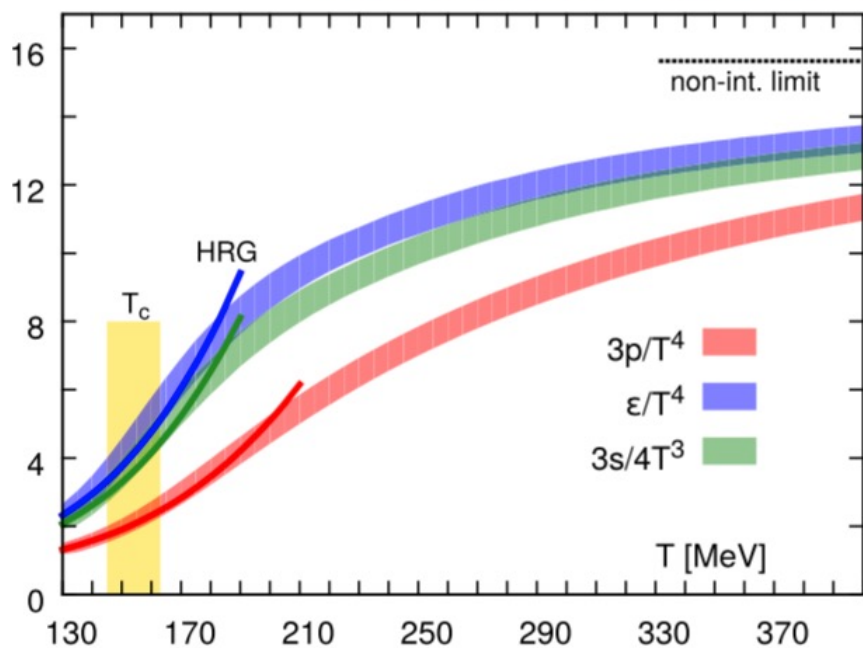
EoS: lattice QCD

# Equation of State

- Equation of State

– Lattice QCD

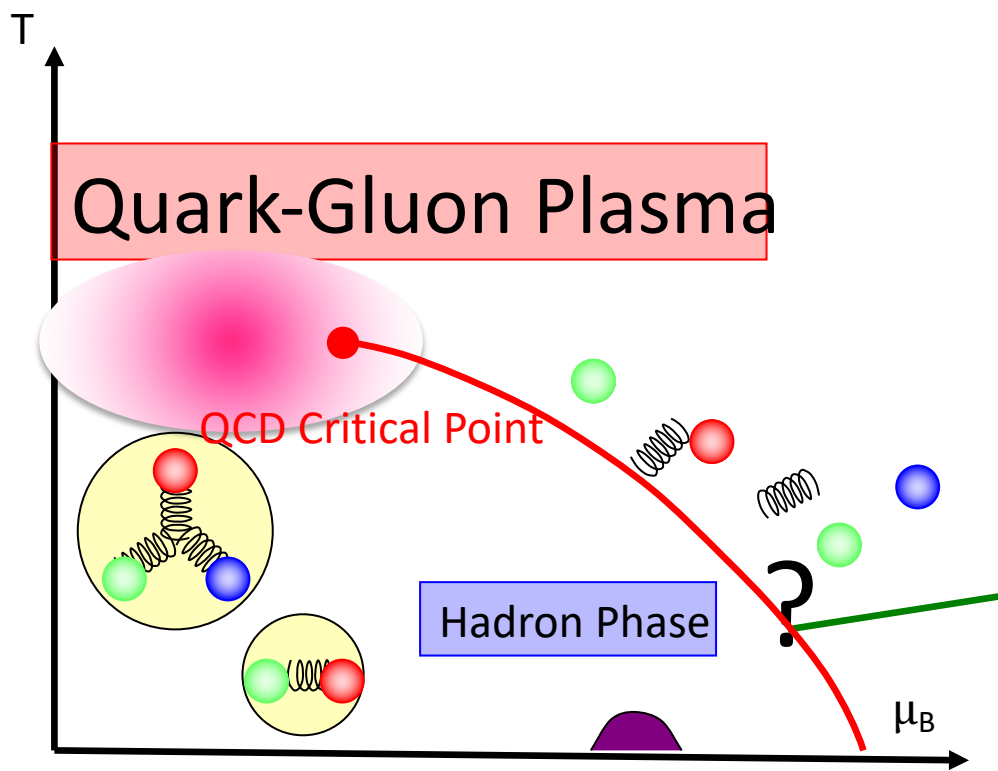
*HotQCD, PRD90, 094503 (2014)*



(2+1) flavor, Highly improved staggered quark action

$N_t=6,8,10,12, N_s=4N_t \rightarrow$  continuum limit

Parametrization of EoS



Neutron star



C. NONAKA

$T_c \sim 155$  MeV

finite  $\mu$ : sign problem

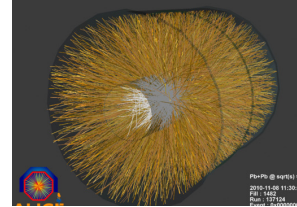
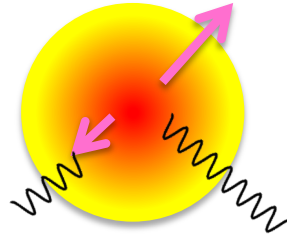
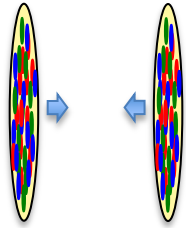


# QGP Bulk Properties

Experimental data

collisions

hydrodynamics



Hydrodynamics

QCD phase diagram

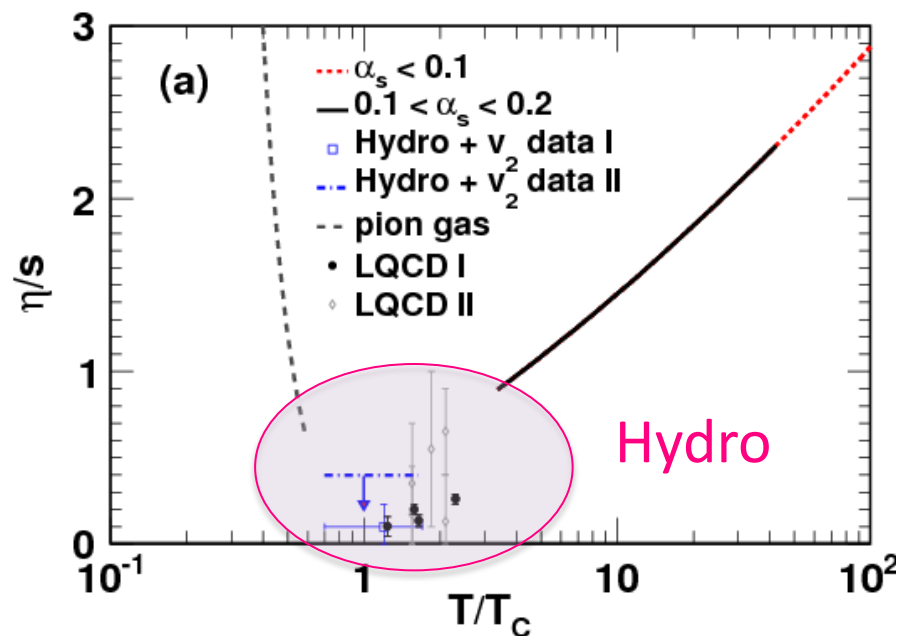
EoS: lattice QCD

Shear and bulk  
viscosities

# Property of QGP

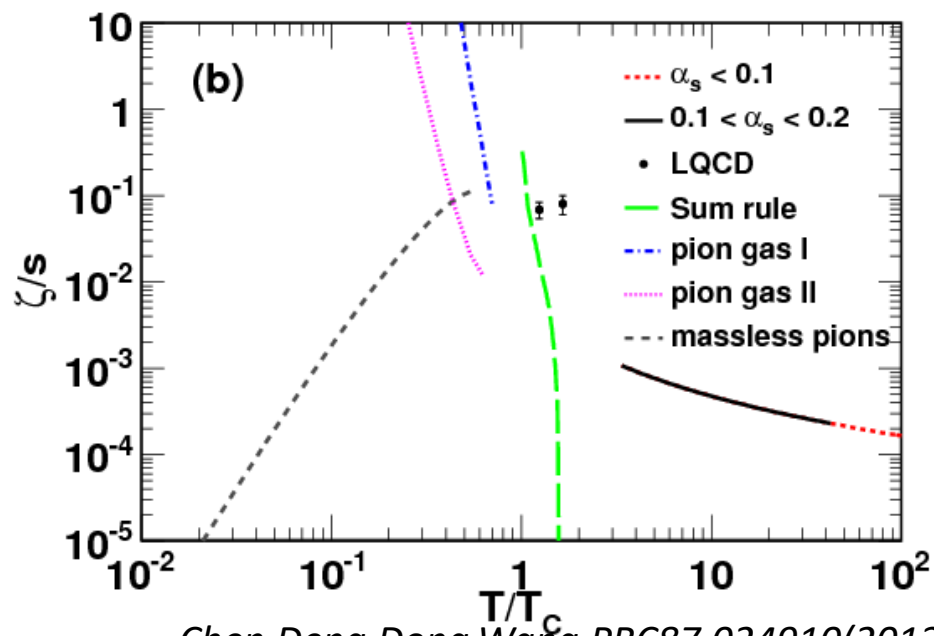
- Current Status for transport coefficients

shear viscosity



- Shear viscosity takes the minimum around  $T_c$ . Cf.  $\eta/s = 1/4\pi$  AdS/CFT
- Hydrodynamic model constant  $\eta/s$

bulk viscosity



*Chen, Deng, Dong, Wang, PRC87,024910(2013)*

- Bulk viscosity
- Temperature dependence is unclear.
- Hydrodynamic model vanishing

Detailed feature of shear and bulk viscosities

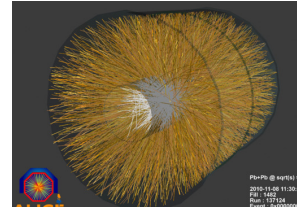
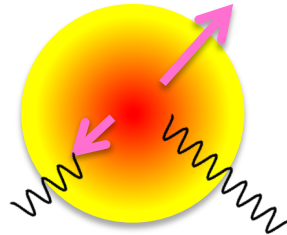
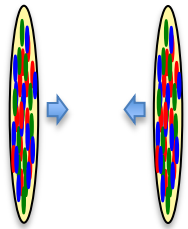


# QGP Bulk Properties

Experimental data

collisions

hydrodynamics



Hydrodynamics

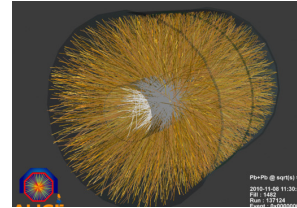
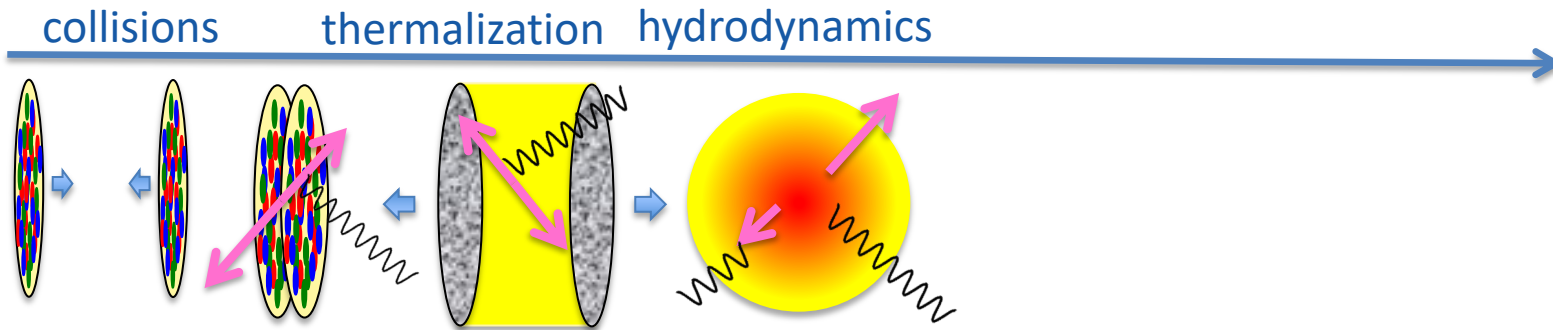
QCD phase diagram

EoS: lattice QCD

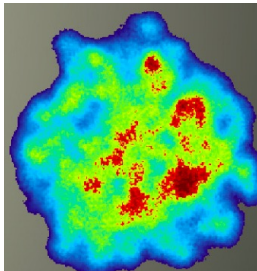
Shear and bulk viscosities

# Initial Condition

Experimental data



Initial conditions



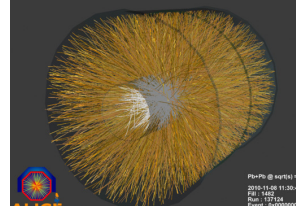
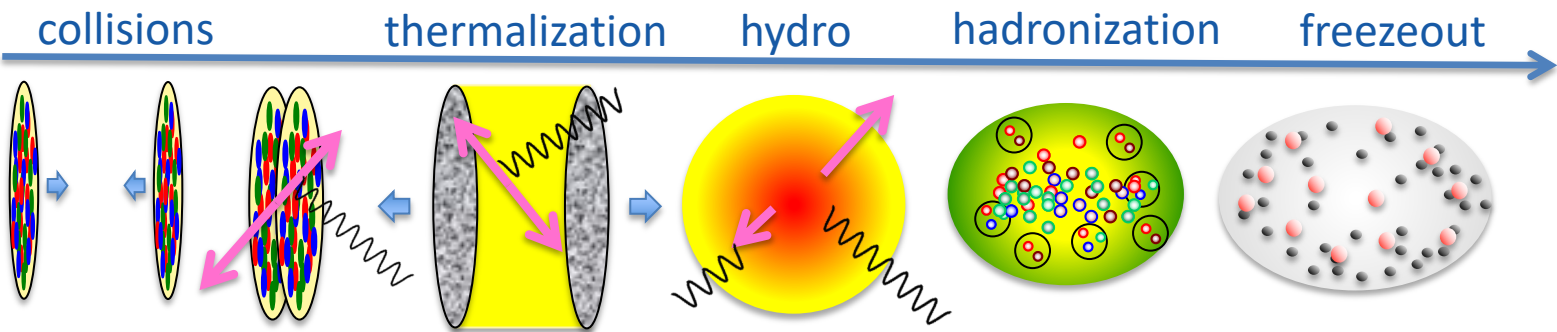
Energy (entropy) density distributions

Hydrodynamics

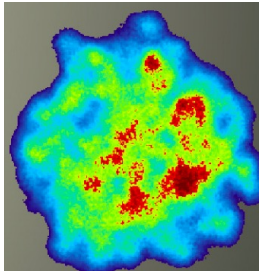
QCD phase diagram  
EoS: lattice QCD  
Shear and bulk viscosities

# From Fluid to Particle

Experimental data



Initial conditions



Energy (entropy) density distributions

Hydrodynamics

QCD phase diagram  
 EoS: lattice QCD  
 Shear and bulk viscosities

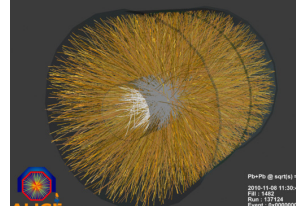
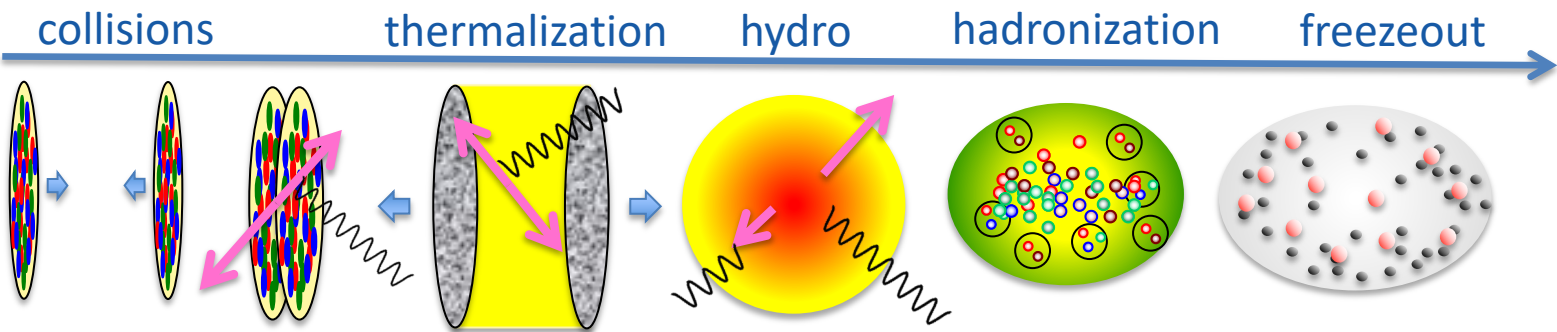
Final state interactions

Hadron based event generator

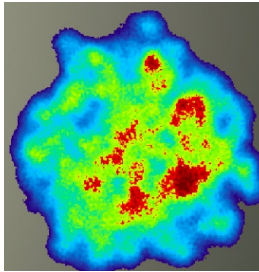


# Quantitative Analyses

Experimental data



Initial conditions



Energy (entropy) density distributions

Hydrodynamics

QCD phase diagram  
EoS: lattice QCD  
Shear and bulk viscosities

New hydrodynamics code

Final state interactions

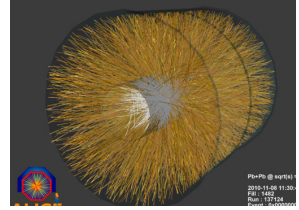
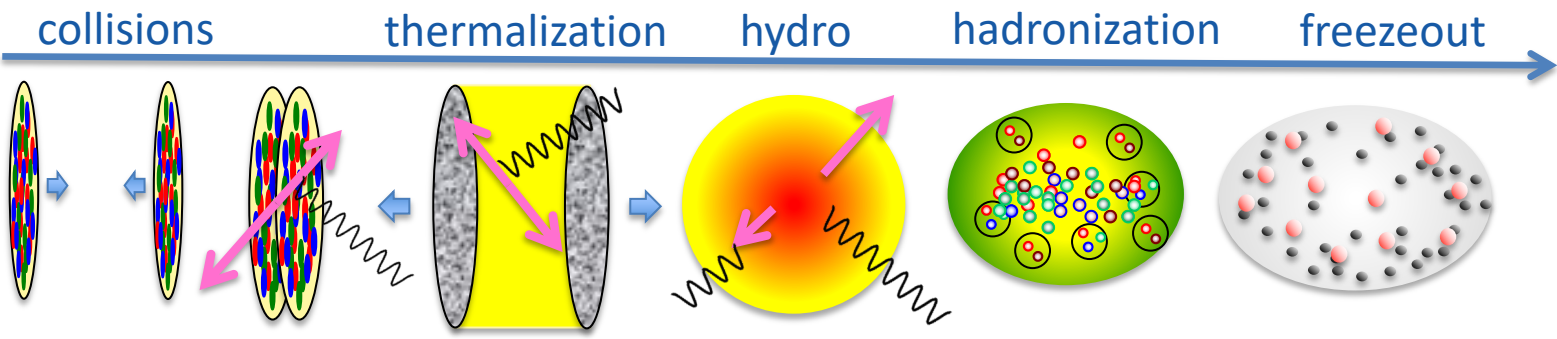
Hadron based event generator

$$\partial_\mu T^{\mu\nu} = 0$$

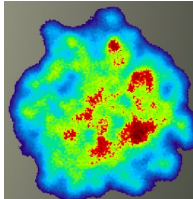
Akamatsu et al, JCP256,34(2014)  
Okamoto, Akamatsu, Nonaka, EPJC76,579(2016)  
Okamoto and Nonaka, EPJC77,383(2017)

# Quantitative Analyses

Experimental data



Initial conditions



TRENTO

Phenomenological model  
Parametrization

Moreland et al., PRC92,011901(2015)  
Ke et al., PRC96,044192(2017)

Hydrodynamics

QGP bulk property  
EoS: lattice QCD  
Shear and bulk viscosities

New hydrodynamics code

$$\partial_\mu T^{\mu\nu} = 0$$

Final state interactions

Hadron based event generator

UrQMD

Bass et al., Prog.Part.Nucl.Phys.(1998)  
Bleicher et al., J.Phys.G25,1859(1999)

**Application to analyses of RHIC and LHC data**

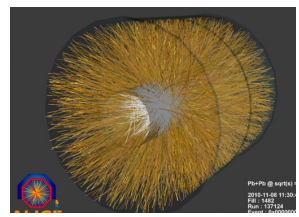
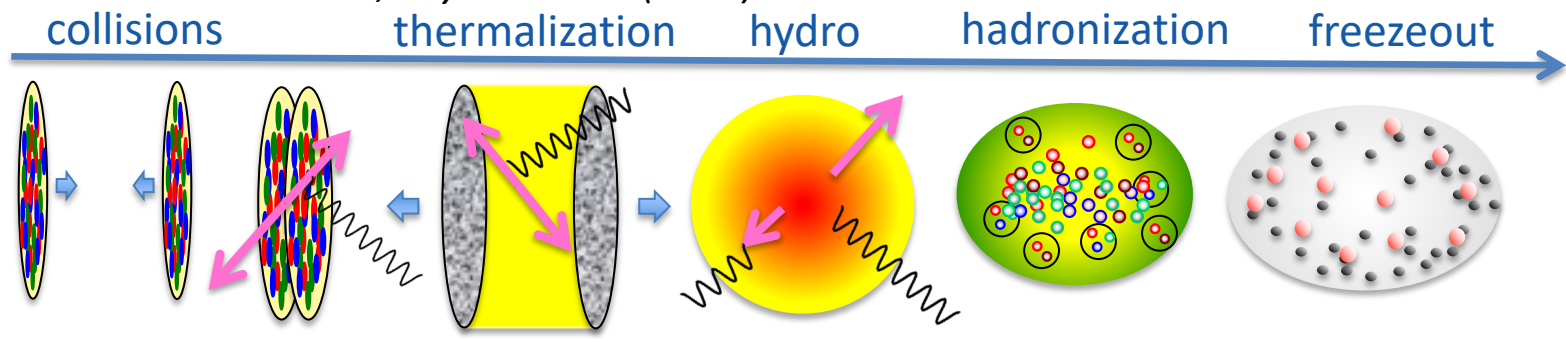


C. NONAKA

# Bulk Property of QGP

Experimental data

Okamoto and Nonaka, Phys.Rev. C98 (2018) 054906

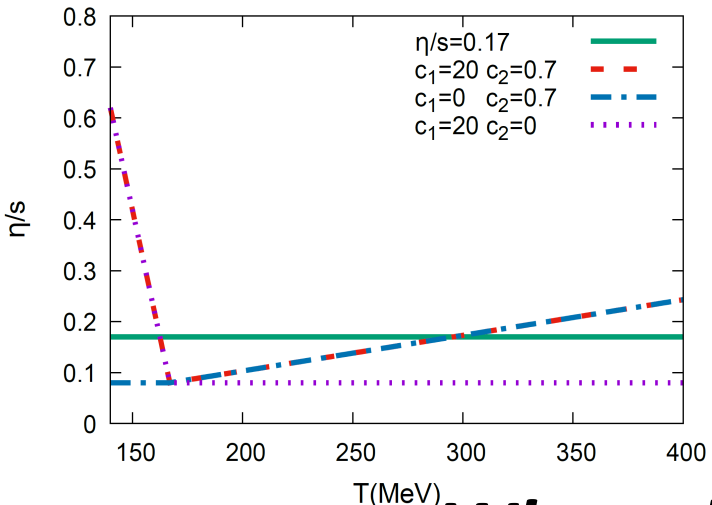


**Our Model**



**Experimental data**

Shear viscosity



ALICE Pb+Pb  $\sqrt{s_{NN}} = 2.76$  TeV, LHC

- ✓ Rapidity distributions
- ✓  $P_T$  distributions
- ✓ Mean  $P_T$
- ✓ Collective flows  $v_2$  and  $v_3$

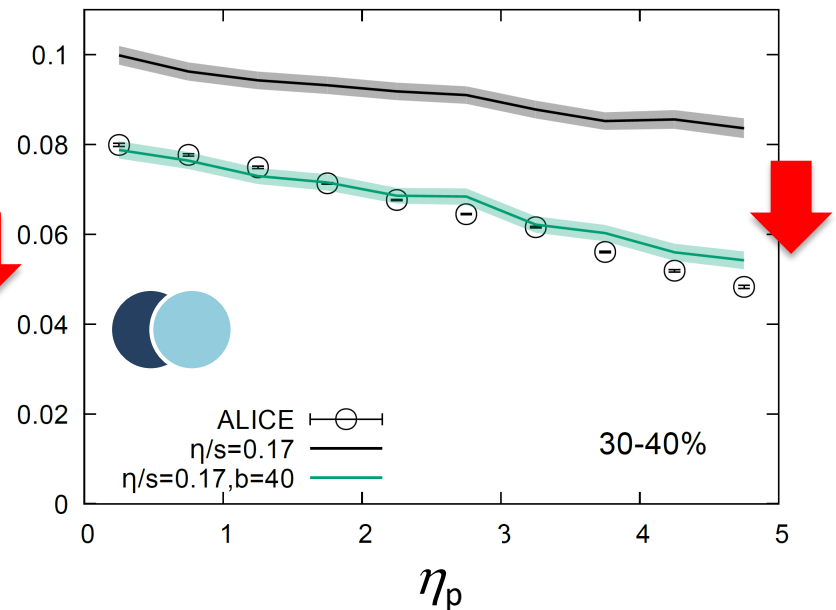
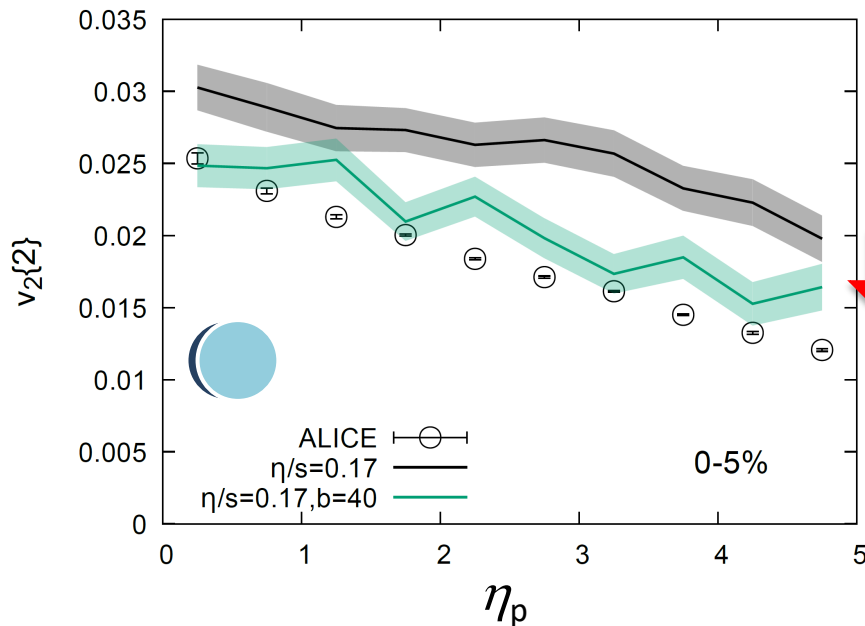
Bulk viscosity

$$\zeta = b\eta \left( \frac{1}{3} - c_s^2 \right)^2 \quad b = 40$$

**What physical observable is interesting?**

# Effect on Collective Flow

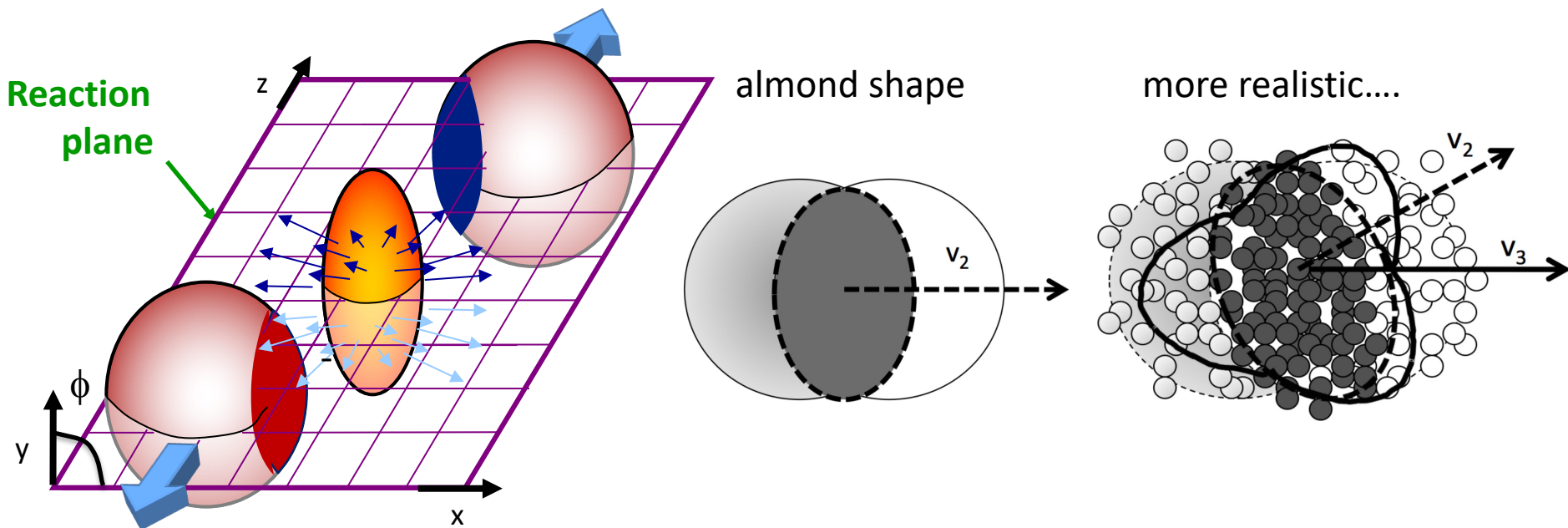
- Collective flow as a function of  $\eta_p$



- (3+1)-d calculation
- $v_n$  with bulk viscosity is much closer to the ALICE data: amplitude and slope
- Effect of bulk viscosity at forward rapidity is large.

**Finite bulk viscosity**

# Collective Flow



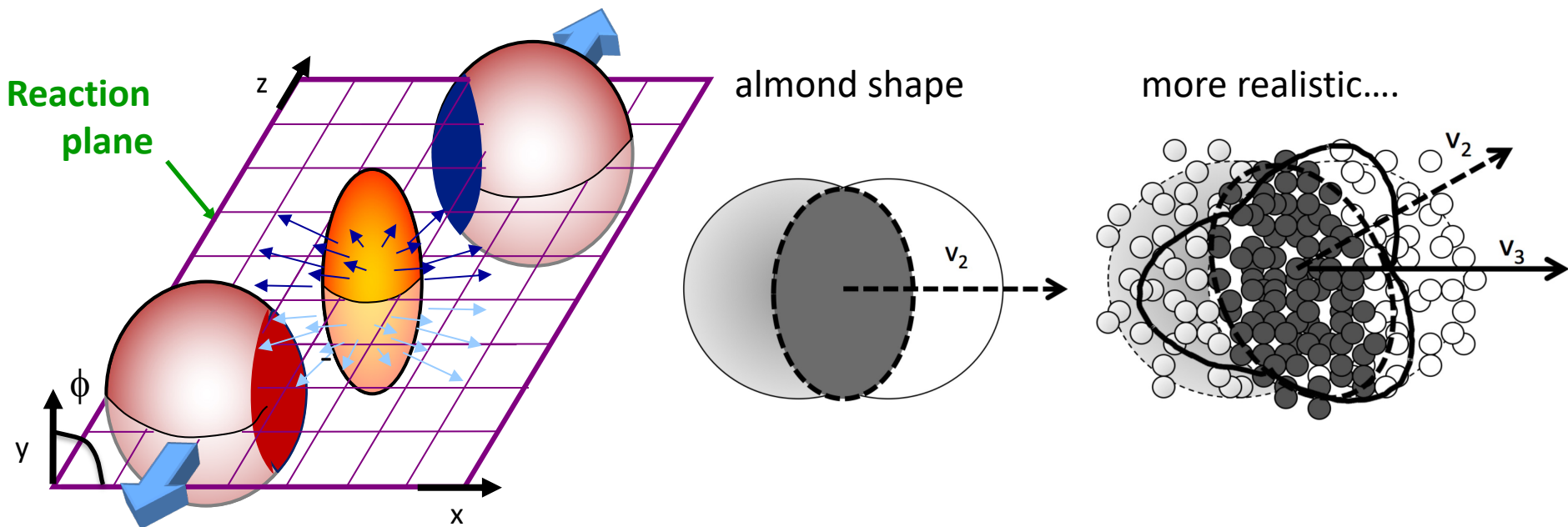
$$\frac{dN}{d\phi} \sim N_0 \left( 1 + \underline{2v_1} \cos \phi + \underline{2v_2} \cos 2\phi \right)$$

Directed flow

Elliptic flow



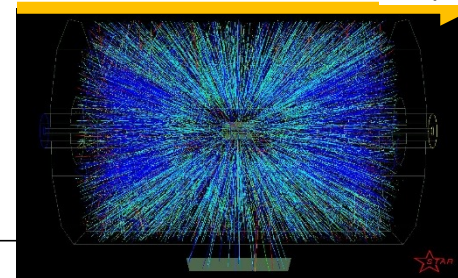
# Collective Flow



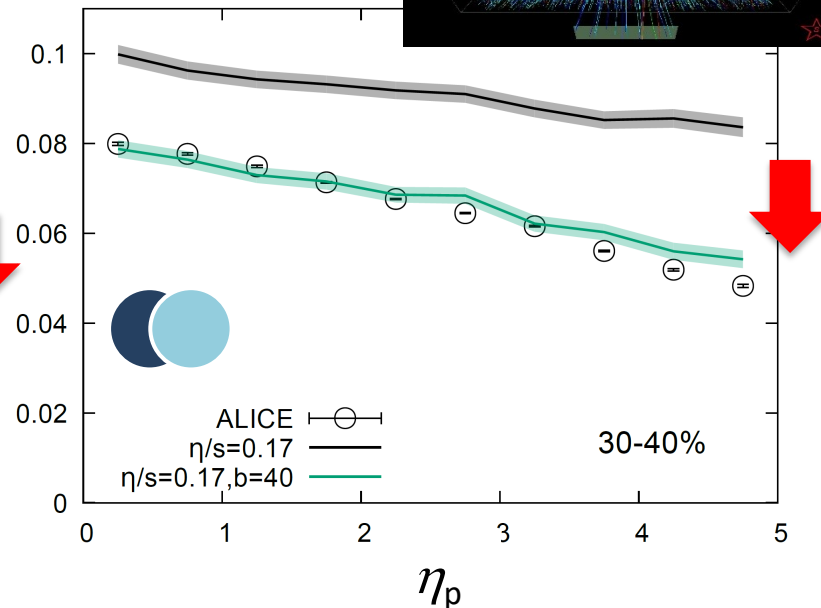
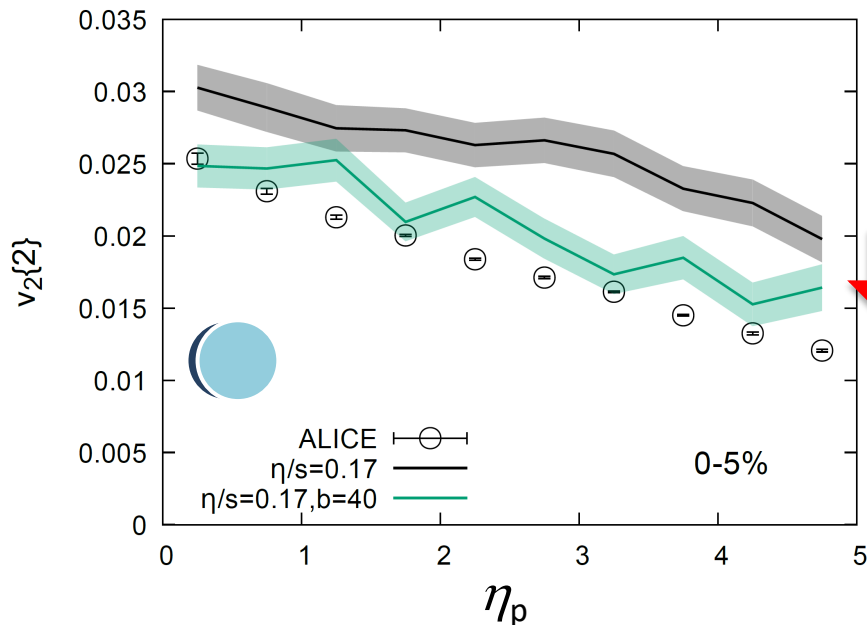
$$\frac{dN}{d\phi} \sim N_0 \left( 1 + \underbrace{2v_1 \cos \phi}_{\text{Directed flow}} + \underbrace{2v_2 \cos 2\phi}_{\text{Elliptic flow}} + 2v_3 \cos 3\phi + 2v_4 \cos 4\phi + \dots \right)$$

# Effect on Collective Flow

pseudorapidity  $\eta_p$



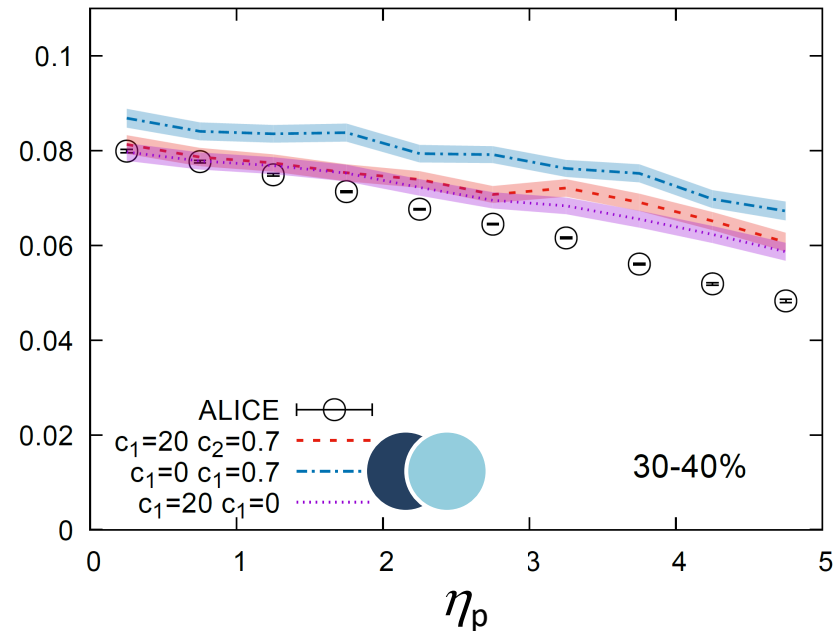
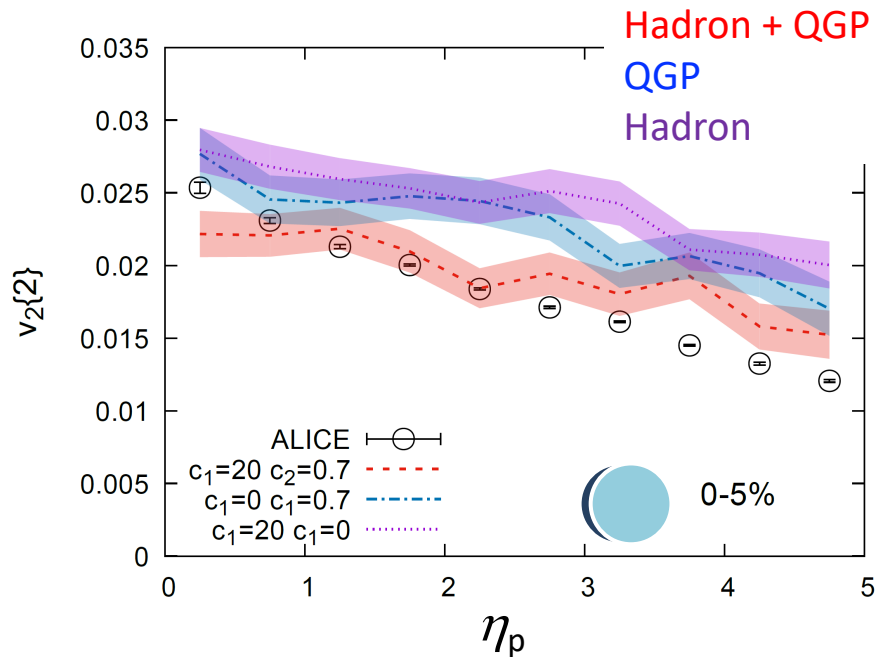
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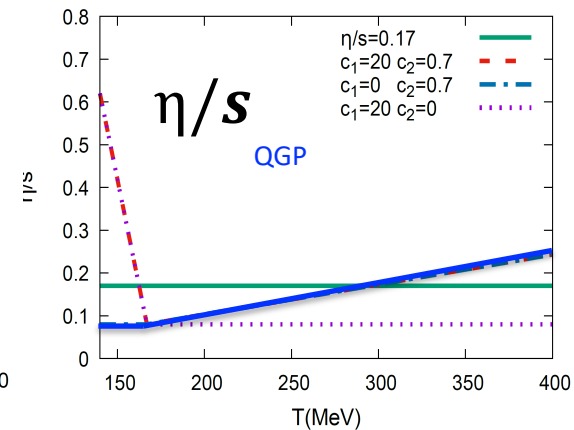
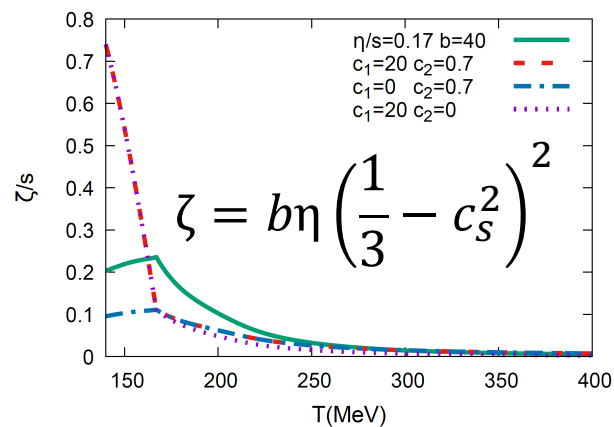
# Temperature Dependent $\eta/s$



QGP phase:  $\eta/s(T)$

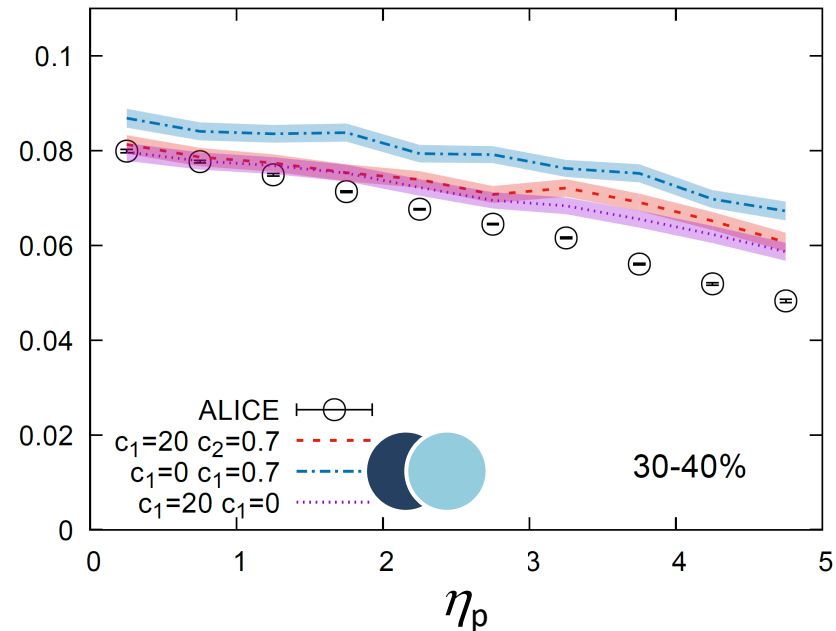
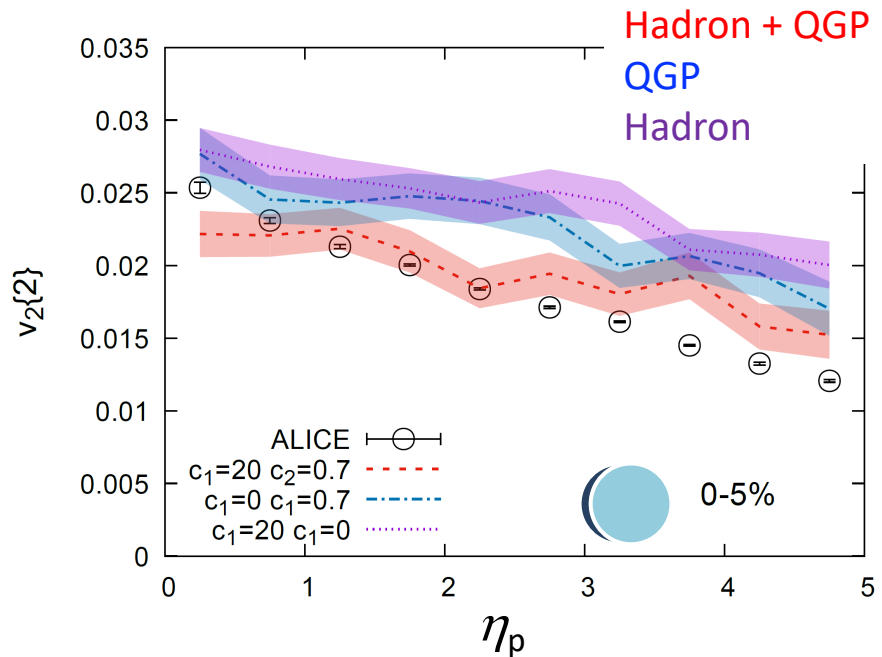
Hadron phase:  $\eta/s=0.08$

In both centrality classes  
 $v_2(\eta_p)$  is larger.



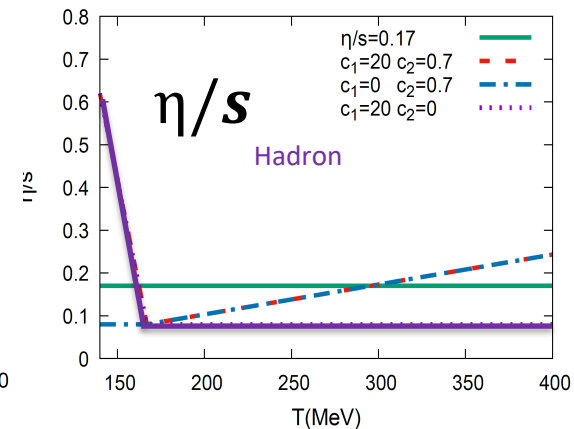
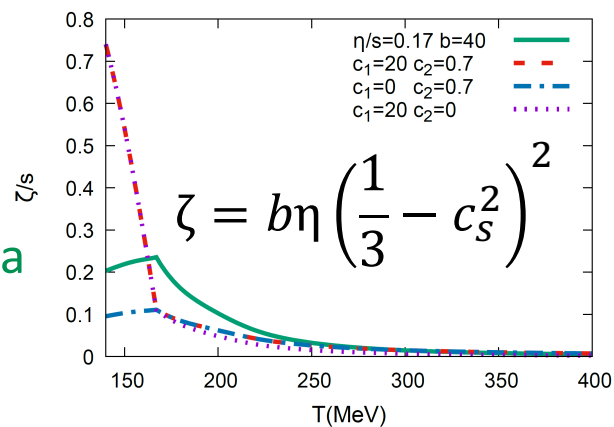
$T_{SW} = 150$  MeV

# Temperature Dependent $\eta/s$



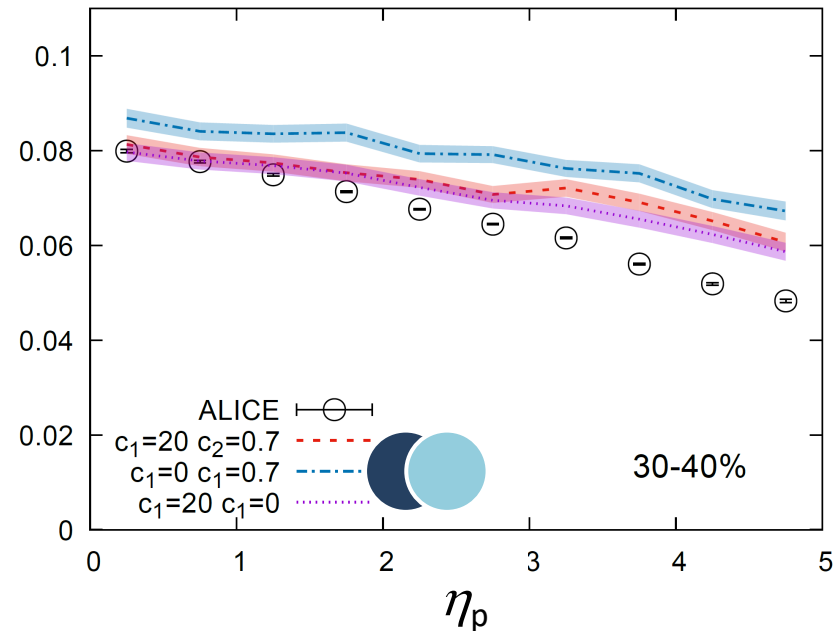
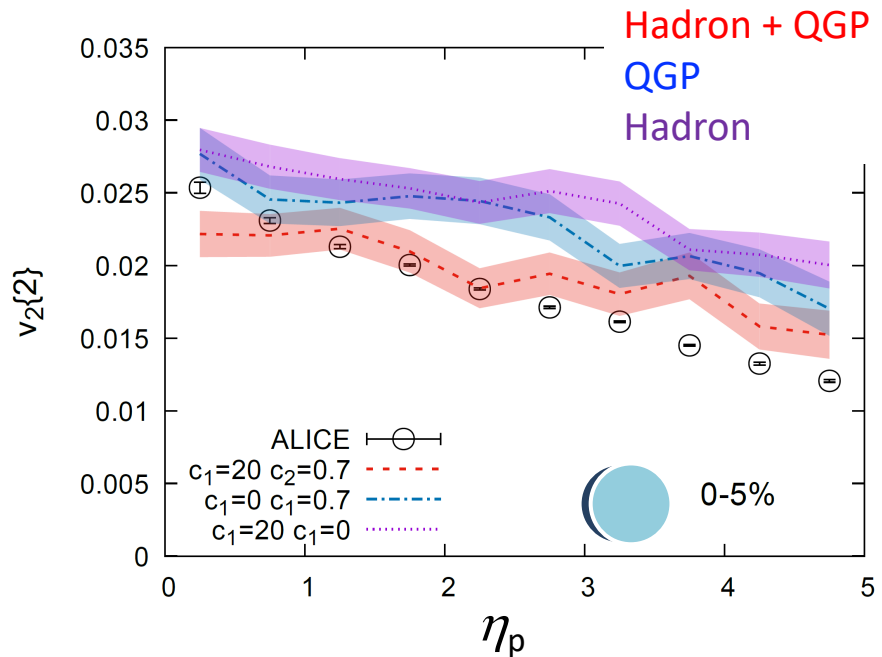
QGP phase:  $\eta/s=0.08$   
Hadron phase:  $\eta/s(T)$

0-5%: larger than ALICE data  
30-40%: close to ALICE data



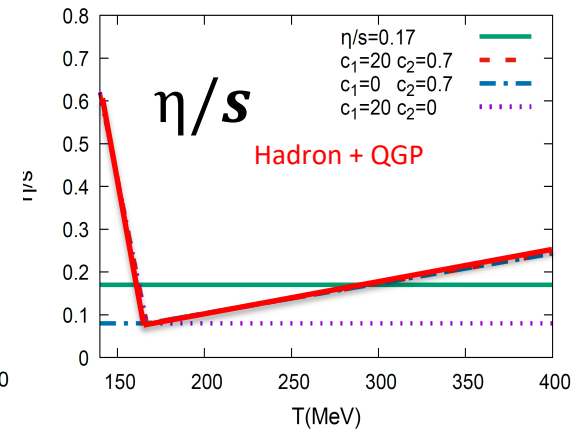
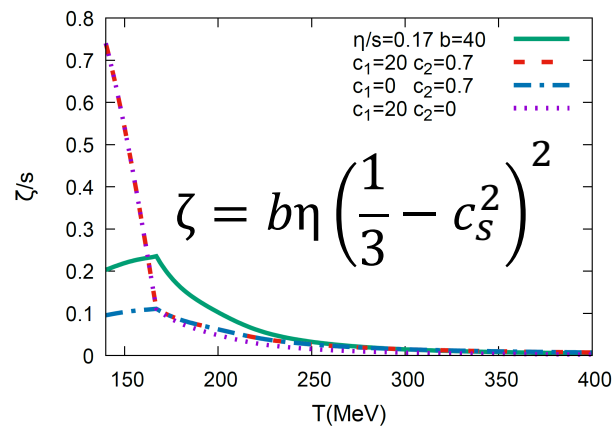
$T_{SW} = 150$  MeV

# Temperature Dependent $\eta/s$



QGP phase:  $\eta/s(T)$   
Hadron phase:  $\eta/s(T)$

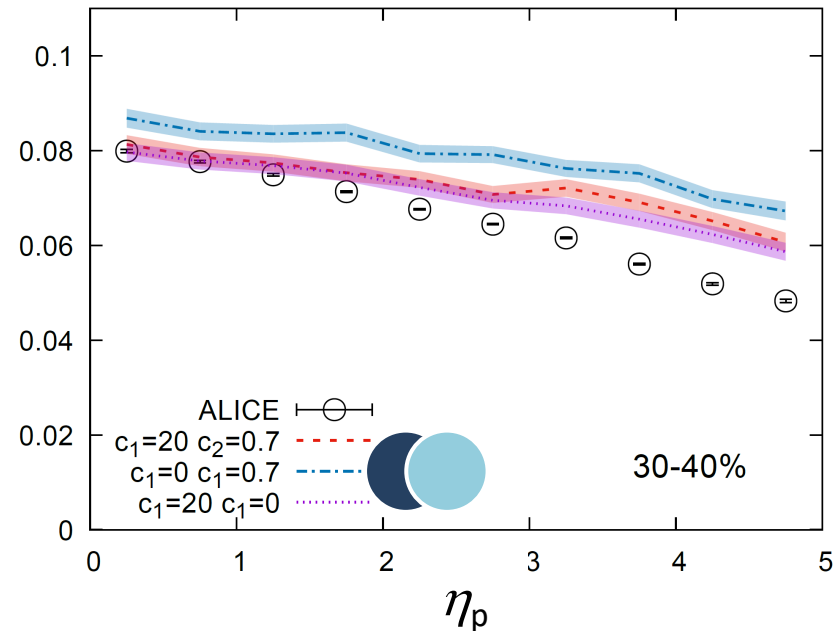
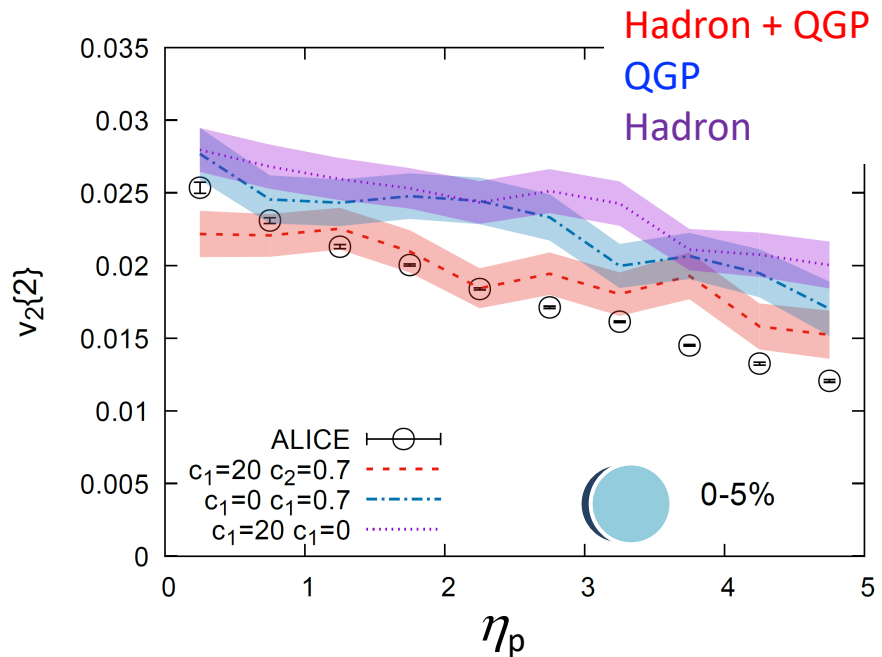
In both central classes,  
close to ALICE data



$$T_{\text{SW}} = 150 \text{ MeV}$$

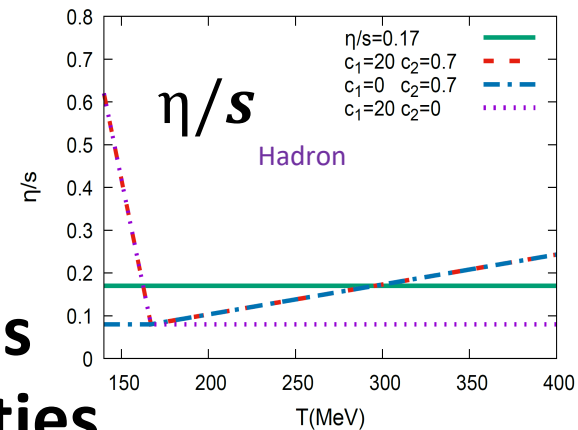


# Temperature Dependent $\eta/s$



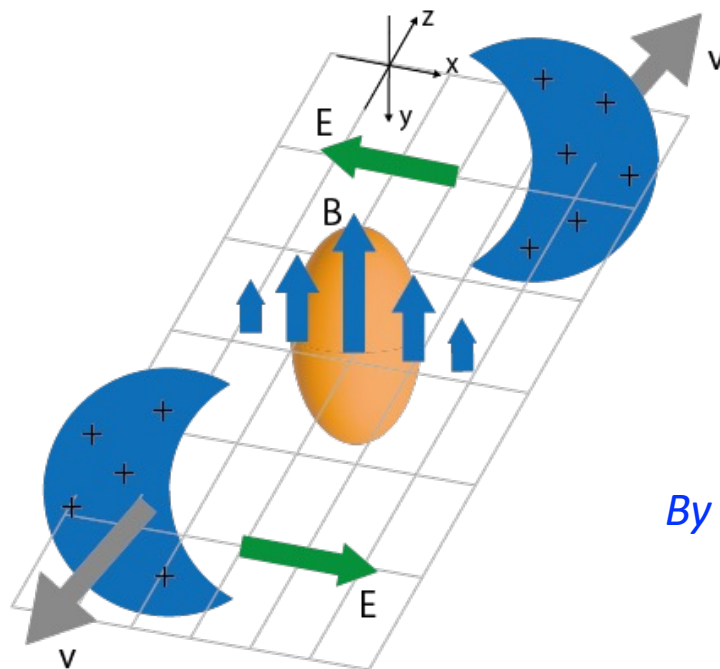
- 0-5 % centrality  
 $\eta/s$  of QGP and hadron phases is important.
- 30-40 % centrality  
 $\eta/s$  of hadron phase is important.

**Central dependence of  $v_2(\eta_p)$  reveals temperature dependence of viscosities.**



# Electromagnetic Field in HIC

- Strong Electromagnetic field ?
  - Au + Au ( $\sqrt{s_{NN}} = 200$  GeV) :  $10^{14}$  T  $\sim 10 m_{\pi}^2$
  - Pb + Pb ( $\sqrt{s_{NN}} = 2.76$  TeV) :  $10^{15}$  T



By Nakamura

*Nakamura, Miyoshi, C. N. and Takahashi, Phys. Rev. C 107, no.1, 014901 (2023)*

*Nakamura, Miyoshi, C. N. and Takahashi, arXiv:2211.02310 [nucl-th]*

*Nakamura, Miyoshi, C. N. and Takahashi, arXiv:2212.02124 [nucl-th]*



# Relativistic Resistive Magneto-Hydrodynamics

Nakamura, Miyoshi, C. N. and Takahashi, *PRC107*, no. 1, 014901 (2023)

collisions

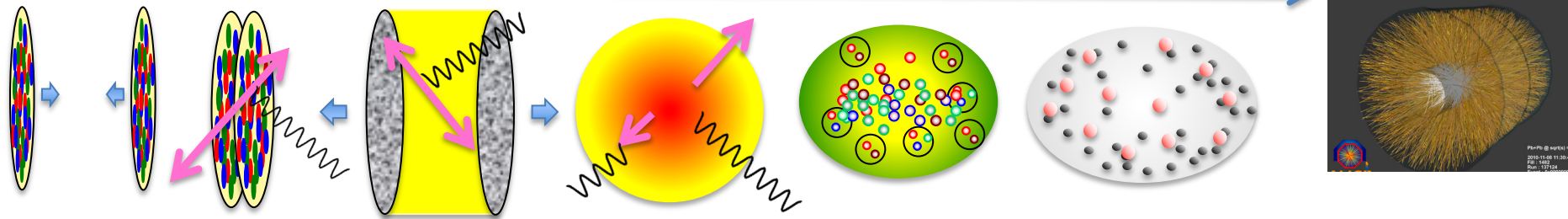
thermalization

hydro

hadronization

freezeout

Experimental data



Initial conditions

Hydrodynamics

Glauber model  
+ approximate solutions of Maxwell eq.

Hydrodynamic eq. + Maxwell eq. + Ohm's law

$$\partial_\mu T^{\mu\nu} = F^{\nu\lambda} J_\lambda \quad J^\mu = \sigma e^\mu$$

# Relativistic Resistive Magneto-Hydrodynamics (RRMHD)

Nakamura, Miyoshi, C.N. and Takahashi, arXiv:2211.02310 [nucl-th]

## ■ RRMHD equation

➤ Conservation law + Maxwell eq. + Ohm's law

$$\partial_\mu T^{\mu\nu} = F^{\nu\lambda} J_\lambda$$

$$J^\mu = J_c^\mu + qu^\mu$$

$e$ : energy density  
 $p$ : pressure  
 $p_{em} = (E^2 + B^2)/2$

$\varepsilon = (e + p)\gamma^2 - p + p_{em}$   
 $m^i = (e + p)\gamma^2 v^i + \epsilon^{ijk} B_j E_k$   
 $\Pi^{ij} = (e + p)\gamma^2 v^i v^j + (p + p_{em})g^{ij} - E^i E^j - B^i B^j$

### Energy Conservation

$$\partial_t \varepsilon + \nabla \cdot m = 0$$

### Momentum conservation

$$\partial_t m^i + \nabla \cdot \Pi^i = 0$$

### Faraday's law

$$\partial_t \vec{B} + \nabla \times \vec{E} = 0$$

### Ohm's law

$$\vec{J} = q\vec{v} + \sigma\gamma[\vec{E} + \vec{v} \times \vec{B} - (\vec{v} \cdot \vec{E})\vec{v}]$$

### Ampere's law

$$\partial_t \vec{E} - \nabla \times \vec{B} = -\vec{J}$$

$$=: \vec{J}_c \quad \rightarrow \quad \partial_t \vec{E} - \nabla \times \vec{B} = q\vec{v}$$

$$\rightarrow \quad \partial_t \vec{E} = \vec{J}_c \quad \text{演算子分離}$$

- Integration with quasi-analytic solutions

$$\vec{E}_\perp = -\vec{v} \times \vec{B} + (E_\perp^0 + \vec{v} \times \vec{B}) \exp(-\sigma\gamma t)$$

$$\vec{E}_\parallel = E_\parallel^0 \exp(-\sigma t/\gamma)$$

Komissarov, Mon. Not. R. Astron. Soc. 382, 995-1004 (2007)



# RRMHD eq. in the Milne Coordinates

*Nakamura, Miyoshi, C. N, and Takahashi, arXiv:2211.02310 [nucl-th]*

- Milne Coordinates

- Expansion system in the longitudinal coordinates  $(\tau, x, y, \eta_s)$

- Strong velocity in the longitudinal direction  $\rightarrow 0$
- Saves the number of cells in the collision axis direction

$$\tau = \sqrt{t^2 - z^2}$$

$$\eta_s = \frac{1}{2} \ln \frac{t+z}{t-z}$$

RRMHD eq.

$$\partial_\tau(\tau U) + \partial_i(\tau F^i) = \tau S$$

$$U = \begin{pmatrix} D \\ m_j \\ \varepsilon \\ B^j \\ E^j \\ q \end{pmatrix}, F^i = \begin{pmatrix} Dv^i \\ \Pi^{ji} \\ m^i \\ \varepsilon^{jik} E_k \\ \varepsilon^{jik} B_k \\ J^i \end{pmatrix}, S = \begin{pmatrix} 0 \\ \frac{1}{2} T^{ik} \partial_j g_{ik} \\ -\frac{1}{2} T^{ik} \partial_0 g_{ik} \\ 0 \\ J_c^i \\ 0 \end{pmatrix}$$

**Newly developed RRMHD code in Milne coordinates**



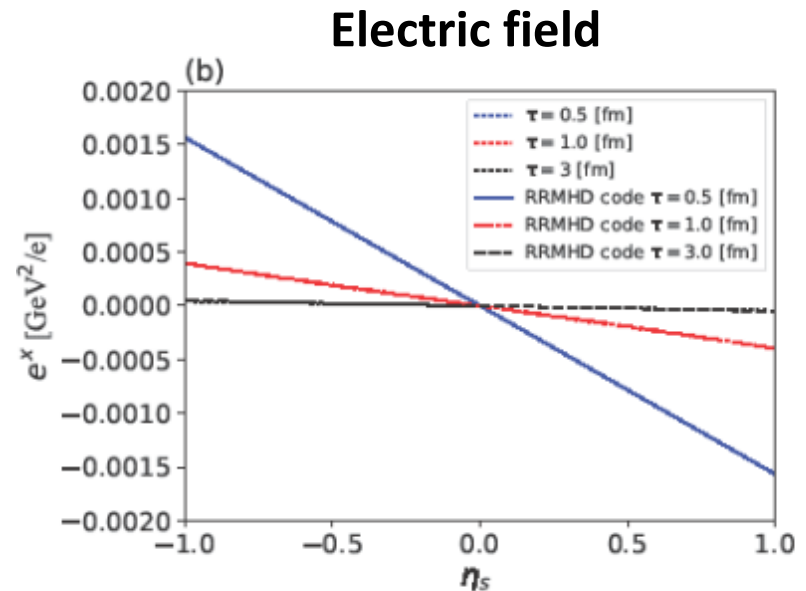
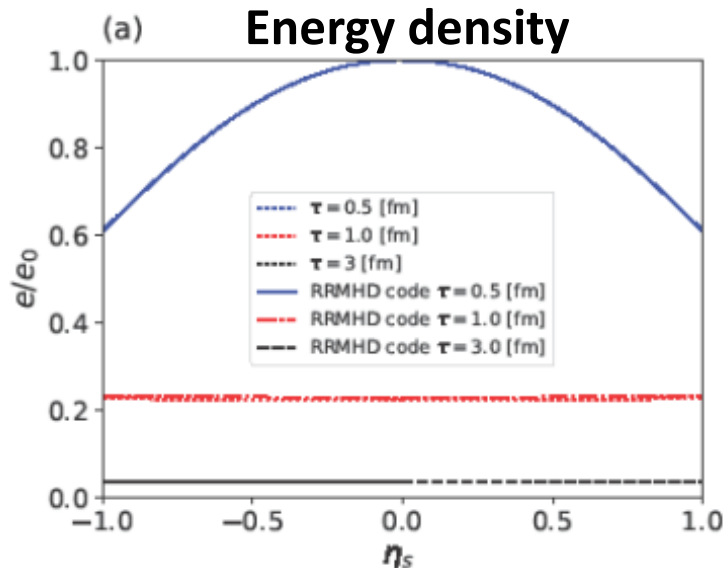
# Validation of the Code

Nakamura, Miyoshi, C. N. and Takahashi, arXiv:2211.02310 [nucl-th]

- RRMHD in the Milne coordinates

## New Test Problem

- (1+1) dimensional expansion system  $u^\mu = (\cosh Y, 0, 0, \sinh Y)$ 
  - Comparison between quasi-analytical solution and RRMHD simulation



Solid line : RRMHD code  
Dashed line: quasi-analytical solution

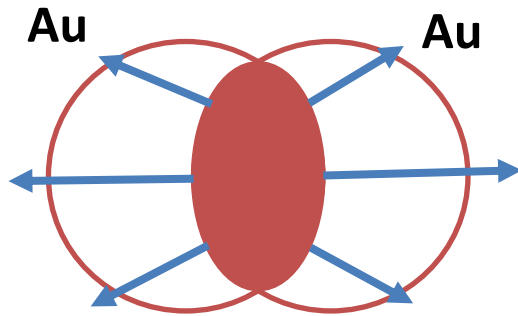
➔ Application to Heavy Ion Collisions



# Symmetric and Asymmetric Systems

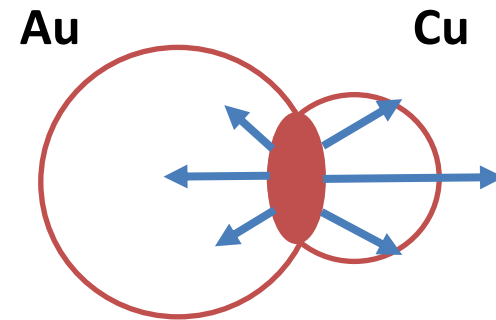
RHIC  $\sqrt{s_{NN}} = 200$  GeV

## ■ Au-Au collisions



- Symmetric pressure gradient
- Almond-shaped medium

## ■ Cu-Au collisions

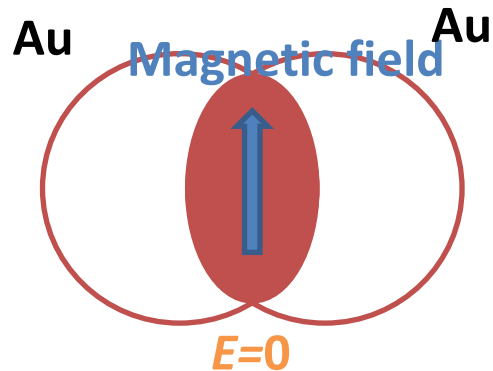


- Asymmetric pressure gradient
- Distorted Almond-shaped medium



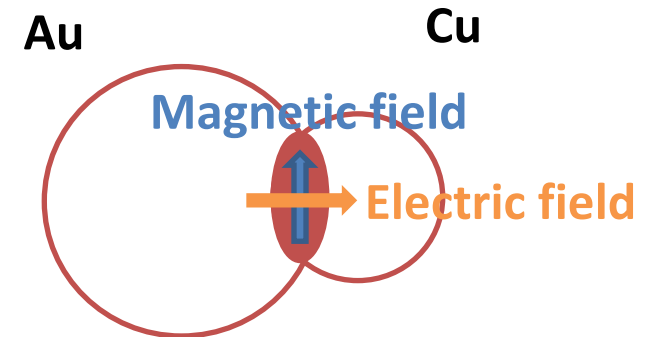
# Electromagnetic Field in Symmetric and Asymmetric Systems

## ■ Au-Au collisions



- Magnetic field
  - Strong magnetic field
- Electric field
  - No electric field

## ■ Cu-Au collisions



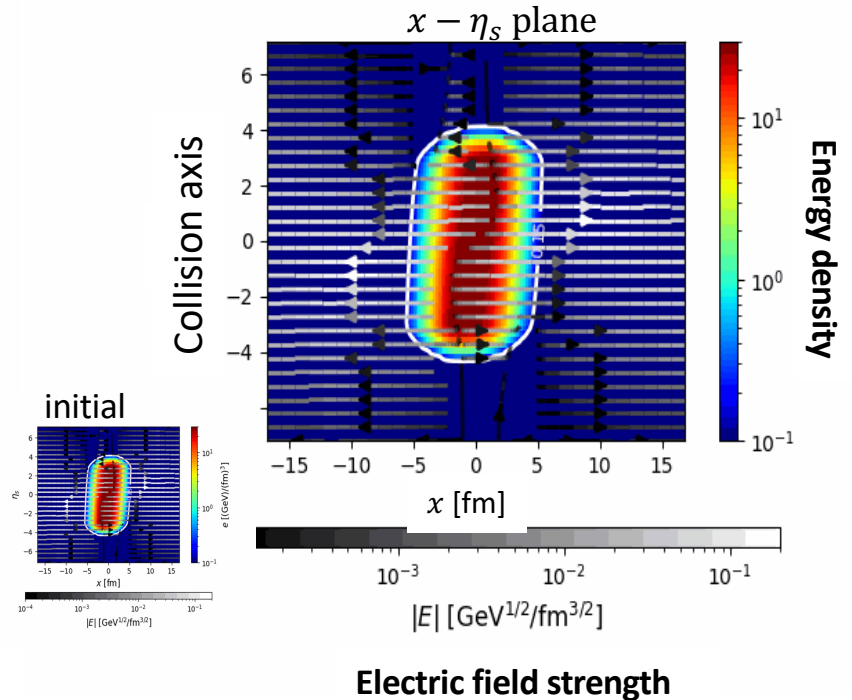
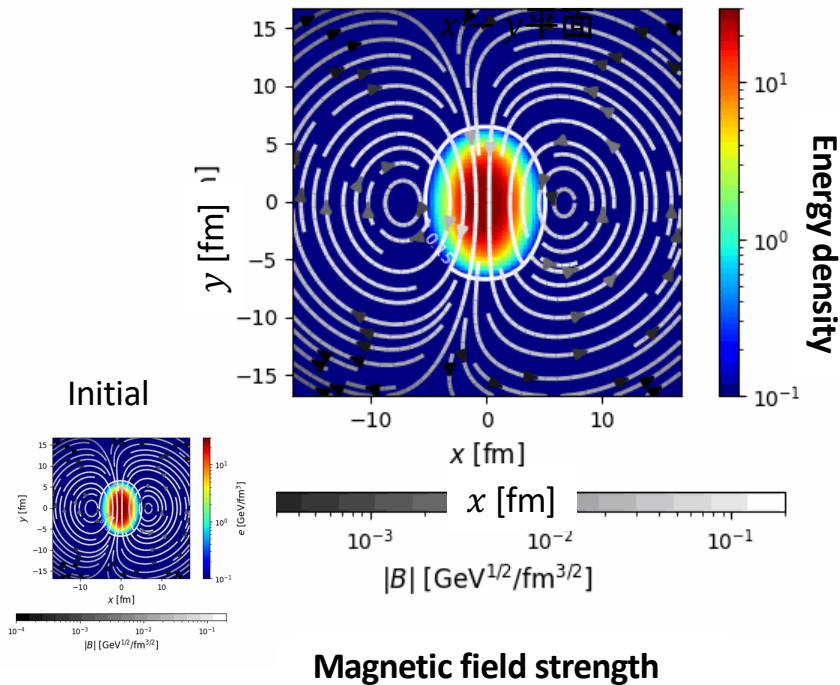
- Magnetic field
  - Strong magnetic field
- Electric field
  - $E \neq 0$  due to the asymmetry of the charge distribution

# Space-time Evolution

Nakamura, Miyoshi, C. N. and Takahashi, PRC 107, no.1, 014901 (2023)

## Au+Au collision system

First calculation in HIC with RRMHD code



Analysis of Heavy Ion Collisions

# Directed Flow

Nakamura, Miyoshi, C. N. and Takahashi, PRC 107, no.1, 014901 (2023)

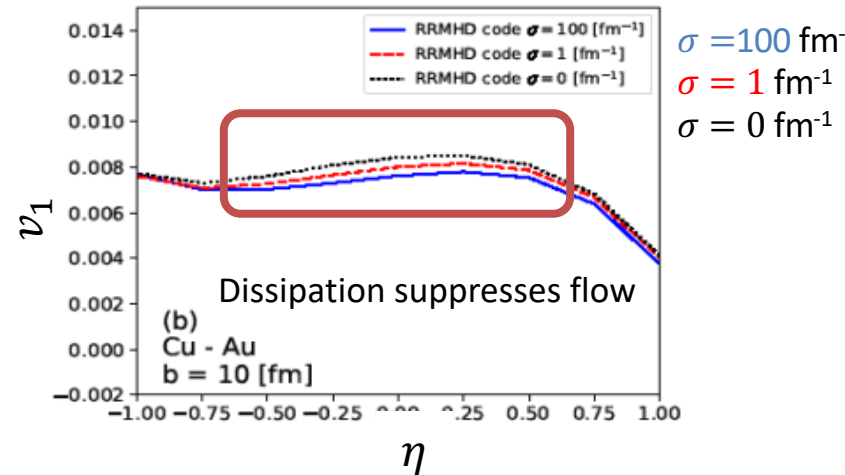
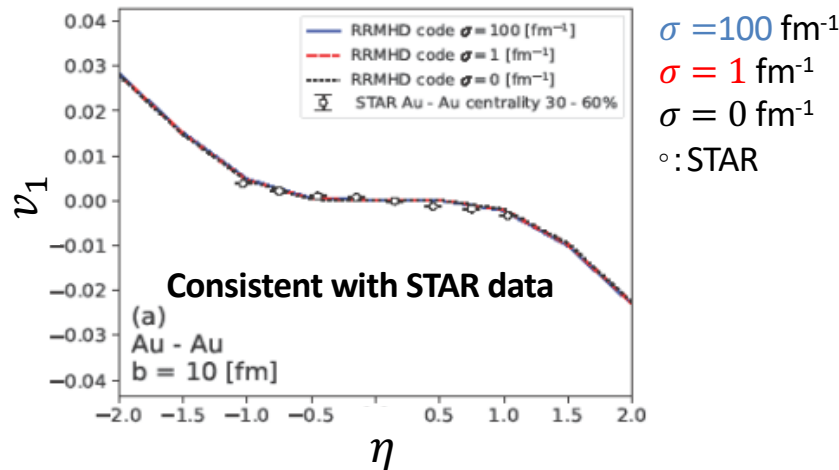
- $v_1 := \langle \cos(\phi - \Psi_1) \rangle \sim \left\langle \frac{p_x}{p_T} \right\rangle$

$$\eta = \frac{1}{2} \ln \frac{|p| + p_z}{|p| - p_z}$$

- Au-Au collisions ( $\sqrt{s_{NN}} = 200$  GeV)
  - Parameter fixed in initial condition from comparison with STAR data

- Cu-Au collisions ( $\sqrt{s_{NN}} = 200$  GeV)
  - Decreases with conductivity
  - Dissipation suppresses flow in the Cu direction

STAR Collaboration, Phys. Rev. Lett. 101 (2008), 252301



# Energy Transfer by Ohm Dissipation

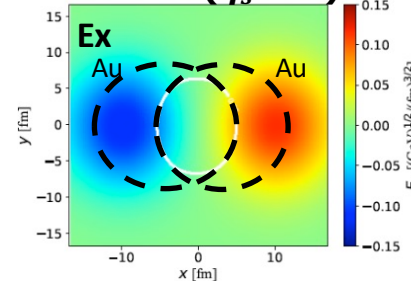
Nakamura, Miyoshi, C. N. and Takahashi, PRC 107, no.1, 014901 (2023)

## Energy Transfer

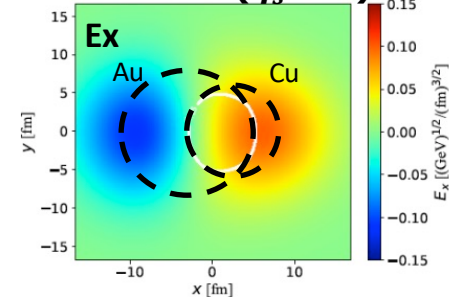
$$D(u) := j^\mu e_\mu = \gamma[\mathbf{j} \cdot (\mathbf{E} + \mathbf{v} \times \mathbf{B}) - q(\mathbf{v} \cdot \mathbf{E})]$$

energy of the electromagnetic field  $\rightarrow$  Thermal energy  
Kinetic energy

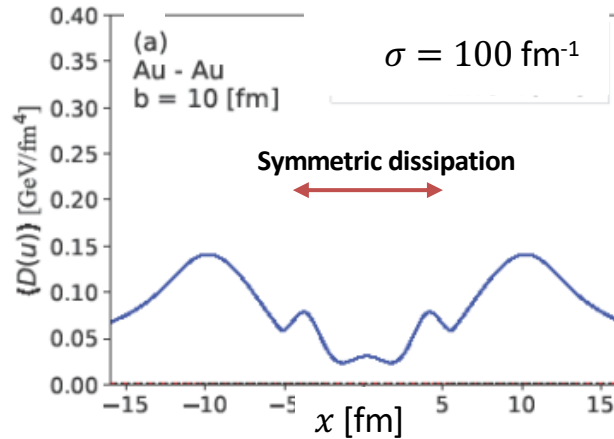
Au + Au ( $\eta_s = 0$ )



Cu + Au ( $\eta_s = 0$ )

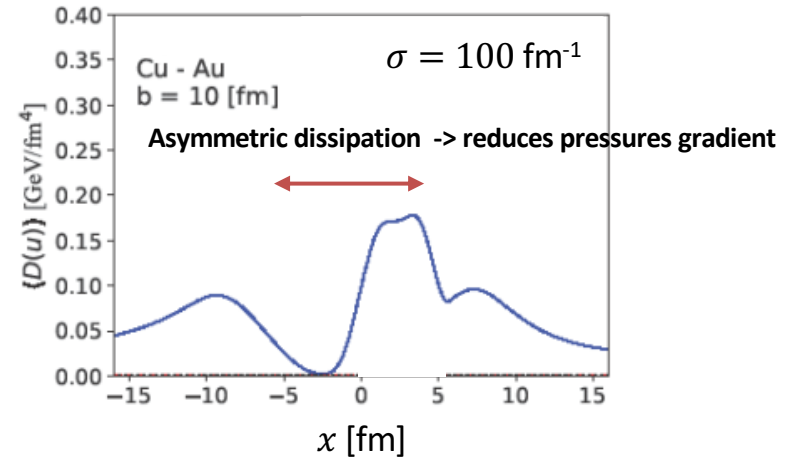


Au+Au collisions



no contribution to  $v_1$

Cu+Au collisions



contribution to  $v_1$



# Directed Flow

Nakamura, Miyoshi, C. N. and Takahashi, PRC 107, no.1, 014901 (2023)

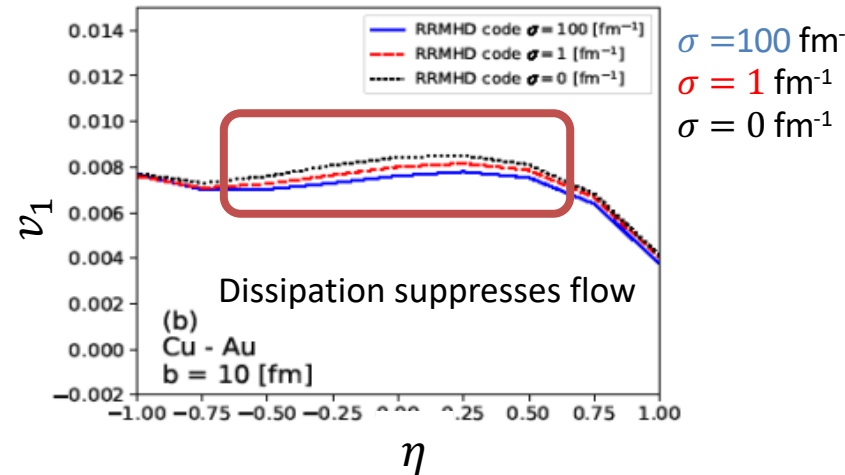
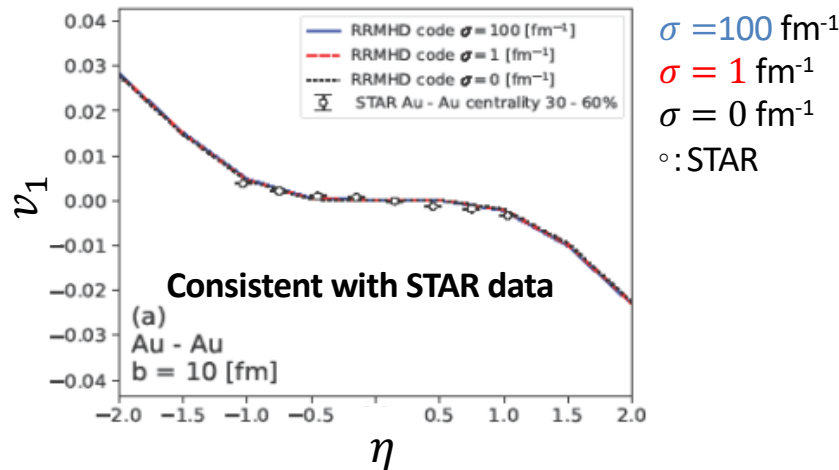
- $v_1 := \langle \cos(\phi - \Psi_1) \rangle \sim \left\langle \frac{p_x}{p_T} \right\rangle$

$$\eta = \frac{1}{2} \ln \frac{|p| + p_z}{|p| - p_z}$$

- Au-Au collisions ( $\sqrt{s_{NN}} = 200$  GeV)
  - Parameter fixed in initial condition from comparison with STAR data

- Cu-Au collisions ( $\sqrt{s_{NN}} = 200$  GeV)
  - Decreases with conductivity
  - Dissipation suppresses flow in the Cu direction

STAR Collaboration, Phys. Rev. Lett. 101 (2008), 252301





# Charge Dependence of $\Delta v_2$ : Au + Au

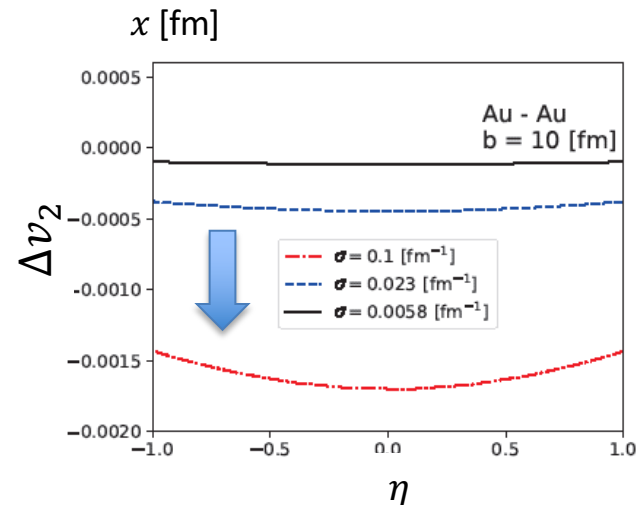
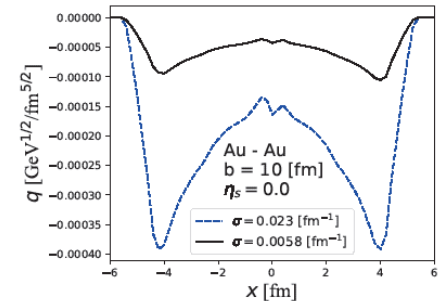
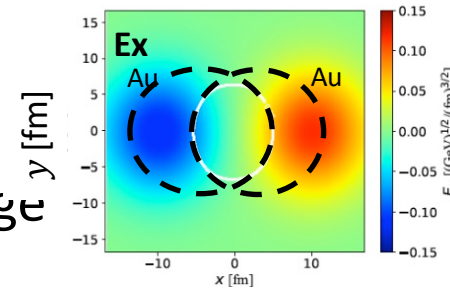
Nakamura, Miyoshi, C. N. and Takahashi, PRC 107, no.1, 014901 (2023)

- $\Delta v_2 = v_2^{\pi^+}(\eta) - v_2^{\pi^-}(\eta)$

## – Negative Elliptic Flow

- Contribution of negative charge  $\tau$  on freezeout hypersurface
- Symmetric structure: initial electric field to the collision axis
- Electric conductivity dependence is observed even in the symmetry system.

Charge distribution on freezeout hypersurface



$\sigma = 0.1 \text{ fm}^{-1}$   
 $\sigma = 0.023 \text{ fm}^{-1}$   
 $\sigma = 0.0058 \text{ fm}^{-1}$

$$\eta = \frac{1}{2} \ln \frac{|p| + p_z}{|p| - p_z}$$

# Charge Dependence of $\Delta v_2$ : Cu + Au

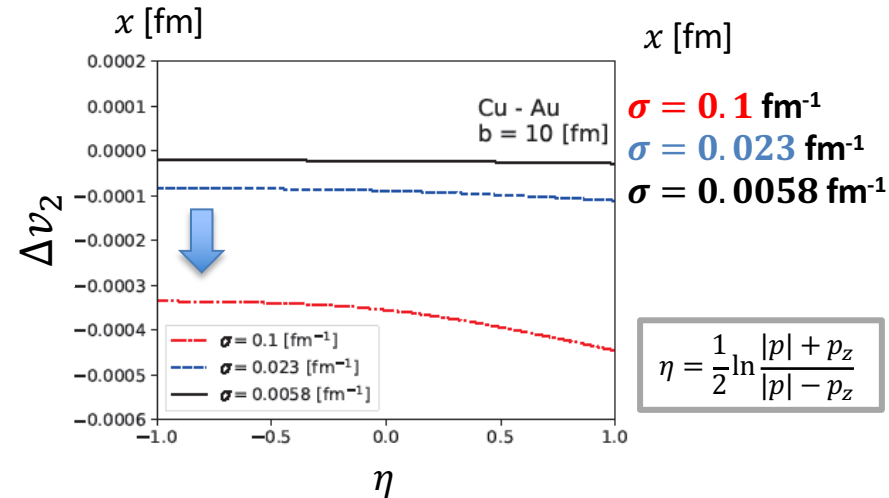
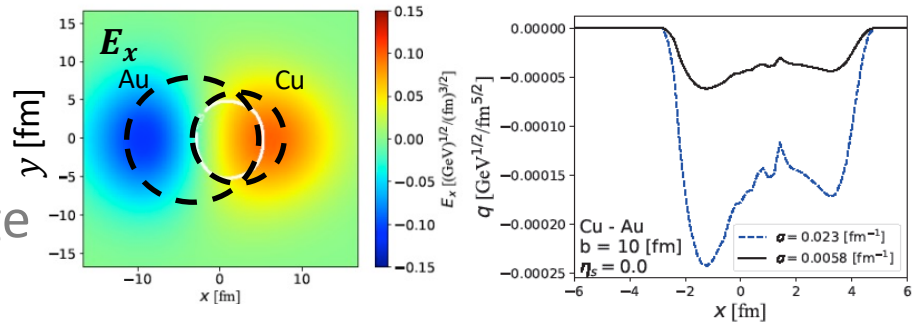
Nakamura, Miyoshi, C. N. and Takahashi, PRC 107, no.1, 014901 (2023)

- $\Delta v_2 = v_2^{\pi^+}(\eta) - v_2^{\pi^-}(\eta)$

## – Negative Elliptic Flow

- Contribution of negative charge on freezeout hypersurface
- Asymmetric structure: initial electric field to the collision axis
- Electric conductivity dependence is observed.

Charge distribution on freezeout hypersurface



$\Delta v_2$ : initial electromagnetic field distribution  
electrical conductivity



# Charge Dependence of $\Delta v_1$ : Au + Au

Nakamura, Miyoshi, C. N. and Takahashi, arXiv:2212.02124 [nucl-th]

- $\Delta v_1 = v_1^{\pi^+}(\eta) - v_1^{\pi^-}(\eta)$

- Clear dependence of charge conductivity

- Proportion to electric conductivity
- Negative charge induced in the opposite direction of fluid flow

suppression of  $v_1$  of negative charge

- $\Delta v_1$  with finite  $\sigma$  is consistent with STAR data

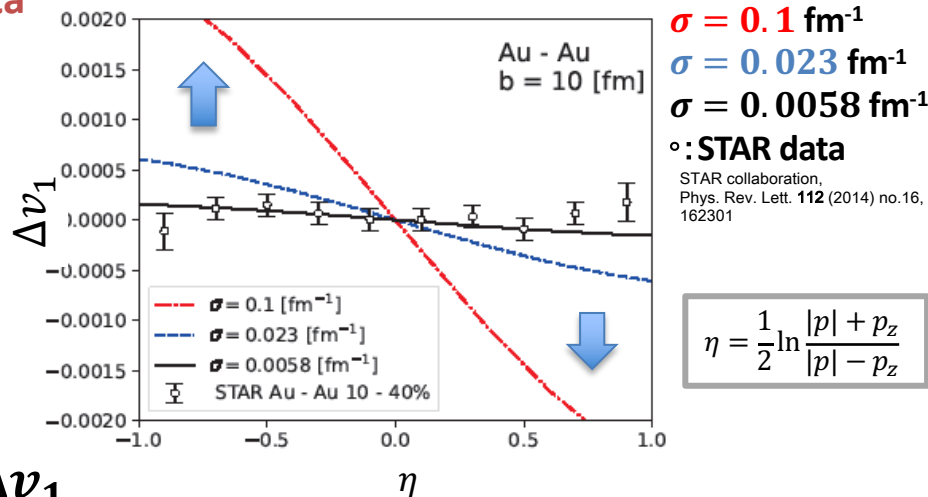
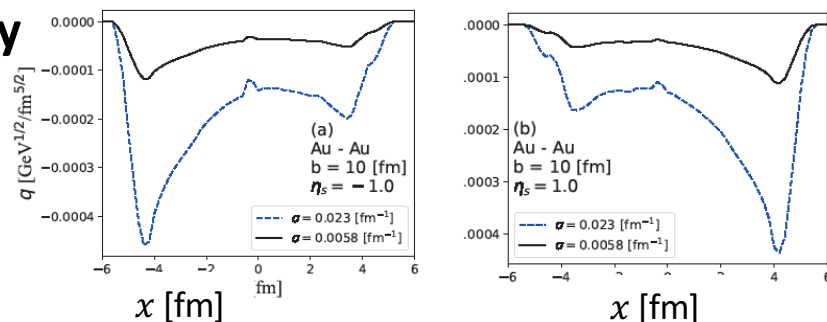
- $\sigma = 0.0058 \text{ fm}^{-1}$

ex.  $\sigma_{LQCD} = 0.023 \text{ fm}^{-1}$

Gert Aarts, et al. from lattice QCD  
*Phys. Rev. Lett.*, 99:022002, 2007.

✓ QGP electrical conductivity from high-precision measurement of  $\Delta v_1$

Charge distribution on freezeout hypersurface



# Charge Dependence of $\Delta v_1$ : Au + Au

Nakamura, Miyoshi, C. N. and Takahashi, arXiv:2212.02124 [nucl-th]

- $\Delta v_1 = v_1^{\pi^+}(\eta) - v_1^{\pi^-}(\eta)$

- Electric field created by initial condition

- $\Delta v_1$  is finite at  $\eta = 0$

- Asymmetry structure to  $\eta = 0$

- Proportion to electric conductivity

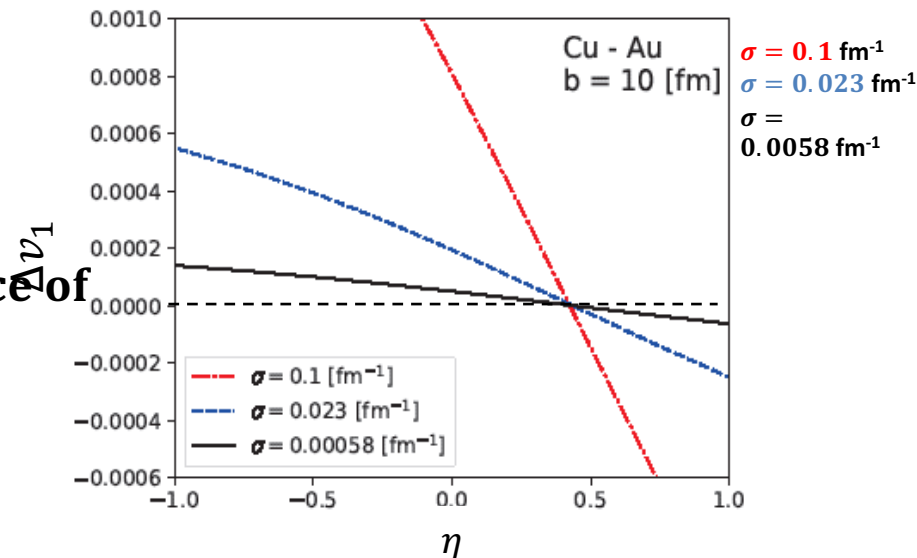
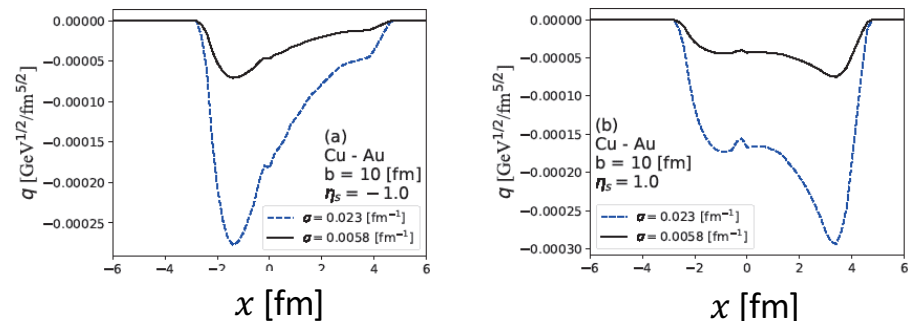
- $\Delta v_1$  vanishes at  $\eta = 0.5$ .

- ✓ Electrical conductivity  $\leftarrow \Delta v_1$  at  $\eta = 0$

- ✓ Initial electrical field from  $\eta$  dependence of

$\Delta v_1$

Charge distribution on freezeout hypersurface



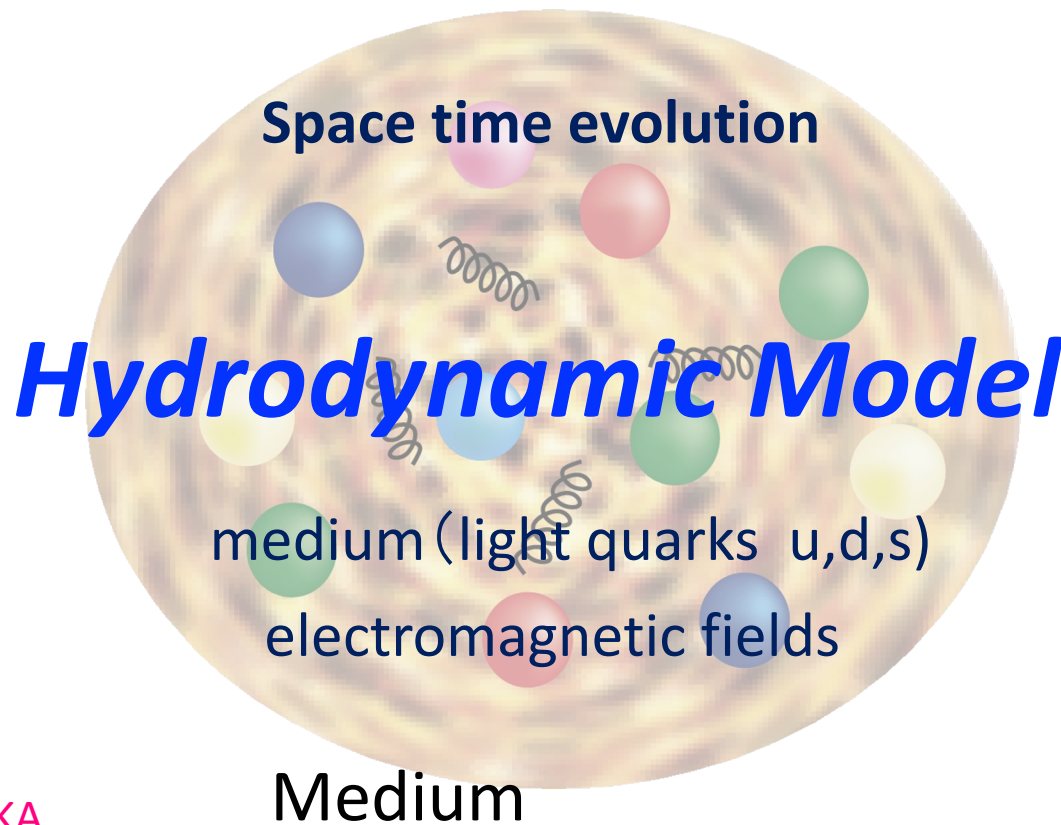
Asymmetric system had advantage in explore of QGP electrical conductivity.



# Summary

- **Tools for analyses of relativistic heavy ion collisions**
  - New relativistic viscous hydrodynamics code
  - New relativistic resistive hydrodynamic model

## Quantitative Analysis on QGP bulk property



temperature  
dependence of  
viscosities

Electric  
conductivity



# Summary

- **Tools for analyses of relativistic heavy ion collisions**
  - Application to other physical observables
    - Jets, heavy quarks, photons, electromagnetic probes...

