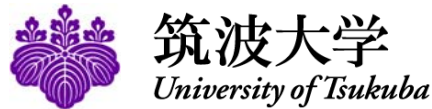


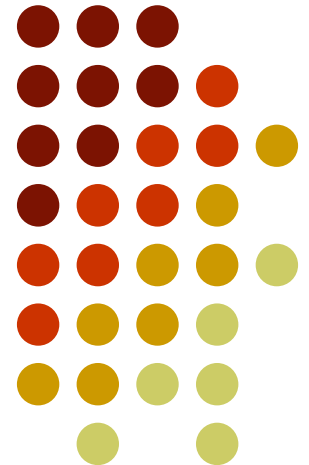
A constant mode in charmonium correlators in finite temperature lattice QCD

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*Tokyo Metropolitan University, Tokyo,
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Introduction



J/ψ suppression is one of the most promising probe to find the QGP formation in HIC experiment.

Lattice QCD studies of charmonium spectral function suggest the survival of J/ψ state above T_c ($1.5T_c$?)

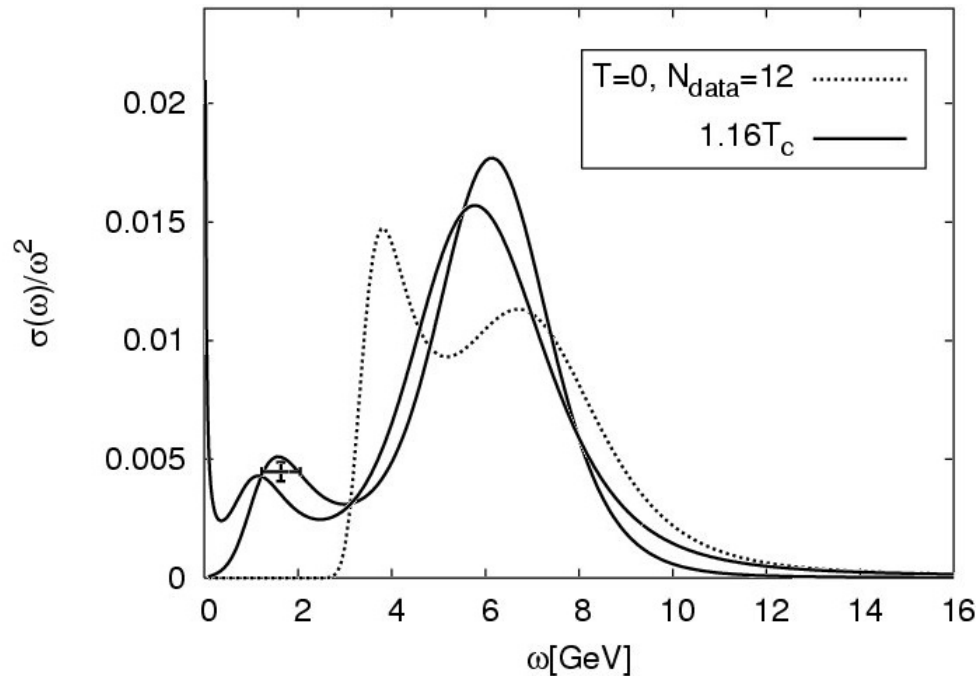
■ Indirect (sequential) J/ψ suppression

total yield of J/ψ =
direct production of J/ψ (60%)
+ decay from higher states, ψ' & χ_c (40%)

L. Antoniazzi et al. (E705 Collab.), PRL 70, 383, (1993).

→ A part of the J/ψ suppression may be observed at $T_{\text{dis.}}(\psi' \text{ or } \chi_c)$ when $T_{\text{dis.}}(\psi' \text{ or } \chi_c) < T_{\text{dis.}}(J/\psi)$

χ_c states on the lattice



*S. Datta et al.,
PRD69, 094507 (2004).
A. Jakovac et al.,
PRD75, 014506 (2007).*



FIG. 19: The scalar spectral function for $\beta = 6.1$ at $T = 1.16T_c$ and at zero temperature reconstructed using $N_{data} = 12$. At finite temperature two default models $m(\omega) = 0.01$ and $m(\omega) = 0.038\omega^2$ have been used.

Quenched QCD at $T > 0$



Lattice setup

anisotropic lattices : $20^3 \times N_t$

$1/a_s = 2.03(1) \text{ GeV}$, $a_s/a_t = 4$

Clover quark action with tadpole imp. on anisotropic lattice

H. Matsufuru et al., PRD64, 114503 (2001).

$r_s=1$ to reduce cutoff effects in higher energy states

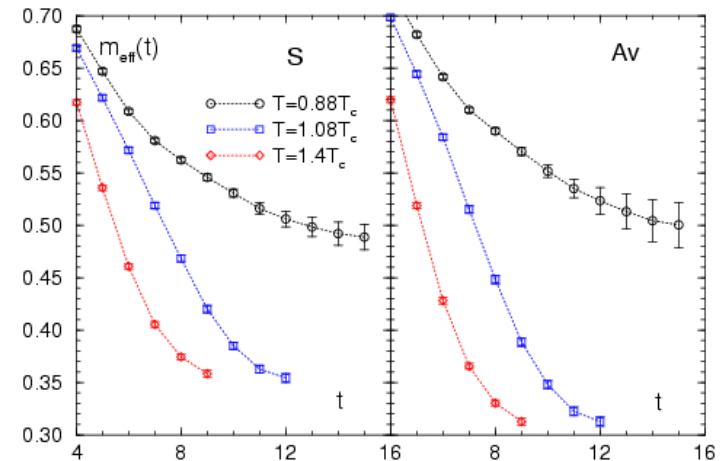
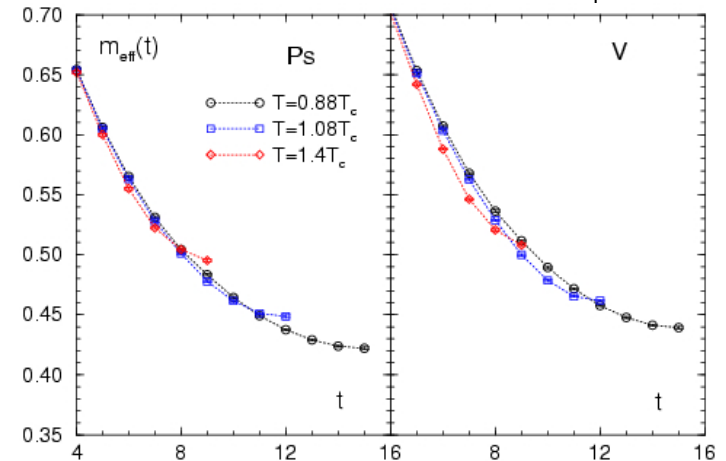
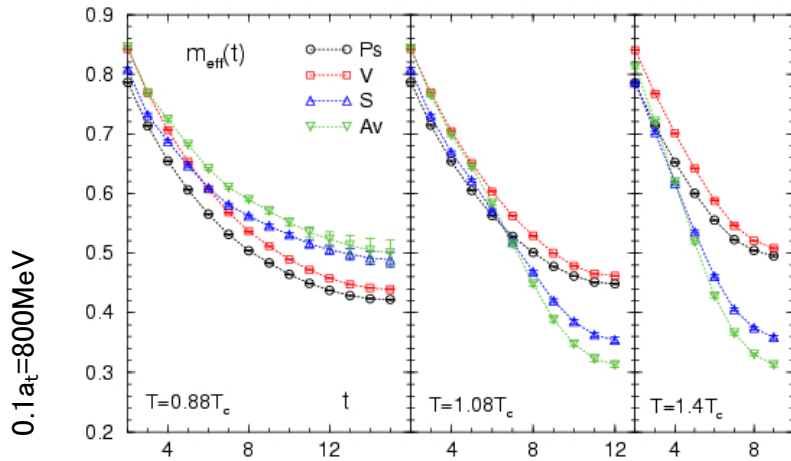
F. Karsch et al., PRD68, 014504 (2003).

quark mass is tuned with $M_{J/\psi} (= 3097 \text{ MeV})$

N_τ	160	32	26	20
T/T_c	~ 0	0.88	1.08	1.4
# of conf.	60	300	300	300

equilib. is 20K sweeps
each config. is separated
by 500 sweeps

Quenched QCD at $T > 0$



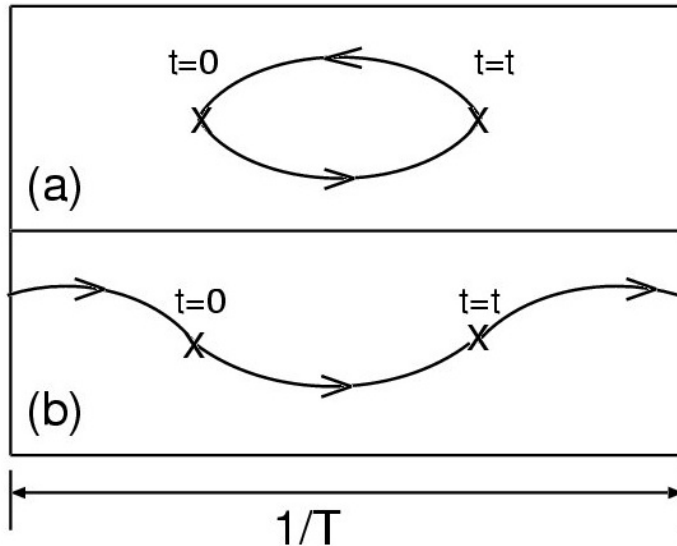
- small change in S-wave states
→ survival of J/ψ & η_c at $T > T_c$
- drastic change in P-wave states
→ dissociation of χ_c just above T_c (?)

*S. Datta et al.,
PRD69, 094507 (2004). etc...*



A constant mode

Now we consider the meson correlator with $p=0$ & $m_{q1}=m_{q2}$



$$\exp(-m_q t) \times \exp(-m_q t) = \exp(-2m_q t)$$

m_q is quark mass
or single quark energy

$$\exp(-m_q t) \times \exp(-m_q(L_t - t)) = \exp(-m_q L_t)$$

$L_t =$ temporal extent

- in imaginary time formalism

$$L_t = 1/\text{Temp.}$$

gauge field : periodic b.c.

quark field : anti-periodic b.c.

- in confined phase: m_q is infinite

→ the effect appears

only in deconfined phase

Pentaquark (KN state):
two pion state:
→ Dirichlet b.c.

*c.f. T.T.Takahashi et al.,
PRD71, 114509 (2005).*



Physical interpretation

Spectral function at high temp. limit

$$\rho_{\Gamma}(\omega) = \Theta(\omega^2 - 4m_q^2) \frac{N_c}{8\pi\omega} \sqrt{\omega^2 - 4m_q^2} [1 - 2n_F(\omega/2)] \\ \times [\omega^2 (a_H^{(1)} - a_H^{(2)}) + 4m^2 (a_H^{(2)} - a_H^{(3)})] \\ + 2\pi\omega\delta(\omega) N_c [(a_H^{(1)} + a_H^{(2)})I_1 + (a_H^{(2)} - a_H^{(3)})I_2]$$

*F. Karsch et al.,
PRD68, 014504 (2003).
G. Aarts et al.,
NPB726, 93 (2005).*

	Γ	$a_H^{(1)} + a_H^{(2)}$	$a_H^{(2)} - a_H^{(3)}$
Ps	γ_5	0	0
V	γ_i	2	2
S	1	0	-2
AV	$\gamma_i\gamma_5$	2	-4

constant mode remains
in the continuum & infinite volume

The constant term is related to some transport coefficients.

From Kubo-formula, for example, a derivative of the SPF in the V channel is related to the electrical conductivity σ .

$$\sigma = \frac{1}{6} \frac{\partial}{\partial \omega} \rho_V(\omega) \Big|_{\omega=0}$$

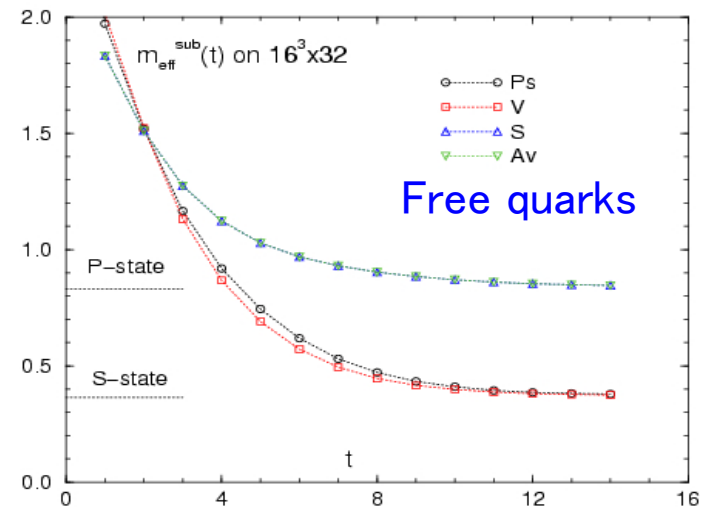
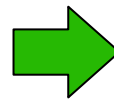
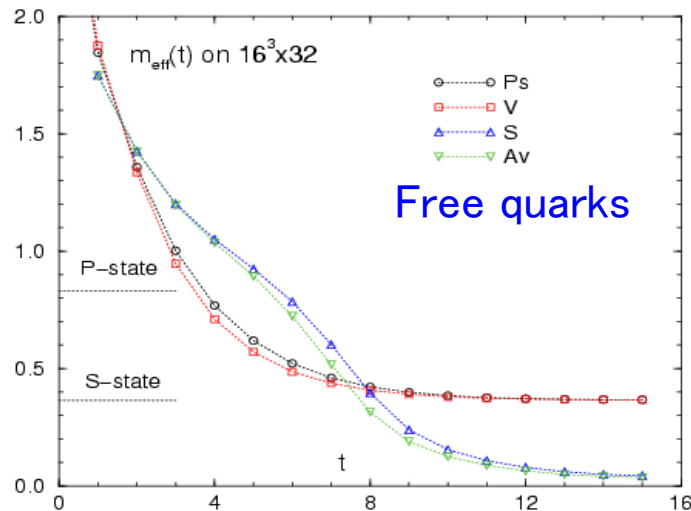


Removing the constant mode

An analysis to avoid the constant mode

Midpoint subtracted correlator

$$\bar{C}(t) = C(t) - C(N_t/2) \quad \rightarrow \quad \frac{\bar{C}(t)}{\bar{C}(t+1)} = \frac{\sinh^2 \left[\frac{1}{2} m_{eff}^{sub}(t) (N_t/2 - t) \right]}{\sinh^2 \left[\frac{1}{2} m_{eff}^{sub}(t) (N_t/2 - t - 1) \right]}$$



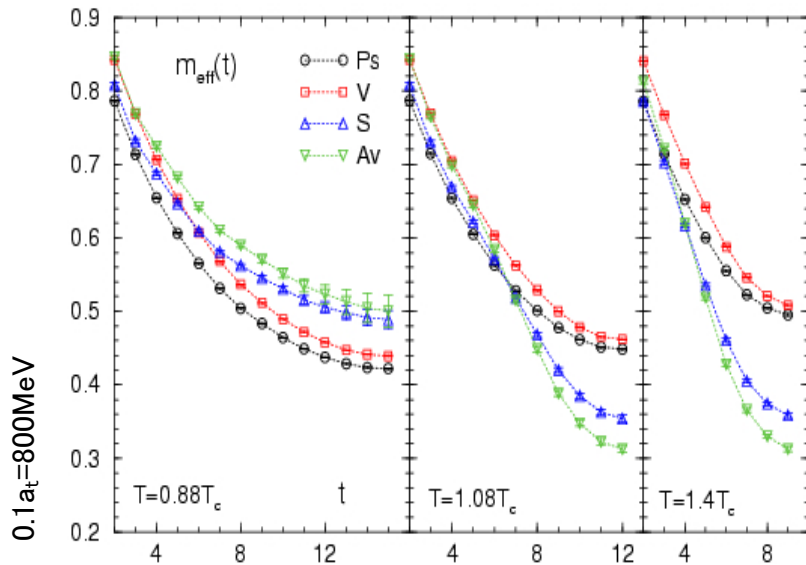


Midpoint subtraction analysis

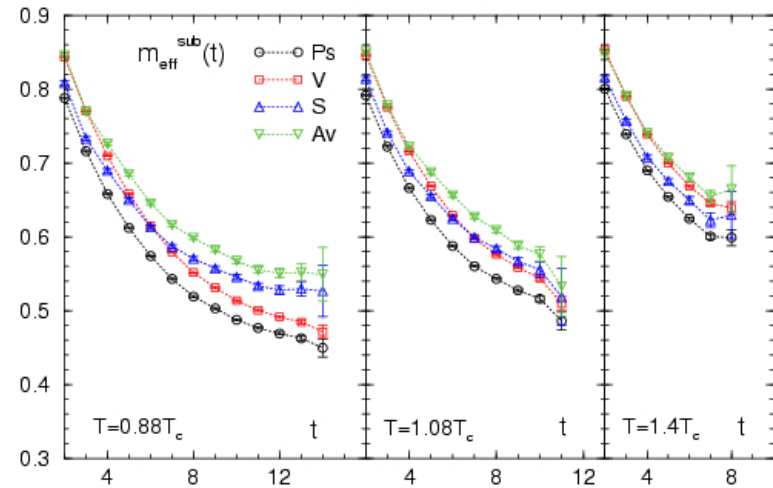
$$\bar{C}(t) = C(t) - C(N_t/2) \quad \frac{\bar{C}(t)}{\bar{C}(t+1)} = \frac{\sinh^2 \left[\frac{1}{2} m_{eff}^{sub}(t) (N_t/2 - t) \right]}{\sinh^2 \left[\frac{1}{2} m_{eff}^{sub}(t) (N_t/2 - t - 1) \right]}$$



usual effective masses at $T > 0$



subtracted effective mass



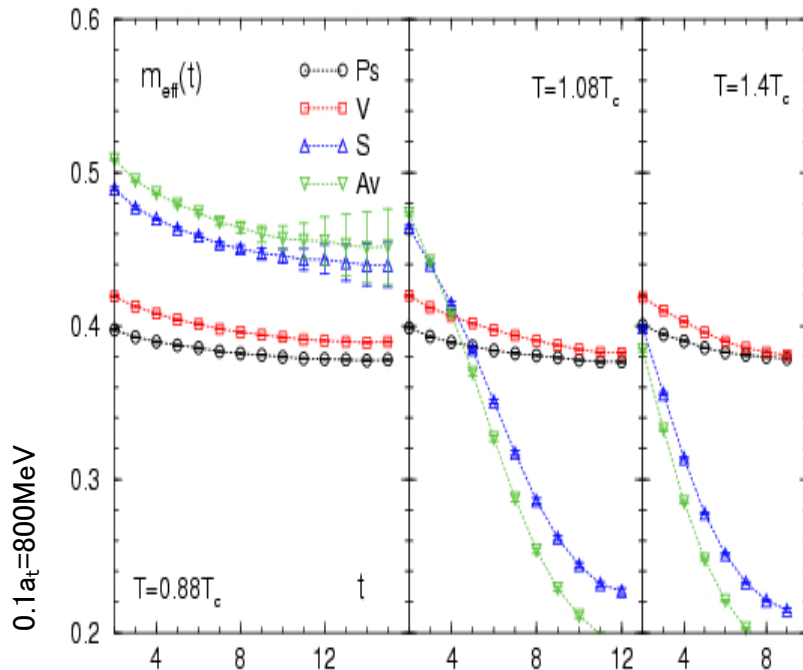
the drastic change in P-wave states disappears in $m_{eff}^{sub}(t)$

→ the change is due to the constant mode

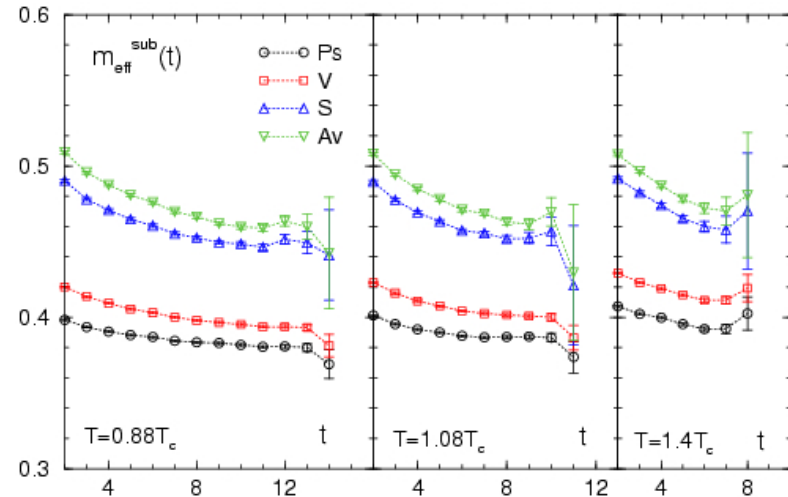
Results with extended op.



usual effective mass



subtracted effective mass



- extended op. enhances overlap with const. mode
- small constant effect is visible in V channel
- no large change above T_c in $m_{\text{eff}}^{\text{sub}}(t)$



MEM analysis fails ?

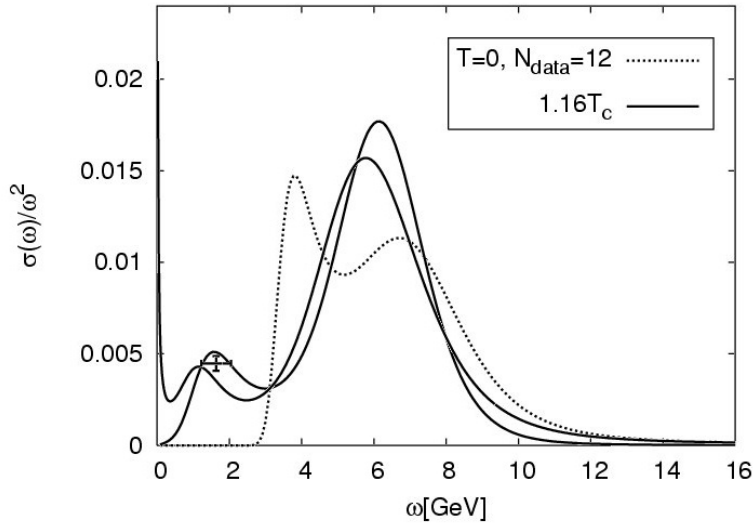
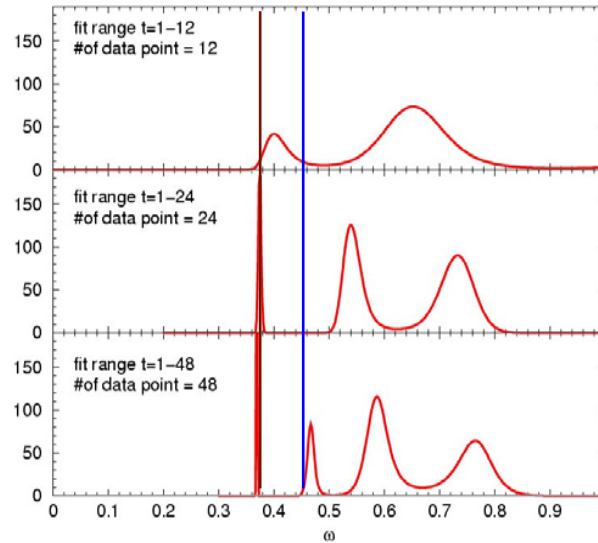


FIG. 19: The scalar spectral function for $\beta = 6.1$ at $T = 1.16T_c$ and at zero temperature reconstructed using $N_{data} = 12$. At finite temperature two default models $m(\omega) = 0.01$ and $m(\omega) = 0.038\omega^2$ have been used.

A. Jakovac et al., PRD75, 014506 (2007).
(also S. Datta et al., PRD69, 094507 (2004).)

MEM test using T=0 data



data
for $T/T_c=1.2$

data
for $T/T_c=0.6$

data
for $T/T_c=0$

MEM sometimes fails
when (# or quality) of data point
is not sufficient.



Conclusion

- There is the constant mode in charmonium correlators above T_c
- The drastic change in χ_c states is due to the constant mode
→ the survival of χ_c states above T_c , at least $T=1.4T_c$.

The result may affect the scenario of J/ψ suppression.

In the MEM analysis,
one has to check consistency of the results
using, e.g., midpoint subtracted correlators.

$$\bar{C}(t) = C(t) - C(N_t/2)$$
$$\bar{C}(t) = \int_0^\infty d\omega \rho_\Gamma(\omega) K^{sub}(\omega, t),$$
$$K^{sub}(\omega, t) = \frac{\sinh^2(\frac{\omega}{2}(N_t/2 - t))}{\sinh(\omega N_t/2)}$$