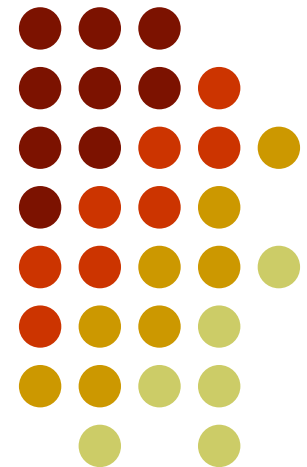


QCD thermodynamics with $N_f=2+1$ near the continuum limit at realistic quark masses

Takashi Umeda (BNL)
for the RBC – Bielefeld Collaboration



Lattice 2006 Tucson AZ 23–28 July 2006



Motivation & Approach



Quantitative study of QCD thermodynamics
from first principle calculation (Lattice QCD)

T_c , EoS, phase diagram, small μ , etc...



from recent studies, we know
these quantities strongly depend on m_q & N_f

Our aim is QCD thermodynamics with 2+1 flavor
at almost realistic quark masses

e.g. pion mass $\approx 150\text{MeV}$, kaon mass $\approx 500\text{MeV}$

- Choice of quark action
 - Improved Staggered quark action
- Continuum limit
 - $N_t = 4, 6, (8) \rightarrow a \approx 0.24, 0.17, (0.12) \text{ fm}$

Computers



US/RBRC QCDOC

20.000.000.000.000 ops/sec



BI – apeNEXT

5.000.000.000.000 ops/sec

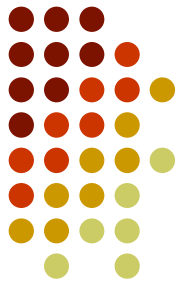


- critical temperature
- equation of state
- hadron properties in matter

today: ~ TFlops

<http://www.quark.phy.bnl.gov/~hotqcd>

Choice of Lattice Action



Improved Staggered action : p4fat3 action

Karsch, Heller, Sturm (1999)

- gluonic part : Symanzik improvement scheme
 - remove cut-off effects of $O(a^2)$
 - tree level improvement $O(g^0)$
- fermion part : improved staggered fermion
 - remove cut-off effects & improve rotational sym.
 - improve flavor symmetry by smeared 1-link term

$$S_F(N_\tau, N_\sigma) = \sum_{n, \hat{n}} \sum_{\mu} \eta(n_\mu) \bar{\chi}_n \left(\frac{3}{8} \left[\frac{1}{1+6\omega} \left(\leftarrow \circ \rightarrow + \omega \sum_{\nu \neq \mu} \left[\begin{array}{c} \uparrow \downarrow \\ \leftarrow \circ \rightarrow \\ \uparrow \downarrow \end{array} \right] \right) \right] \right) \left(\frac{1}{48} \sum_{\nu \neq \mu} \left[\begin{array}{c} \uparrow \downarrow \\ \leftarrow \circ \rightarrow \\ \uparrow \downarrow \\ \leftarrow \circ \rightarrow \\ \uparrow \downarrow \end{array} \right] \right) \chi_{n'} + m_q \sum_n \bar{\chi}_n \chi_n$$

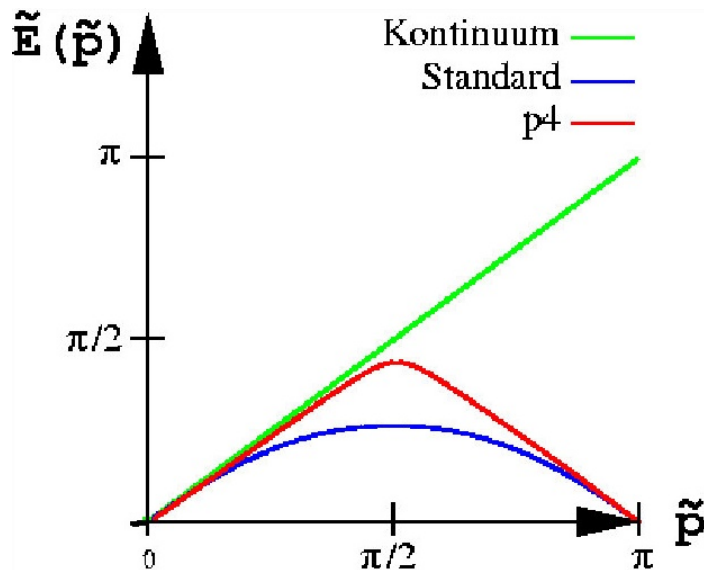
fat3

p4

Properties of the $p4$ -action

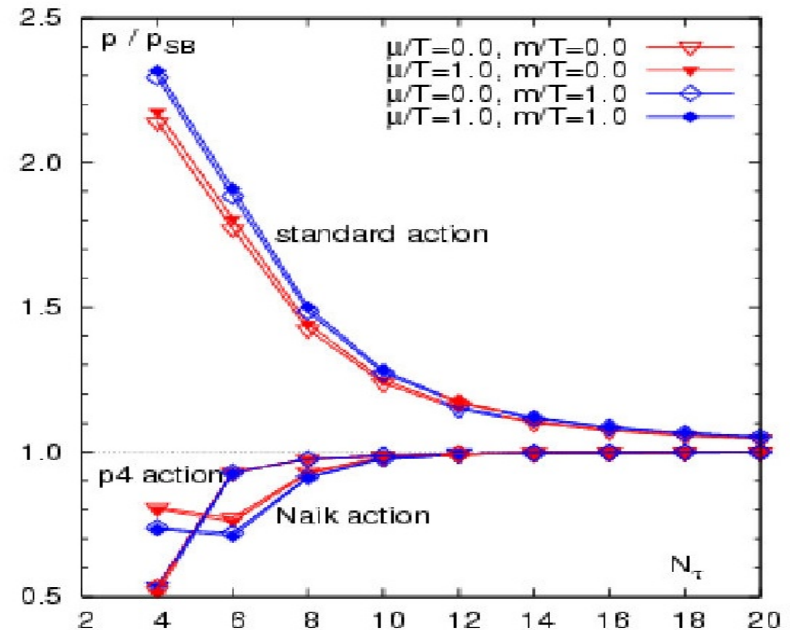


Dispersion relation



The free quark propagator is rotational invariant up to $O(p^4)$

pressure in high T limit



Bulk thermodynamic quantities show drastically reduced cut-off effects

flavor sym. is also improved by fat link

Contents of this talk



- Motivation and Approach
- Choice of lattice action
- **Critical temperature**
 - Simulation parameters
 - Critical β search
 - Scale setting by Static quark potential
 - Critical temperature
- **Spatial string tension**
- Conclusion

Simulation parameters



■ Critical β search at $T > 0$

N_τ	\hat{m}_s	\hat{m}_l	V	# β values	max.# conf.
4	0.1	$0.5 \hat{m}_s$	8^3	10	40,000
		$0.2 \hat{m}_s$	8^3	6	12,000
4	0.065	$0.4 \hat{m}_s$	$8^3, 16^3$	10, 11	30,000, 60,000
		$0.2 \hat{m}_s$	$8^3, 16^3$	8, 7	30,000, 60,000
		$0.1 \hat{m}_s$	$8^3, 16^3$	9, 6	34,000, 50,000
		$0.05 \hat{m}_s$	$8^3, 16^3$	8, 5	30,000, 42,000
6	0.0040	$0.4 \hat{m}_s$	16^3	11	20,000
		$0.2 \hat{m}_s$	16^3	9	60,000
		$0.1 \hat{m}_s$	16^3	7	60,000

(* conf. = 0.5 MD traj.

to check
 m_s dependence for T_c

■ $T=0$ scale setting at $\beta_c(N_t)$

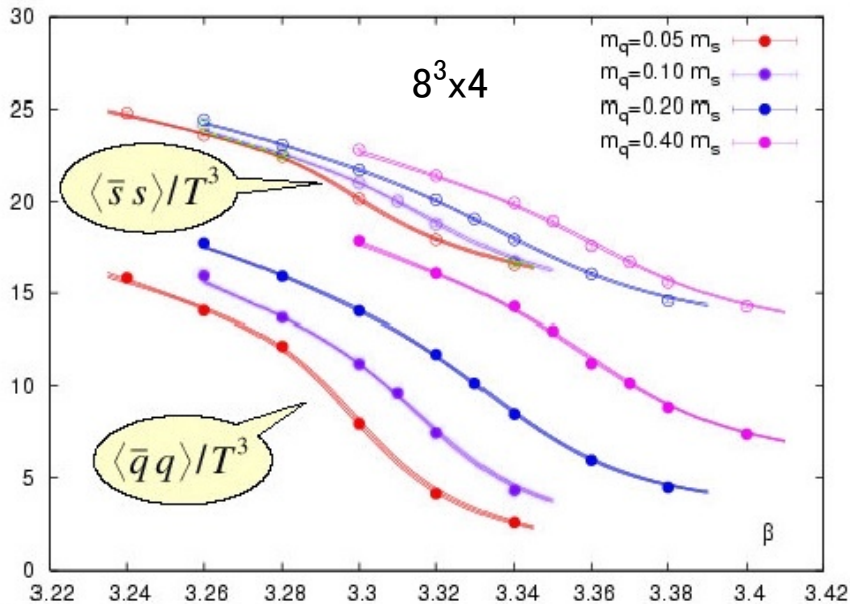
N_τ	\hat{m}_s	\hat{m}_l	β	# conf.	m_{ps}/m_v	a [fm]
4	0.1	$0.5 \hat{m}_s$	3.409	600	0.520(2)	0.2273(4)
		$0.2 \hat{m}_s$	3.371	238	0.372(5)	0.2336(7)
4	0.065	$0.4 \hat{m}_s$	3.362	500	0.410(2)	0.2312(7)
		$0.2 \hat{m}_s$	3.335	400	0.303(7)	0.2365(6)
		$0.1 \hat{m}_s$	3.310	750	0.212(7)	0.2458(5)
		$0.05 \hat{m}_s$	3.300	400	0.154(5)	0.2475(8)
6	0.0040	$0.4 \hat{m}_s$	3.500	294	0.461(4)	0.1558(7)
		$0.2 \hat{m}_s$	3.470	500	0.343(6)	0.1617(5)
		$0.1 \hat{m}_s$	3.455	410	0.248(4)	0.1670(5)

(* conf. = 5 MD traj.
after thermalization

Critical β search

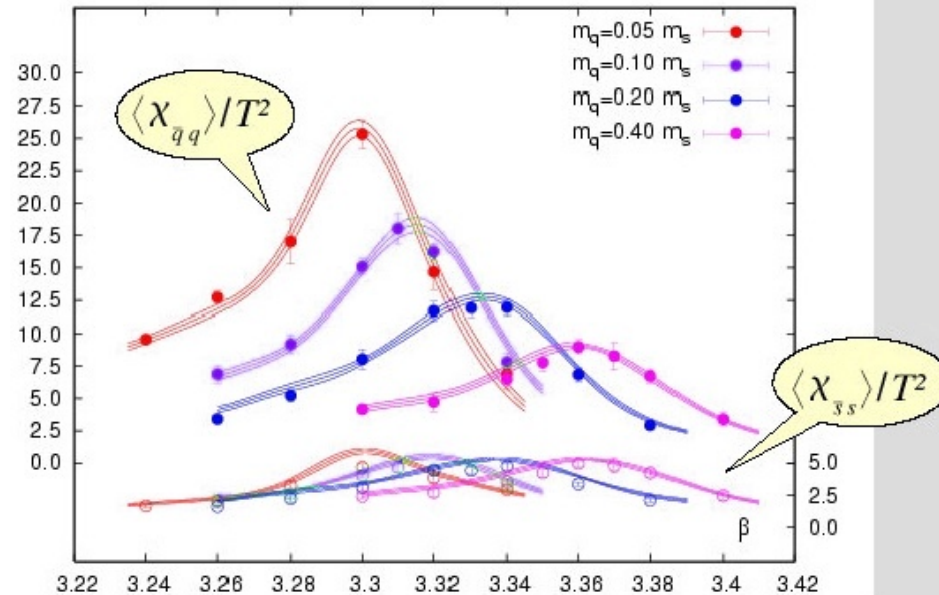


chiral condensate:



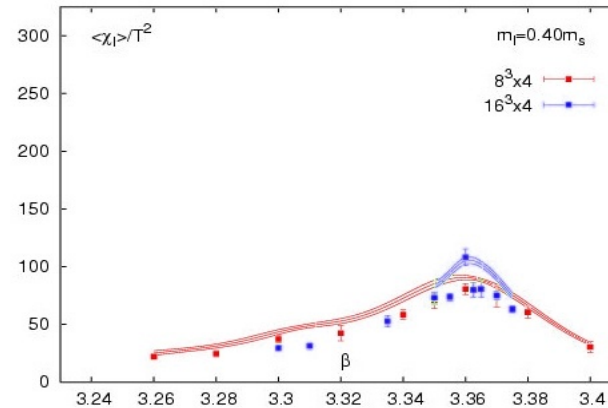
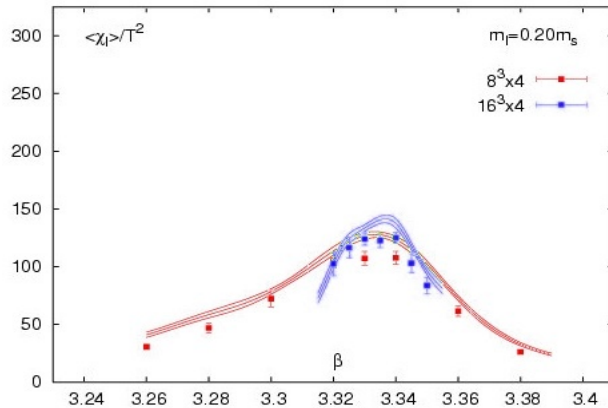
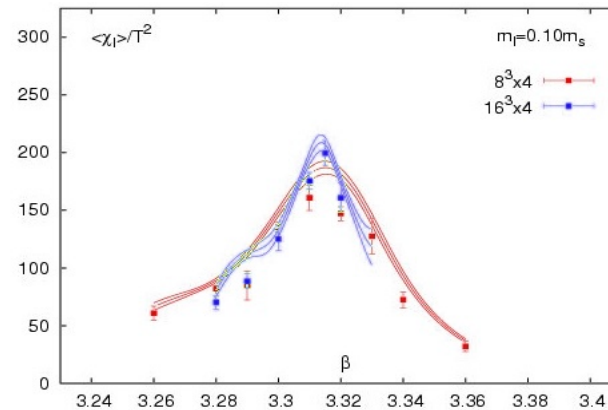
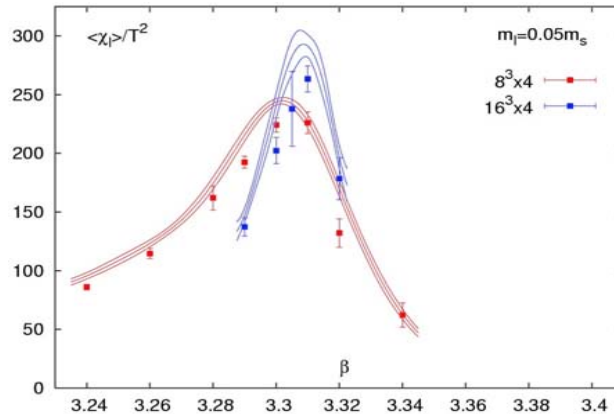
chiral susceptibility:

$$\langle \chi_{\bar{q}q} \rangle \equiv \langle (\bar{q}q)^2 \rangle - \langle \bar{q}q \rangle^2$$



- multi-histogram method (Ferrenberg–Swendsen) is used
- β_c are determined by peak positions of the susceptibilities
- Transition becomes stronger for smaller light quark masses

Volume dependence of β_c



- No large change in peak height & position
 - consistent with crossover transition rather than true transition

Uncertainties in β_c



- Statistical error

→ jackknife analysis for peak-position of susceptibility

- We can find a difference between β_I and β_L

→ small difference but statistically significant

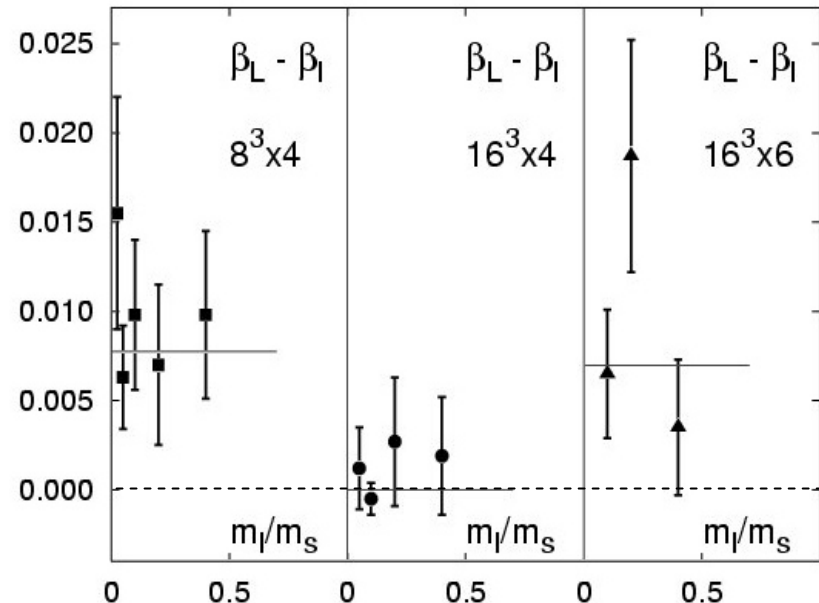
β_I : peak position of χ_I

β_L : peak position of χ_L

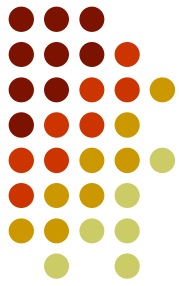
- the difference is negligible at $16^3 \times 4$ ($N_s/N_t=4$)

- no quark mass dependence

- the difference at $16^3 \times 6$ are taken into account as a systematic error in β_c

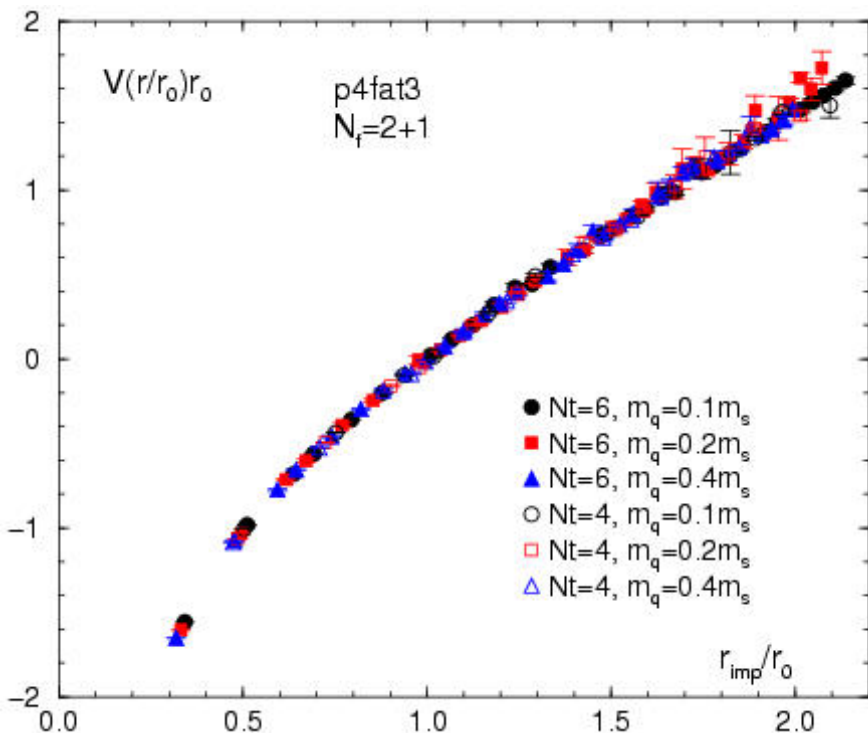


Scale setting at $T=0$



Lattice scale is determined by a static quark potential $V(r)$

$$r^2 \frac{dV_{\bar{q}q}(r)}{dr} \Big|_{r=r_0} = 1.65$$



$$V_3(r) = -\frac{\alpha}{r_I} + \sigma r_I + C$$

$$V_4(r) = -\frac{\alpha}{r} + \sigma r - \alpha' \left(\frac{1}{r_I} - \frac{1}{r} \right) + C$$

where, r_I is the improve dist.

■ statistical error

→ jackknife error

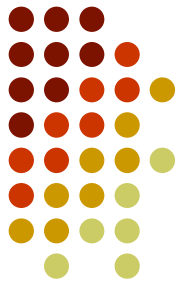
■ systematic errors

→ diff. between $V_3(r)$ & $V_4(r)$

diff. in various fit range

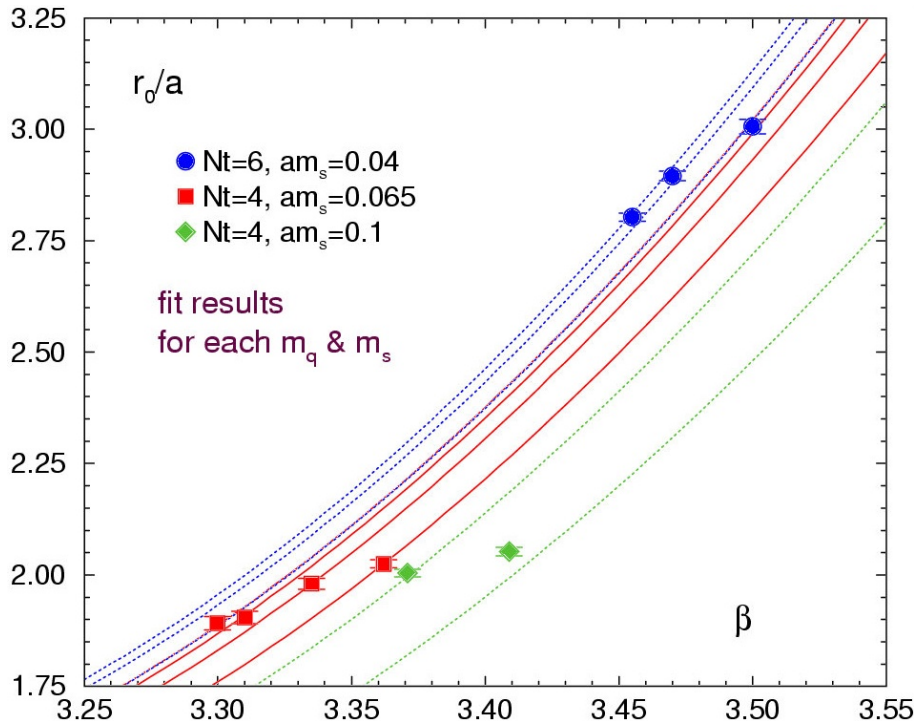
($r_{\text{min}}=0.15-0.3\text{fm}$, $r_{\text{max}}=0.7-0.9\text{fm}$)

β & m_q dependence of r_0



RG inspired ansatz with 2-loop beta-function $R(\beta)$

$$\ln(r_0/a) = A(2m_l + m_s) - \ln(R(\beta)) + B \left(\frac{R(\beta)}{R(3.4)} \right)^2 + C \left(\frac{R(\beta)}{R(3.4)} \right)^4 + D$$



$$A = -1.53(11), \quad B = -0.88(20)$$

$$C = 0.25(10), \quad D = -2.45(10)$$

$$\chi^2/dof = 1.2$$

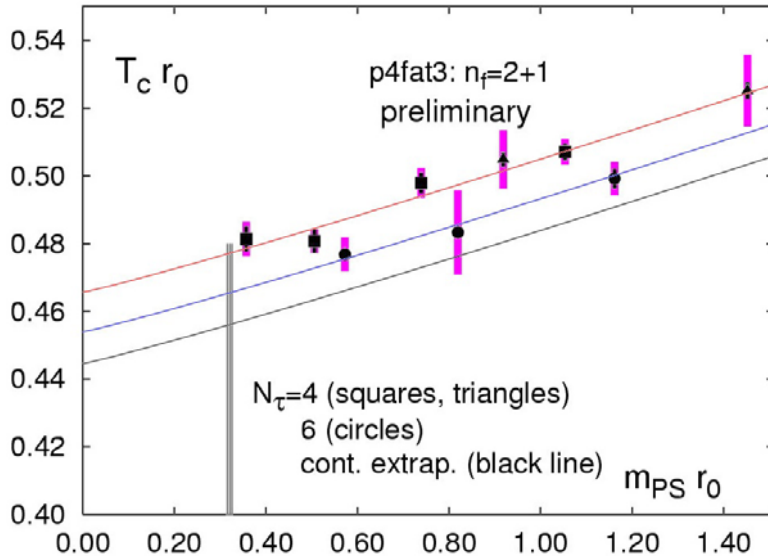
The fit result is used

- (1) correction for the diff. between β_c & simulation β at $T=0$
- (2) conversion of sys. + stat. error in β_c to error of r_0/a

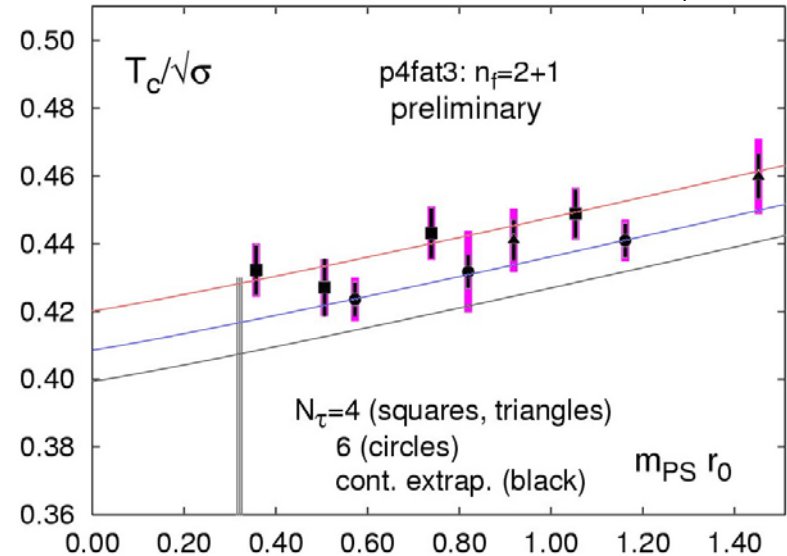
Critical temperature



ratio to Sommer scale:



ratio to string tension:



$$T_c r_0 = A(m_{ps} r_0)^d + B/N_\tau^2 + C \quad (d=1.08 \text{ from } O(4) \text{ scaling})$$

↳ $T_c r_0 = 0.456(7)_{-1}^{+3}$, $T_c / \sqrt{\sigma} = 0.408(7)_{-1}^{+3}$ at phys. point,
(fit form dependence $\rightarrow d=1,2$: upper, lower errors)

Finally we obtain $T_c=192(5)(4)\text{MeV}$ from $r_0=0.469(7)\text{fm}$

Spatial string tension



Important to check theoretical concepts (dim. reduction) at high T

Static quark “potential” from Spatial Wilson loops

$$V_z(r) = \ln W(r, z)/W(r, z + 1) \sim C - \frac{\alpha}{r_I} + \sigma_s r_I$$

$$\sqrt{\sigma_s(T)} = c g^2(T) T \quad (\text{free parameters: } c, \Lambda_\sigma)$$

$g^2(T)$ is given by the 2-loop RG equation

$$g^{-2}(T) = 2b_0 \ln \frac{T}{\Lambda_\sigma} + \frac{b_1}{b_0} \ln \left(2 \ln \frac{T}{\Lambda_\sigma} \right)$$

“ c ” should equal with 3-dim. string tension and should be flavor independent, if dim. reduction works

■ Numerical simulation

$$16^3 \times 4, \quad m_q = 0.1 m_s \text{ fixed}$$

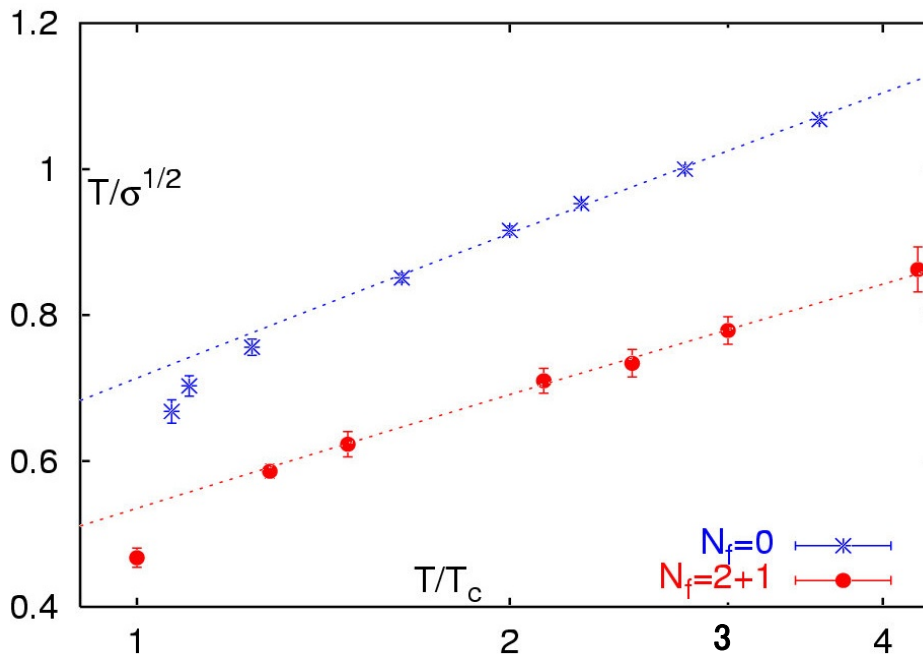
β	3.31	3.41	3.46	3.61	3.68	3.76	3.94
m_q	0.0065	0.0052	0.0040	0.00325	0.0026	0.002	0.001625

Spatial string tension



$$\sqrt{\sigma_s(T)} = cg^2(T)T \quad (\text{free parameters: } c, \Lambda_\sigma)$$

$$g^{-2}(T) = 2b_0 \ln \frac{T}{\Lambda_\sigma} + \frac{b_1}{b_0} \ln \left(2 \ln \frac{T}{\Lambda_\sigma} \right)$$



Quenched case:

G.Boyd et al.

Nucl.Phys.B469('96)419.

$$c = 0.566(13)$$

$$\frac{\Lambda_\sigma}{T_c} = 0.104(9)$$

2+1 flavor case:

$$c = 0.587(41)$$

$$\frac{\Lambda_\sigma}{T_c} = 0.114(27)$$

“c” is flavor independent within error

→ dim. reduction works well even for $T=2T_c$

Summary



$N_f=2+1$ simulation with realistic quark mass at $N_t=4, 6$

■ critical temperature

$$T_c r_0 = 0.456(7), \quad (T_c = 192(5)(4) \text{ MeV from } r_0 = 0.469)$$

– $T_c r_0$ is consistent with previous p4 result
difference in T_c mainly comes from physical value of r_0

– however, our value is about 10% larger than MILC result

MILC collab., Phys. Rev. D71('05) 034504.

– most systematic uncertainties are taken into account
remaining uncertainty is in continuum extrapolation

■ spatial string tension

dimensional reduction works well even for $T=2T_c$