

# Quarkonium correlators on the lattice

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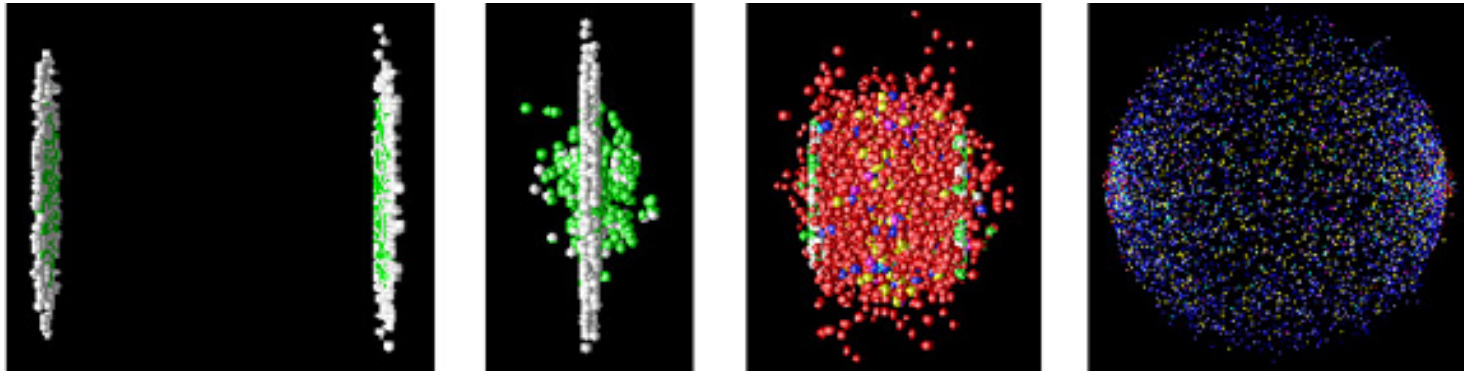


*Joint CATHIE-INT mini-program*  
*INT, Seattle, USA, June 16th 2009*

# Contents of this talk

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from the Phenix group web-site



- Introduction

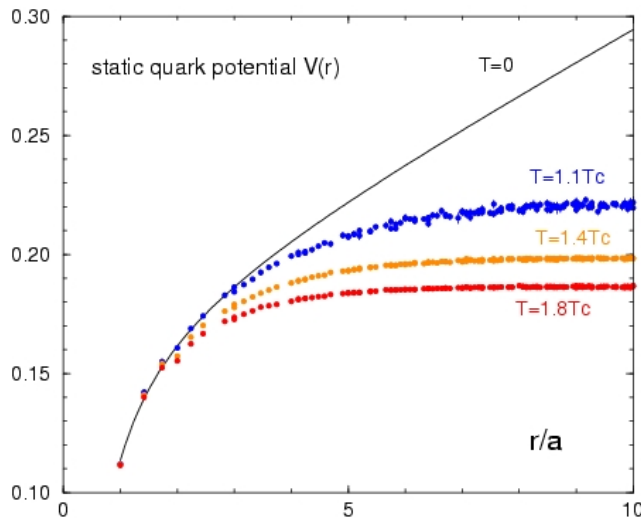
- Quark Gluon Plasma &  $J/\psi$  suppression
- Lattice studies on  $J/\psi$  suppression

- Our approach to study charmonium dissociation

- Charmonium wave functions at  $T > 0$

- Discussion & Summary

# $J/\psi$ suppression as a signal of QGP



Confined phase:

linear raising potential

→ bound state of  $c - \bar{c}$

De-confined phase:

Debye screening

→ scattering state of  $c - \bar{c}$

T.Hashimoto et al.('86), Matsui&Satz('86)

## Lattice QCD calculations:

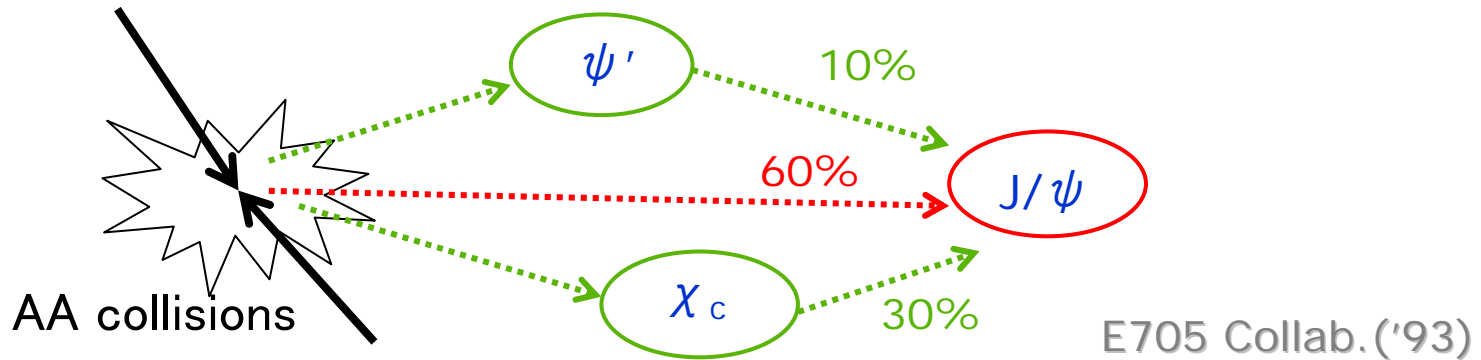
Spectral function by MEM: T.Umeda et al.('02), S.Datta et al.('04),  
Asakawa&Hatsuda('04), A.Jakovac et al.('07), G.Aatz et al.('06)

Wave func.: T.Umeda et al.('00)

B. C. dep.: H.Iida et al. ('06)

→ all calculations conclude that  $J/\psi$  survives till  $1.5T_c$  or higher

# Sequential $J/\psi$ suppression scenario



$J/\psi$ (1S)	: $J^{PC} = 1^{--}$	M=3097MeV	(Vector)
$\psi$ (2S)	: $J^{PC} = 1^{--}$	M=3686MeV	(Vector)
$\chi_{c0}$ (1P)	: $J^{PC} = 0^{++}$	M=3415MeV	(Scalar)
$\chi_{c1}$ (1P)	: $J^{PC} = 1^{++}$	M=3511MeV	(AxialVector)

PDG('06)

It is important to study dissociation temperatures for not only  $J/\psi$  but also  $\psi$  (2S),  $\chi_c$ 's

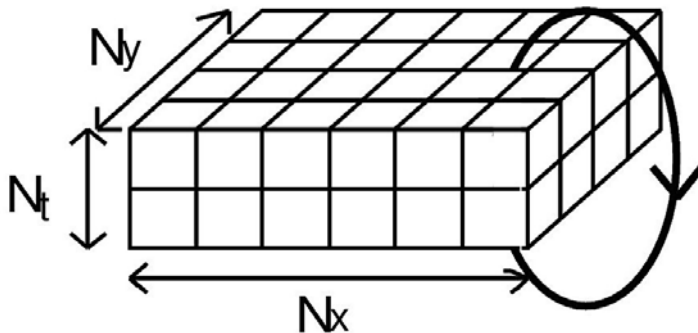
# Hot QCD on the lattice

Lattice QCD enables us to perform  
nonperturbative calculations of QCD

$$\langle X \rangle = \frac{1}{Z_{QCD}} \int Dq(x) D\bar{q}(x) DA_\mu(x) X(q, \bar{q}, A_\mu) e^{-S_{QCD}}$$

Path integral by Monte Carlo integration

QCD action on a lattice



Finite T Field Theory on the lattice

- 4dim. Euclidean lattice
- gauge field  $U_\mu(x) \rightarrow$  periodic B.C.
- quark field  $q(x) \rightarrow$  anti-periodic B.C.
- Temperature  $T=1/(N_t a)$

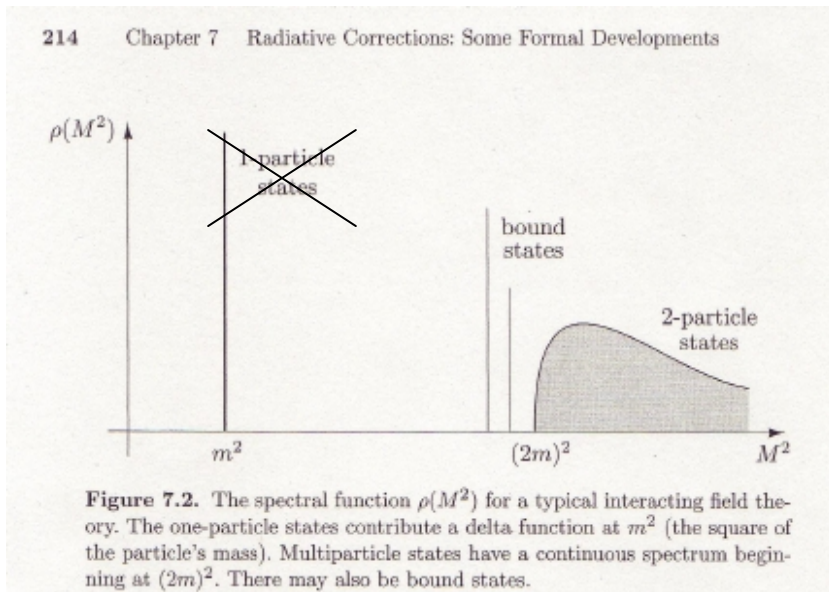
# Spectral function on the lattice

Thermal hadron (charmonium) correlation functions  $C_H(\tau, T)$

$$\begin{aligned}
 C_H(\tau, T) &= \sum_{\vec{r}} \langle J_H(\tau, \vec{r}) J_H^\dagger(0, \vec{0}) \rangle \\
 &= \int_0^\infty d\omega \sigma_H(\omega, T) \frac{\cosh(\omega(\tau - 1/2T))}{\sinh(\omega/2T)}
 \end{aligned}$$

Spectral function  $\sigma_H(\omega, T)$

- discrete spectra
  - bound states
  - charmonium states
- continuum spectra
  - 2-particle states
  - $c \bar{c}$  scattering states
  - melted charmonium

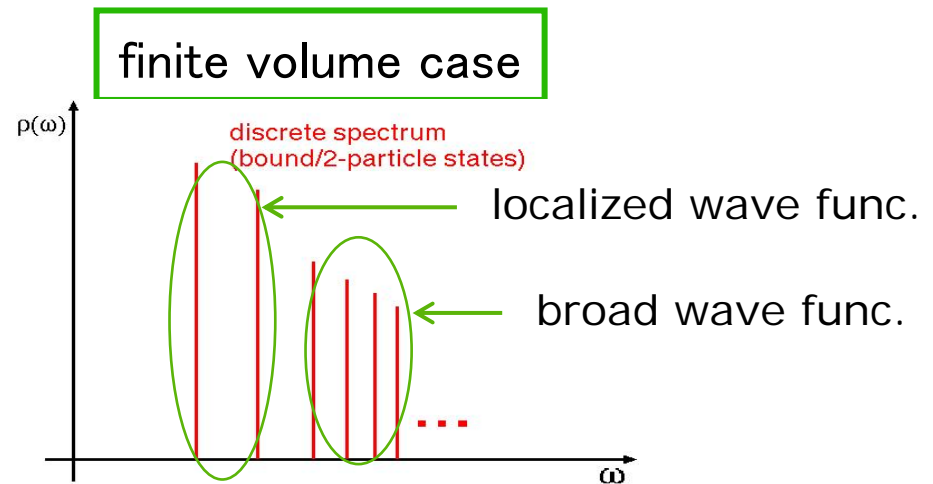
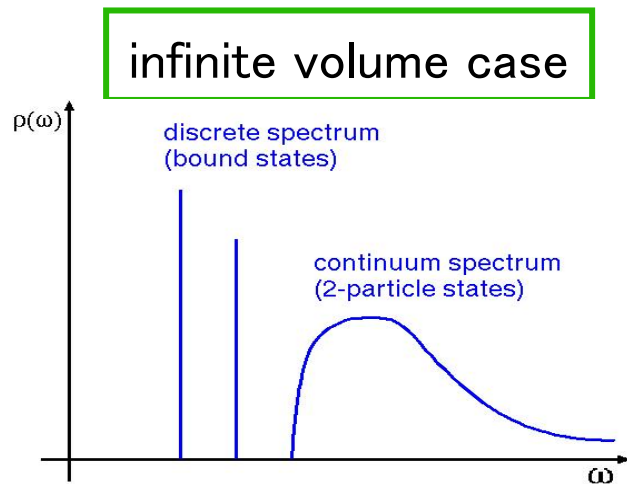


Michael E. Peskin, Perseus books (1995)

# Spectral functions in a finite volume

Momenta are discretized in finite ( $V=L^3$ ) volume

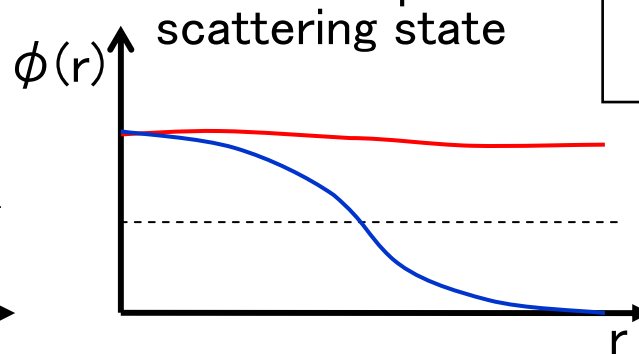
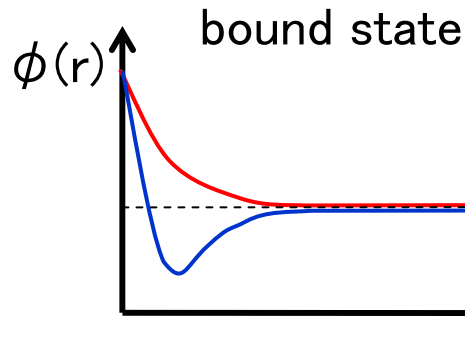
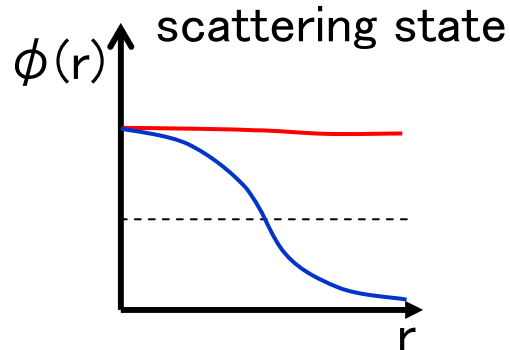
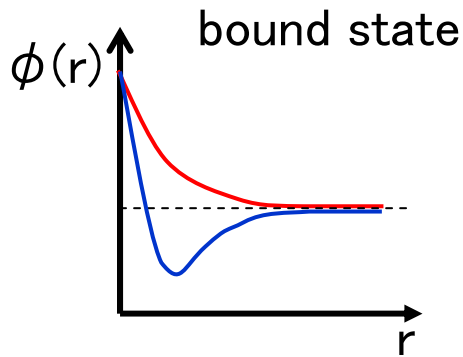
$p_i/a = 2n_i \pi / L$  ( $n_i=0, \pm 1, \pm 2, \dots$ ) for Periodic boundary condition



In a finite volume (e.g. Lattice simulations),  
discrete spectra does not always indicate bound states !

Shape of wave functions may be good signature  
to find out the charmonium melting.

# Bound state or scattering state ?



$\Phi(r)$ : wave function  
 $r$ : c - c distance

— lowest state  
 — next lowest state

examples for

- S-wave
- Periodic B.C

■ localized wave function  
 ■ small vol. dependence

■ broad wave function  
 ■ vol. dependence



# Wave functions at finite temperature

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Temp. dependence of ( Bethe-Salpeter ) “Wave function”

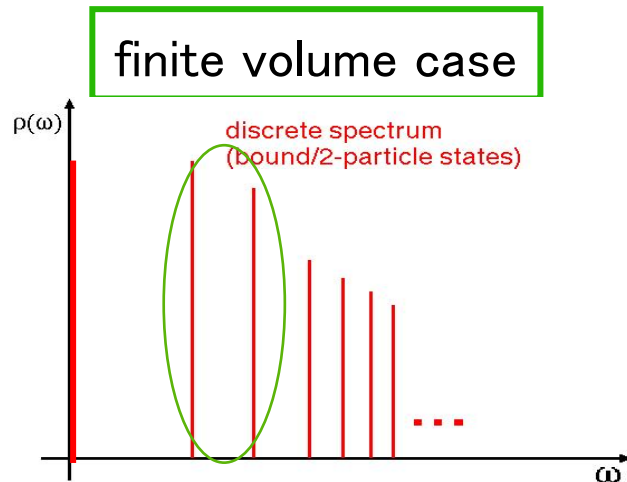
$$BS(\vec{r}, t) = \sum_{\vec{x}} \langle \bar{q}(\vec{x} + \vec{r}, t) \Gamma q(\vec{x}, t) \bar{q}(\vec{0}, 0) \Gamma q(\vec{0}, 0) \rangle$$
$$\Psi(|\vec{r}|, t) = BS(\vec{r}, t) / BS(\vec{r}_0, t)$$

$$\Gamma = \begin{cases} \gamma_5 & (\text{Ps}) \\ \gamma_i & (\text{Ve}) \quad (i = 1, 2, 3) \\ \sum_j (\vec{\partial}_j \gamma_j - \overleftarrow{\partial}_j \gamma_j) & (\text{Sc}) \\ \sum_{j,k} \epsilon_{ijk} (\vec{\partial}_j \gamma_k - \overleftarrow{\partial}_j \gamma_k) & (\text{Av}) \quad (i = 1, 2, 3) \end{cases}$$

Remarks on wave function of quark-antiquark

- gauge variant → Coulomb gauge fixing
- large or small components of quark/antiquark  
→ wave func. for large components

# Technique to calculate wave function at $T > 0$



It is difficult to extract higher states from lattice correlators (at  $T > 0$ ) even if we use MEM !!

It is important to investigate a few lowest states (at  $T > 0$ )

Constant mode can be separated by the Midpoint subtraction

*T. Umeda (2007)*

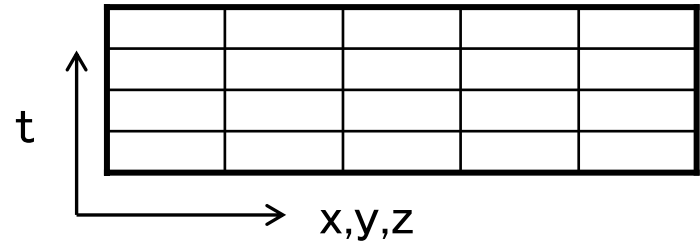
In order to study a few lowest states, the variational analysis is one of the most reliable methods !

$N \times N$  correlation matrix :  $C(t)$

$$C(t)\psi = \lambda(t, t_0)C(t_0)\psi \quad \lambda_i(t, t_0) = e^{-E_i(t-t_0)}$$

# Lattice setup

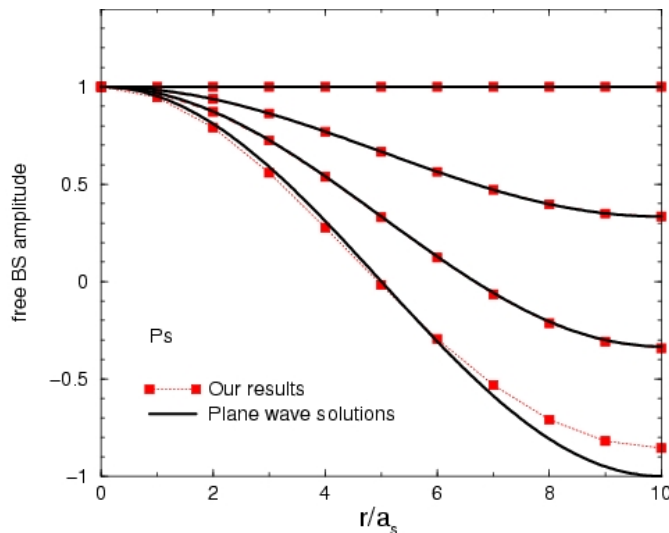
- Quenched approximation ( no dynamical quark effect )
- Anisotropic lattices
  - lattice spacing :  $a_s = 0.0970(5)$  fm
  - anisotropy :  $a_s/a_t = 4$
- $r_s=1$  to suppress doubler effects
- Variational analysis with 4 x 4 correlation matrix



$N_t$	32	26	20	16	12
$T/T_c$	0.88	1.08	1.40	1.75	2.33
# of conf.					
$V=16^3$	300	300	300	300	300
$V=20^3$	300	300	300	300	300
$V=32^3$	—	—	—	—	100

# Wave functions in free quark case

Test with free quarks (  $L_s/a=20$ ,  $ma=0.17$  )  
in case of S-wave channels

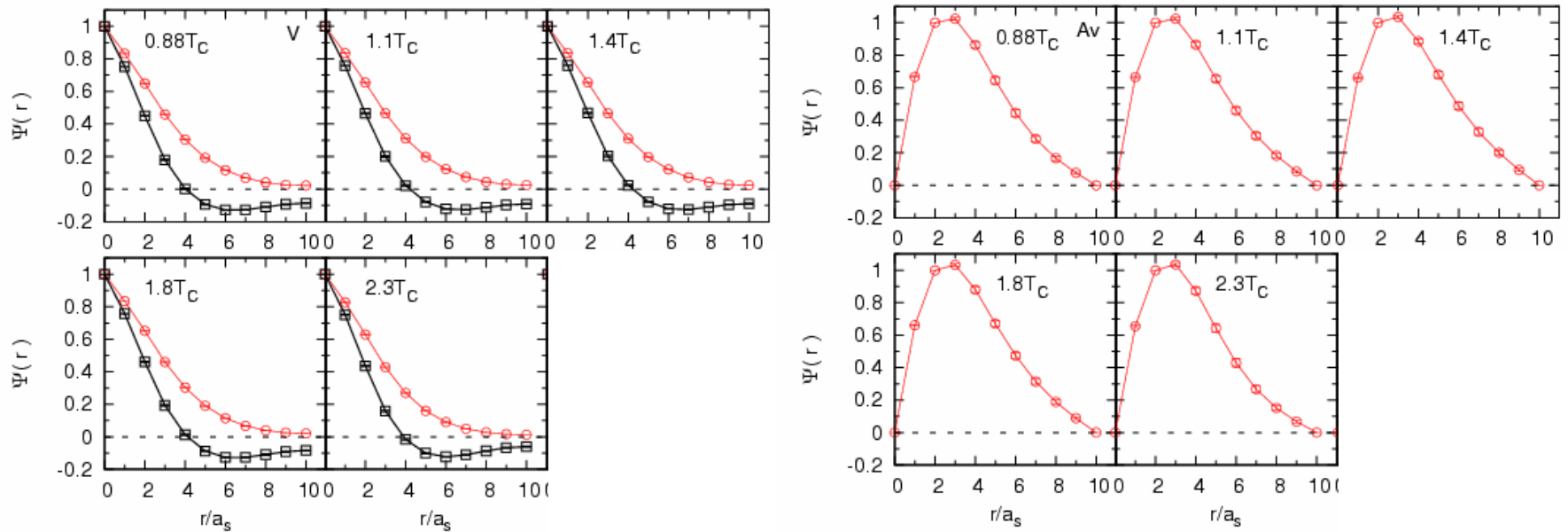


- Free quarks make trivial waves with an allowed momentum in a box

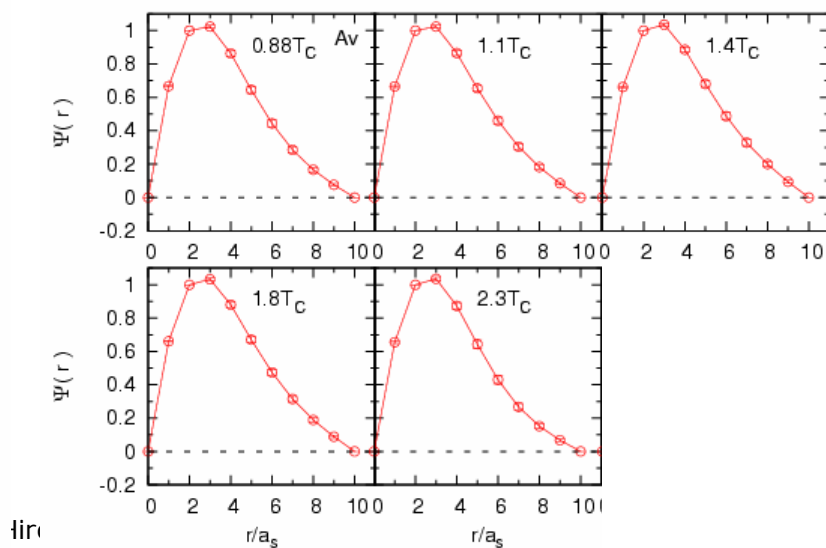
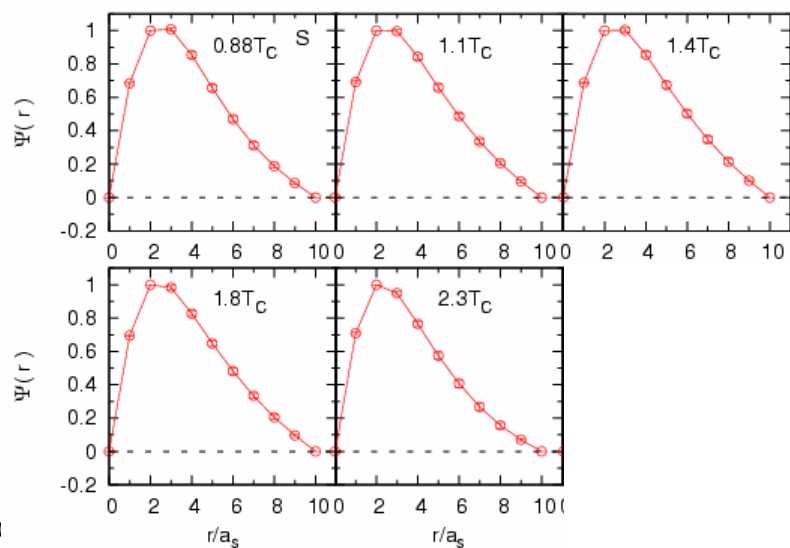
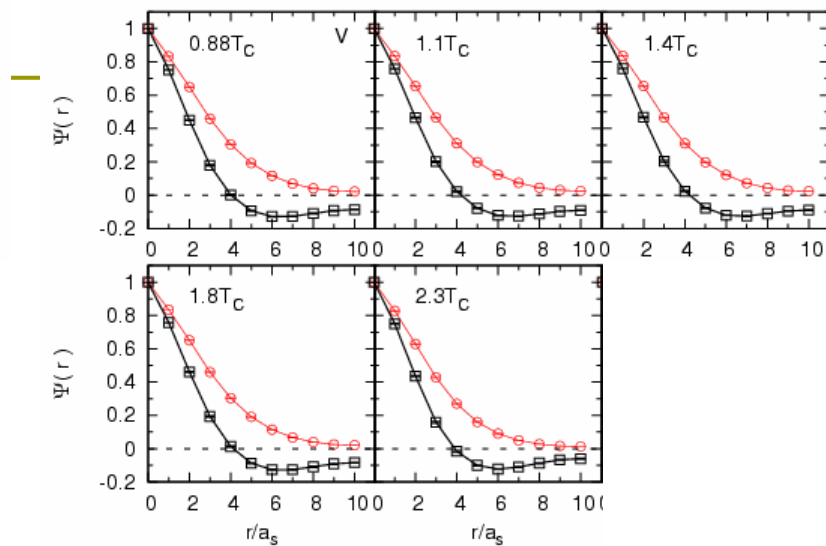
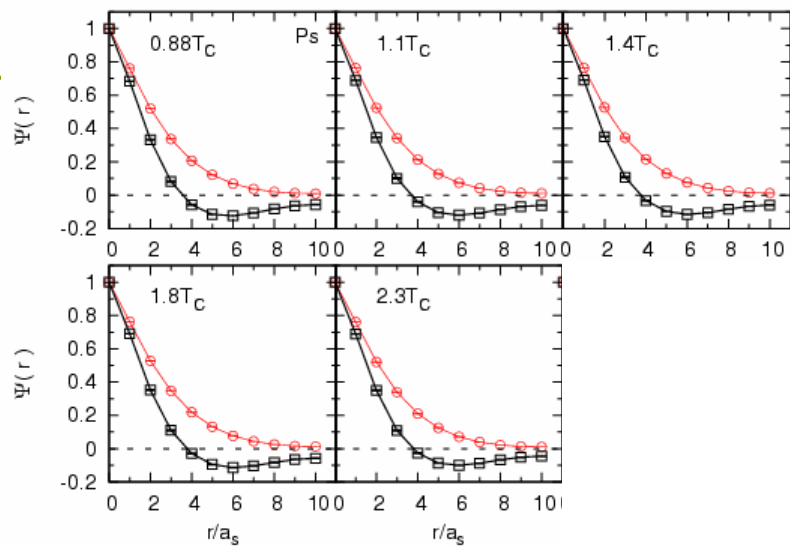
$$\Psi_k(|\vec{r}|, t) = \frac{\sum_{\vec{p}=\vec{k}} \cos(p_1 r_1) \cos(p_2 r_2) \cos(p_3 r_3)}{\sum_{\vec{p}=\vec{k}} 1}$$

- The wave function is constructed with eigen functions of 6 x 6 correlators
- 6 types of Gaussian smeared operators  $\phi(x) = \exp(-A|x|^2)$ ,  
 $A = 0.02, 0.05, 0.1, 0.15, 0.2, 0.25$
- Our method well reproduces the known result ( ! )

# Charmonium wave functions at finite temperatures

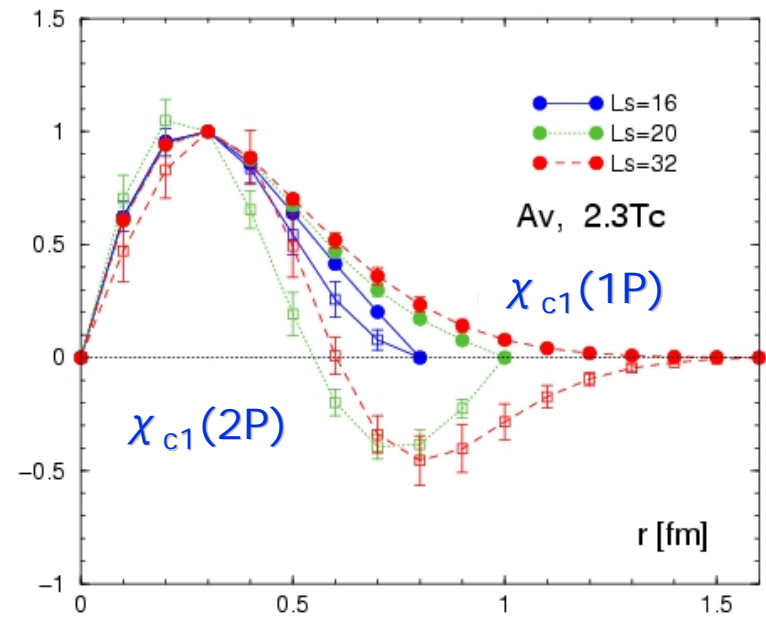
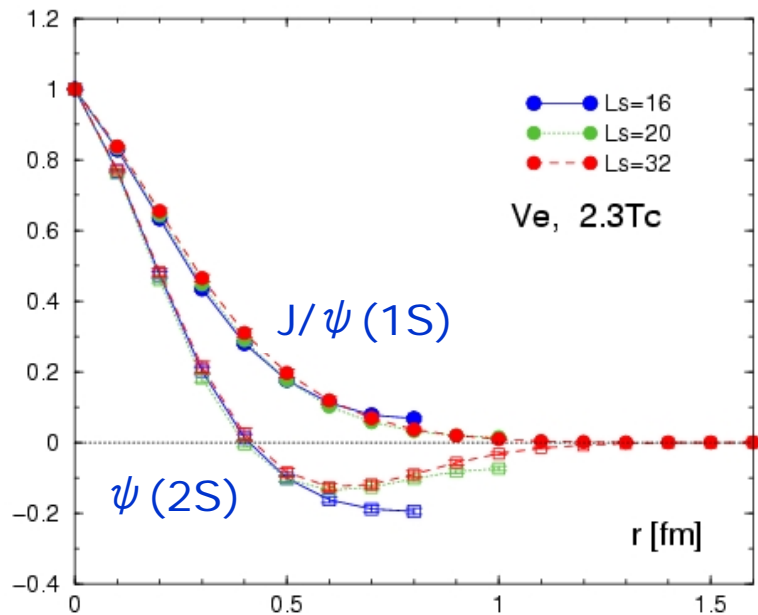


- Small temperature dependence in each channels
- Clear signals of bound states even at  $T=2.3T_C$  (!)
- $(2\text{fm})^3$  may be small for P-wave states.



lit

# Volume dependence at $T=2.3T_c$



- Clear signals of bound states even at  $T=2.3T_c$  (!)
- Large volume is necessary for P-wave states.

# Discussion

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We found localized wave functions  
up to  $2.3T_c$  for S- & P- wave channels.

(1) Does variational analysis does work well ?

wave function for lowest/next-lowest state

+ contributions from higher states ← contaminations

our results suggest

there are no/small broad wave functions  
even in the higher states (!)

(2) Unbound state with localized wave function ?

anyway,

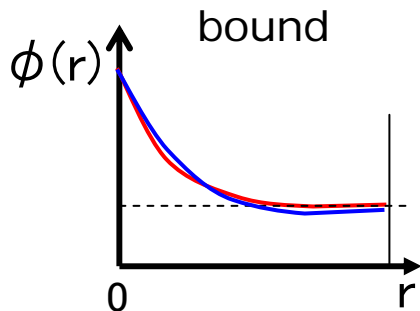
tight wave function is incompatible  
with the  $J/\psi$  suppression (!)

(3) Small changes in wave functions

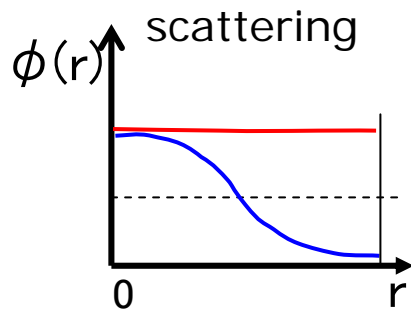
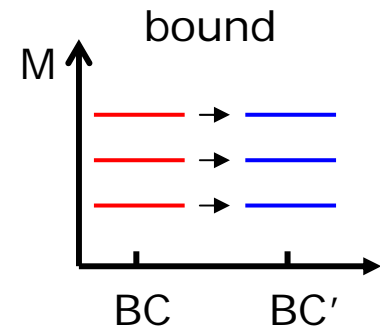
contradict some potential model results with  $T_{\text{dis}}(J/\psi) > T_c$  (!)



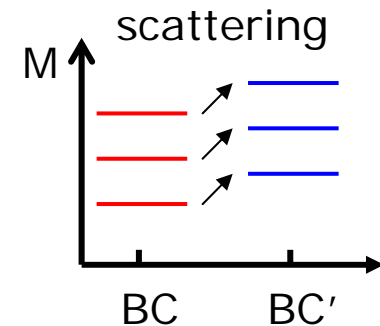
# Boundary condition dependence



The wave functions are localized,  
their energies are insensitive to B.C.



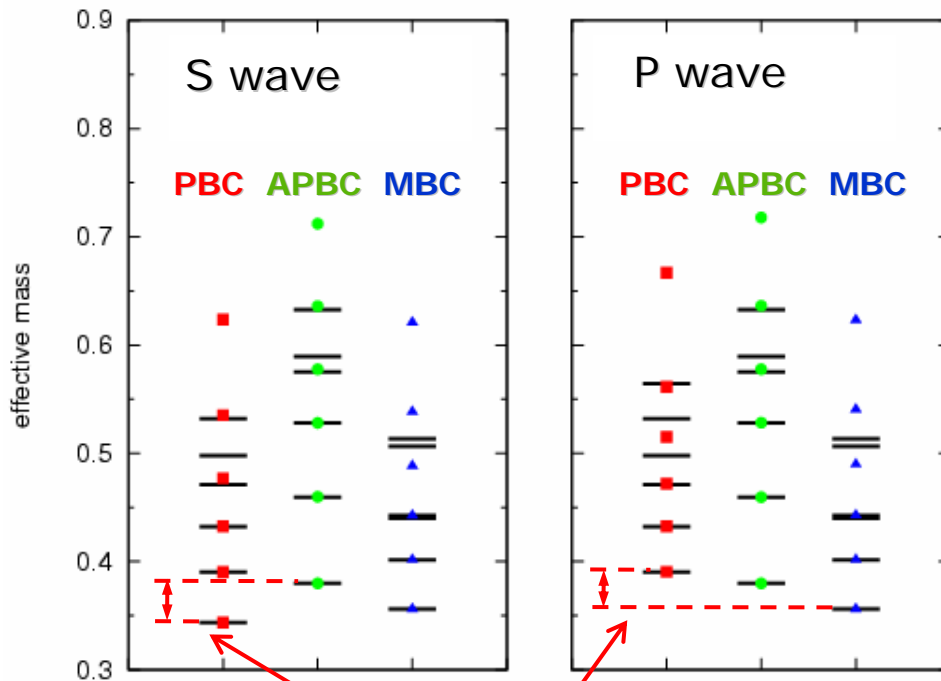
The momenta depends on BC,  
the scattering state energies  
are sensitive to B.C.



*The idea has been originally applied for the charmonium study  
in H. Iida et al., Phys. Rev. D74 (2006) 074502.*

# Variational analysis in free quark case

Test with free quarks (  $L_s/a=20$ ,  $ma=0.17$  )



Mass diff. between the lowest masses in each BC

$$q(x_i + L_i) = b_i q(x_i)$$

$b_i = 1$  : periodic

$b_i = -1$  : anti-periodic

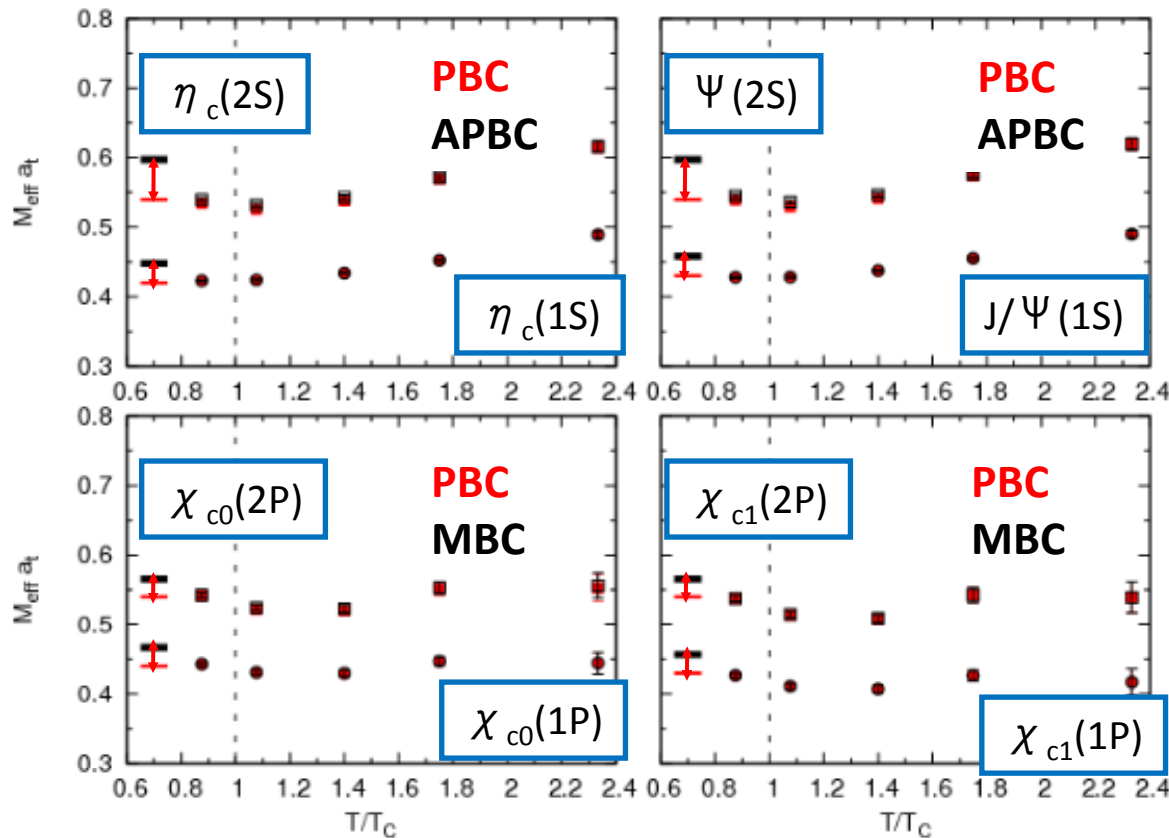
PBC :  $b=(1, 1, 1)$

APBC :  $b=(-1, -1, -1)$

MBC :  $b=(-1, 1, 1)$

an expected diff.  
in  $V=(2\text{fm})^3$   
(free quark case)  
 $\sim 200\text{MeV}$

# Temperature dependence of charmonium spectra



$$q(x_i + L_i) = b_i q(x_i)$$

$b_i = 1$  : periodic  
 $b_i = -1$  : anti-periodic

PBC :  $b = (1, 1, 1)$   
 APBC :  $b = (-1, -1, -1)$   
 MBC :  $b = (-1, 1, 1)$

an expected gap  
 in  $V = (2\text{fm})^3$   
 (free quark case)  
 $\sim 200\text{MeV}$

- No significant differences in the different B.C.
- Analysis is difficult at higher temperature ( $2T_c \sim$ )

# Summary and future plan

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We investigated  $T_{\text{dis}}$  of charmonia from Lattice QCD without Bayesian (MEM) analysis using...

- Bethe-Salpeter “wave function”
- Volume dependence of the “wave function”
- Boundary condition dependence

No evidence for unbound  $c\bar{c}$  quarks up to  $T = 2.3 T_c$

→ The result may affect the scenario of  $J/\psi$  suppression.

## Future plan

- Possible scenarios for the experimental  $J/\psi$  suppression
- Higher Temp. calculations (  $T/T_c=3\sim 5$  )
- Full QCD calculations (  $N_f=2+1$  Wilson is now in progress )