

# 有限温度格子QCDの 新しいアプローチの可能性

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for WHOT-QCD Collaboration

This talk is (partly) based on Phys. Rev. D 79, 051501(R) (2009)

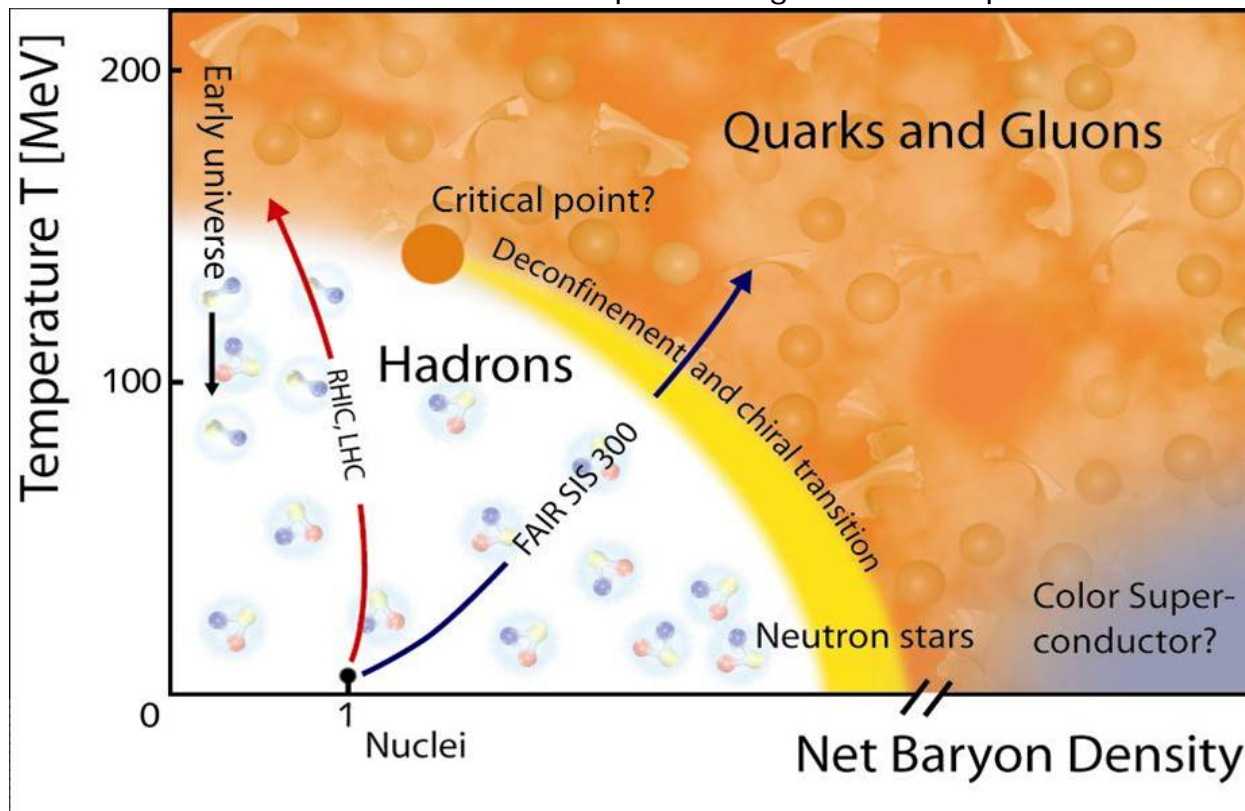
*S. Ejiri, S. Aoki, T. Hatsuda, N. Ishii, K. Kanaya,  
Y. Maezawa, H. Ohno, T.U. (WHOT-QCD Collaboration)*

*GCOE-PD seminar, Kyoto, Japan, 18 Mar. 2009*

# Study of Quark-Gluon Plasma

QCD phase diagram in (Temperature, density)

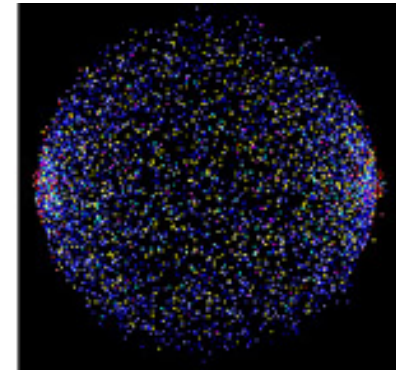
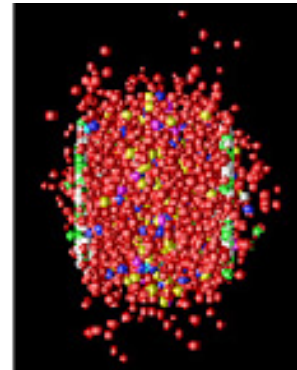
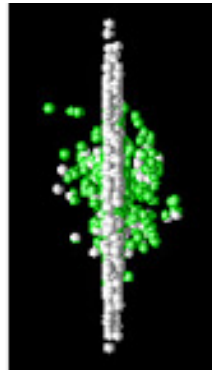
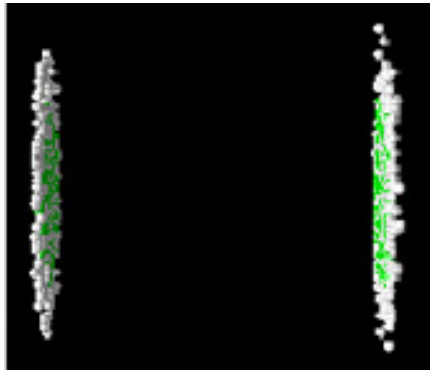
<http://www.gsi.de/fair/experiments/>



# Heavy Ion Collision experiments



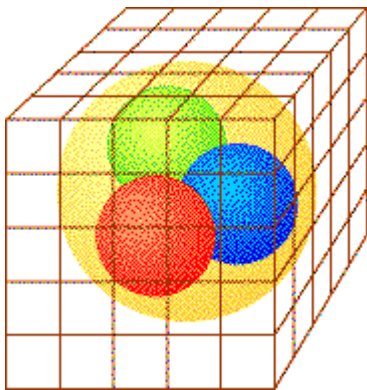
- **SPS** : CERN ( – 2005)  
Super Proton Synchrotron
- **RHIC** : BNL (2000 – )  
Relativistic Heavy Ion Collider
- **LHC** : CERN (2009 - )  
Large Hadron Collider



# Lattice QCD simulations

## Lattice QCD

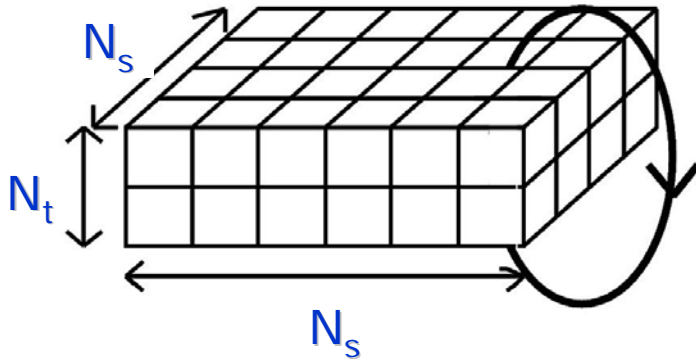
- First principle (nonperturbative) calculation of QCD
- QCD action is defined on the lattice (discretized space-time)
- Path integral is carried out by Monte Carlo Integration



## Contents of this talk

- Introduction
- Problems in the conventional approach
- A new approach for QCD Thermodynamics on lattices
- The EOS calculation by **"T-integration method"**
- Summary

# Hot QCD on the lattice



## Finite T Field Theory on the lattice

- 4dim. Euclidean lattice ( $N_s^3 \times N_t$ )
- gauge field  $U_\mu(x) \rightarrow$  periodic B.C.
- quark field  $q(x) \rightarrow$  anti-periodic B.C.
- Temperature  $T = 1/(N_t a)$

Input parameters :  $\beta (=6/g^2)$  (lattice gauge coupling)  
 (Nf=2+1 QCD)  $am_{ud}$  (light (up & down) quark masses)  
 $am_s$  (strange quark mass)  
 $N_t$  (temperature)

(\* ) lattice spacing "a" is not an input parameter

$$a = a(\beta, am_{ud}, am_s)$$

$$\text{e.g. } (am_\rho)_{\text{lat}} / (0.77 \text{ GeV}) = a [\text{GeV}^{-1}]$$

Temperature  $T = 1/(N_t a)$  is varied by  $a$  at fixed  $N_t$

# Fermions on the lattice

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## Lattice QCD

QCD action is defined on the lattice

### Fermion doubling problem

- naive discretization causes  $2^4$  doublers
- Nielsen-Ninomiya's No-go theorem
  - Doublers appear unless chiral symmetry is broken

- Staggered (KS) fermion → Low cost  
16 doublers = 4 spinors x 4 flavors  
Fourth root trick : still debated
- Wilson fermion → Moderate cost  
adds the Wilson term to kill extra  $2^4-1$  doublers
- Domain Wall fermion → High cost
- Overlap fermion → High cost
- ...

# Problems in QCD Thermo. with KS fermions

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Many QCD thermo. calc. were done with KS fermions.

- Phase diagram

  - $N_f=2$  massless QCD  $\rightarrow$   $O(4)$  critical exponents

  - KS fermion does not exhibit expected  $O(4)$  scaling

  - (Wilson fermion shows  $O(4)$ , but at rather heavy masses)

- Transition temperature (crossover transition in KS studies)

  - Equation of State (  $p/T^4$ ,  $e/T^4$ ,  $s/T^4$ , ... )

  - KS results are not consistent with each other

$N_f=2$ ,  $2+1$  is not 4 !!!

# Integral method to calculate pressure $p/T^4$

$$p = \frac{T}{V} \ln Z \quad \text{for large volume system}$$


Lattice QCD can not directly calculate the partition function  $\ln Z$

however its derivative is possible  $\frac{\partial}{\partial \beta} \ln Z = - \left\langle \frac{\partial S_{QCD}}{\partial \beta} \right\rangle$

One can obtain  $p$  as the integral of derivative of  $p$

high temp. 

$$\left. \frac{p}{T^4} \right|_{\beta_0}^{\beta} = \frac{1}{VT^3} \int_{\beta_0}^{\beta} d\beta' \frac{\partial}{\partial \beta'} \ln Z$$

low temp.  
with  $p \approx 0$  

$$= -N_t^4 \int_{\beta_0}^{\beta} d\beta' \frac{1}{N_s^3 N_t} \left( \left\langle \frac{\partial S_{QCD}}{\partial \beta} \right\rangle_{T>0} - \left\langle \frac{\partial S_{QCD}}{\partial \beta} \right\rangle_{T=0} \right)$$

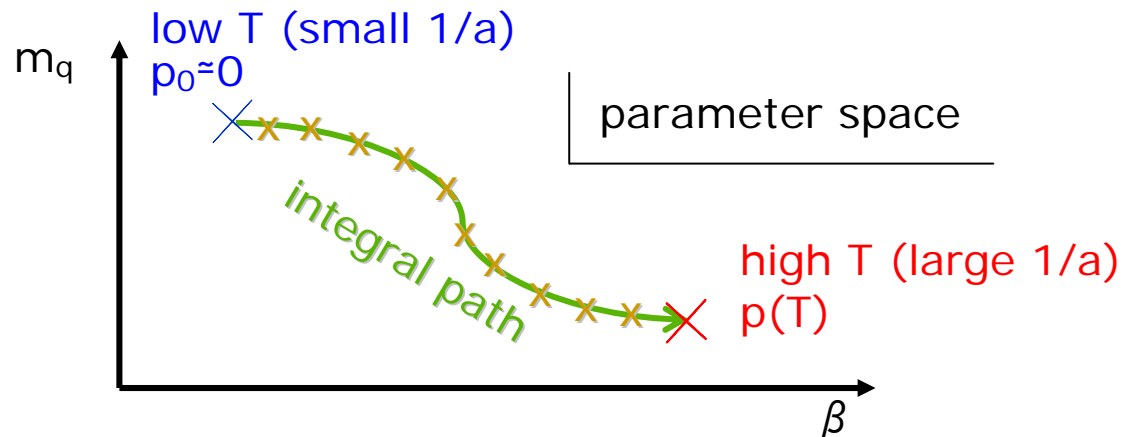
T=0 subtraction



# Line of constant physics (LCP)

In case of  $N_f=2+1$  QCD

there are three (bare) parameters:  $\beta$ ,  $(am_{ud})$  and  $(am_s)$



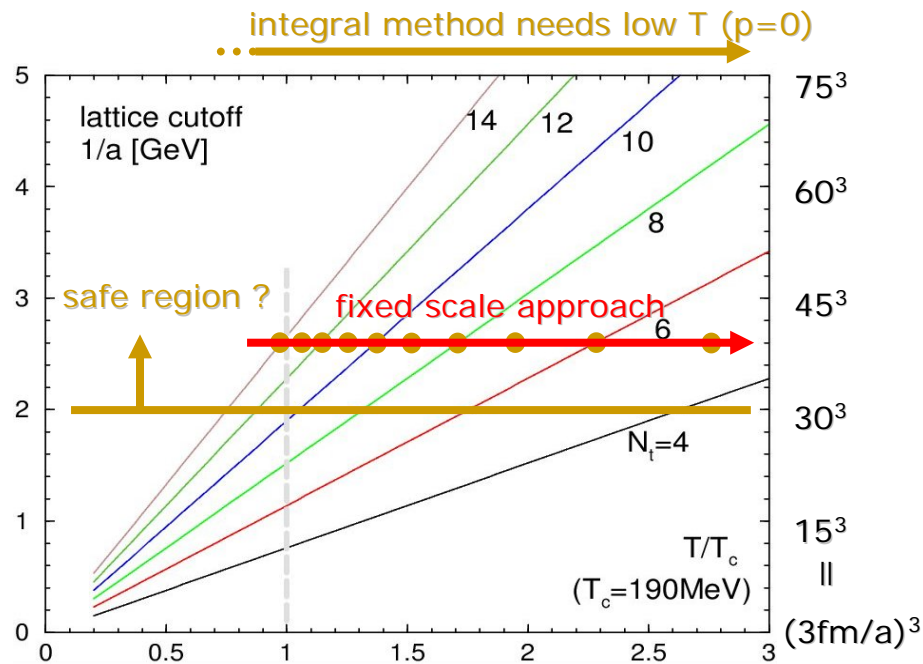
**Line of Constant Physics (LCP)** defined at  $T=0$

QCD Thermodynamics requires huge computational cost !!

Most group adopts KS fermion to study the QCD Thermodynamics.

# Fixed scale approach to study QCD thermodynamics

Temperature  $T=1/(N_t a)$  is varied by  $N_t$  at fixed  $a(\beta, m_{ud}, m_s)$



## Advantages

- LCP is trivially exact
- $T=0$  subtraction is done with a common  $T=0$  sim. ( $T=0$  high. stat. spectrum)
- easy to keep large  $1/a$  at whole  $T$  region
- easy to study  $T$  effect without  $V$ ,  $1/a$  effects

## Disadvantages

- $T$  resolution by integer  $N_t$
- program for odd  $N_t$
- $(1/a)/T = \text{const.}$  is not suited for high  $T$  limit study

# T-integration method to calculate the EOS

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We propose a new method (“**T-integration method**”)  
to calculate the EOS at fixed scales


*T.Umeda et al. (WHOT-QCD) Phys. Rev. D 79, 051501(R) (2009)*

Our method is based on **the trace anomaly** (interaction measure),

$$\frac{\epsilon - 3p}{T^4} = \left( \frac{N_t^3}{N_s^3} \right) a \frac{d\beta}{da} \left\langle \frac{dS}{d\beta} \right\rangle_{sub}$$

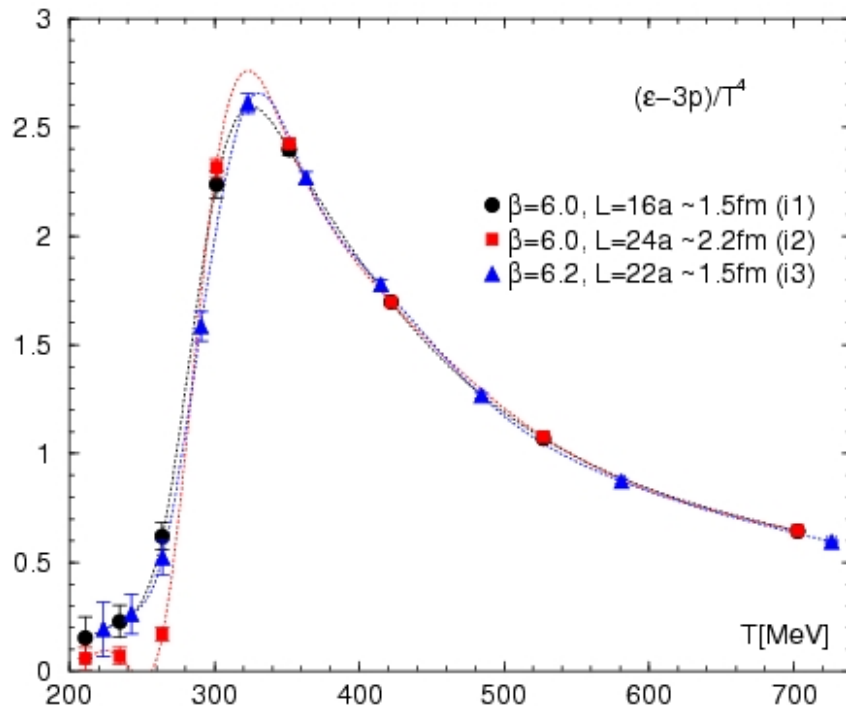
and **the thermodynamic relation**.

$$\frac{\epsilon - 3p}{T^4} = T \frac{\partial(p/T^4)}{\partial T}$$

  $\frac{p}{T^4} = \int_0^T dT' \frac{\epsilon - 3p}{T'^5}$

# Trace anomaly $(\epsilon - 3p)/T^4$ in SU(3) gauge theory

We present results from SU(3) gauge theory as a test of our method



dotted lines : cubic spline

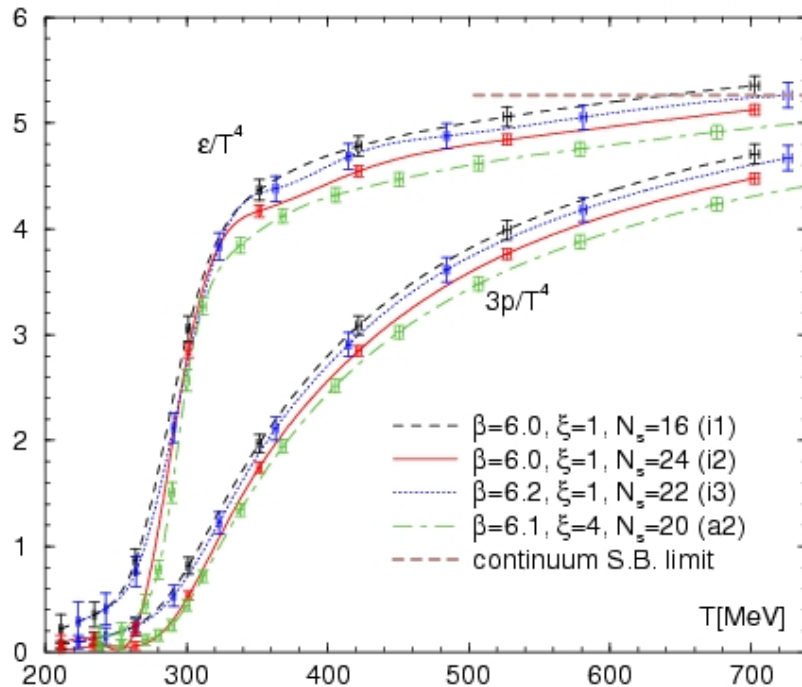
- (1)  $\beta = 6.0, 1/a = 2.1\text{GeV}, V = (1.5\text{fm})^3$
- (2)  $\beta = 6.0, 1/a = 2.1\text{GeV}, V = (2.2\text{fm})^3$
- (3)  $\beta = 6.2, 1/a = 2.5\text{GeV}, V = (1.5\text{fm})^3$

beta function : G.Boyd et al. ('96)  
lattice scale  $r_0$  : R.Edwards et al. ('98)

$$\frac{\epsilon - 3p}{T^4} = \left( \frac{N_t^3}{N_s^3} \right) a \frac{\partial \beta}{\partial a} \left\langle \frac{\partial S_g}{\partial \beta} \right\rangle_{sub}$$

# Trace anomaly $(e - 3p)/T^4$ in SU(3) gauge theory

We present results from SU(3) gauge theory as a test of our method



- (1)  $\beta = 6.0, 1/a=2.1\text{GeV}, V=(1.5\text{fm})^3$
- (2)  $\beta = 6.0, 1/a=2.1\text{GeV}, V=(2.2\text{fm})^3$
- (3)  $\beta = 6.2, 1/a=2.5\text{GeV}, V=(1.5\text{fm})^3$

beta function : G.Boyd et al. ('96)  
lattice scale  $r_0$  : R.Edwards et al. ('98)

$$\text{Integration } \left( \frac{p}{T^4} = \int_0^T dT' \frac{\epsilon - 3p}{T'^5} \right)$$

is performed with the cubic spline of  $(e-3p)/T^4$

Our fixed scale approach with "T-integration method" works well !!

# Summary and future plans

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- A new approach to study the QCD Thermodynamics is proposed.
  - "T-integral method" to calculate the EOS works well.
  - There are many advantages in the approach.
- We have already generated  $T>0$  configurations using CP-PACS/JLQCD parameter ( $N_f=2+1$  Clover+RG,  $1/a=3\text{GeV}$ , pion mass  $\sim 500\text{MeV}$ )
- Our final goal is to study thermodynamics on the physical point (pion mass  $\sim 140\text{MeV}$ ) with  $N_f=2+1$  Wilson quarks (PACS-CS) or exact chiral symmetry with  $N_f=2+1$  Overlap quarks (JLQCD)
- We are looking for new ideas to study QGP physics in our approach. ( density correlations, J/psi suppression, finite density...)

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# Backup slides

# Recent lattice calculations of EOS

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Hot-QCD Collab. (2007~)	RBC-Bielefeld:	$N_t=4,6,8$	Staggered (p4) quark pion mass $\sim 220\text{MeV}$ , $N_f=2+1$ <i>Phys. Rev. D77 (2008) 014511</i>
	MILC:	$N_t=4,6,8$	Staggered (Asqtad) quark pion mass $\sim 220\text{MeV}$ , $N_f=2+1$ <i>Phys. Rev. D75 (2007) 094505</i>
	Wuppertal:	$N_t=4,6$	Staggered (stout) quark pion mass $\sim 140\text{MeV}$ , $N_f=2+1$ <i>JHEP 0601 (2006) 089</i>
	CP-PACS:	$N_t=4,6$	Wilson (MFI Clover) quark pion mass $\sim 500\text{MeV}$ , $N_f=2$ <i>Phys. Rev. D64 (2001) 074510</i>



# Introduction

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## Physics in Lattice QCD at finite temperature

- Phase diagram in  $(T, \mu, m_{ud}, m_s)$
- Transition temperature
- Equation of state  $(e, p, s, \dots)$
- Excitation spectrum
- Transport coefficients (shear/bulk viscosity)
- Finite chemical potential
- etc...

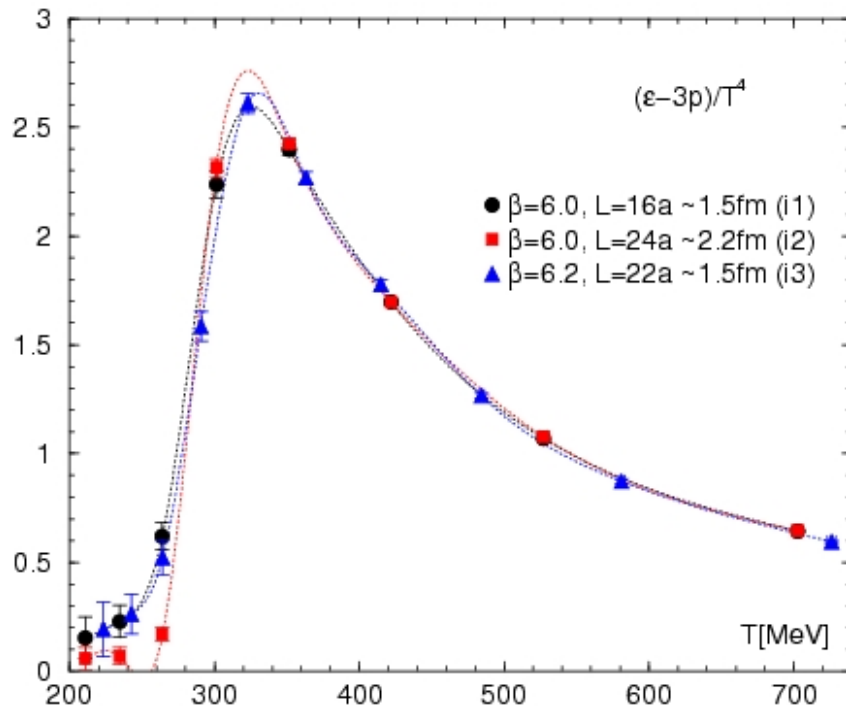


These are important to study

- Quark Gluon Plasma in Heavy Ion Collision exp.
- Early universe
- Neutron star
- etc...

# Trace anomaly $(\epsilon - 3p)/T^4$ on isotropic lattices

$$\frac{\epsilon - 3p}{T^4} = \left( \frac{N_t^3}{N_s^3} \right) a \frac{\partial \beta}{\partial a} \left\langle \frac{\partial S_g}{\partial \beta} \right\rangle_{sub}$$



dotted lines : cubic spline

- (1)  $\beta = 6.0, 1/a = 2.1 \text{ GeV}, V = (1.5 \text{ fm})^3$
- (2)  $\beta = 6.0, 1/a = 2.1 \text{ GeV}, V = (2.2 \text{ fm})^3$
- (3)  $\beta = 6.2, 1/a = 2.5 \text{ GeV}, V = (1.5 \text{ fm})^3$

beta function : G.Boyd et al. ('96)  
lattice scale  $r_0$  : R.Edwards et al. ('98)

- Excellent agreement between (1) and (3)  
→ scale violation is small  
 $1/a = 2 \text{ GeV}$  is good
- Finite volume effect appears below & near  $T_c$   
→ volume size is important  
 $V = (2 \text{ fm})^3$  is necessary.

# Simulation parameters (isotropic lattices)

We present results from SU(3) gauge theory as a test of our method

- plaquette gauge action on  $N_s^3 \times N_t$  lattices
- Jackknife analysis with appropriate bin-size

To study scale- & volume-dependence,  
we prepare 3-type of lattices.

(1)  $\beta = 6.0, V = (16a)^3$   
 $1/a = 2.1 \text{ GeV}$

$\beta$	$N_s$	$N_t$	T[MeV]	conf.
6.0	16	16	$\sim 0$	350k
6.0	16	10	210	350k
6.0	16	9	230	250k
6.0	16	8	260	200k
6.0	16	7	300	100k
6.0	16	6	350	50k
6.0	16	5	420	50k
6.0	16	4	530	50k
6.0	16	3	700	50k

(2)  $\beta = 6.0, V = (24a)^3$   
 $1/a = 2.1 \text{ GeV}$

$\beta$	$N_s$	$N_t$	T[MeV]	conf.
6.0	24	16	$\sim 0$	150k
6.0	24	10	210	250k
6.0	24	9	230	200k
6.0	24	8	260	150k
6.0	24	7	300	100k
6.0	24	6	350	50k
6.0	24	5	420	50k
6.0	24	4	530	50k
6.0	24	3	700	50k

(3)  $\beta = 6.2, V = (22a)^3$   
 $1/a = 2.5 \text{ GeV}$

$\beta$	$N_s$	$N_t$	T[MeV]	conf.
6.2	22	22	$\sim 0$	250k
6.2	22	13	220	350k
6.2	22	12	240	350k
6.2	22	11	270	350k
6.2	22	10	290	250k
6.2	22	9	320	200k
6.2	22	8	360	200k
6.2	22	7	420	100k
6.2	22	6	490	100k
6.2	22	5	580	50k
6.2	22	4	730	50k