

The hyperfine splitting of charmonium in lattice QCD

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Charmonium hyperfine splitting on the lattice

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Introduction (1)

Charmonium spectrum on lattice QCD

lattice cutoff $1/a \gg m_{\text{charm}}$ is necessary

→ large computational cost !

- NRQCD
 - HQET
 - Fermilab
- } Recent progress will be presented
in the next talk !

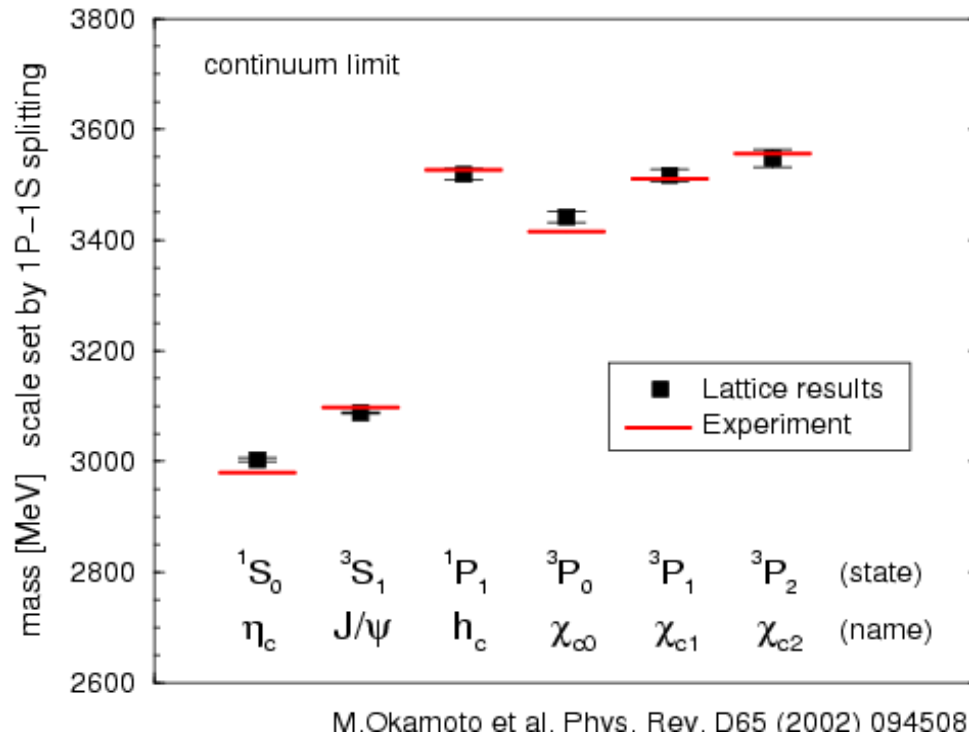
- **Anisotropic lattice** → next page

where the condition may be relaxed

to $1/a_t \gg m_{\text{charm}}$, $1/a_s \sim m_{\text{charm}}$?

Introduction (2)

M.Okamoto et al. (CP-PACS), Phys. Rev. D65 (2002)094508



- quenched anisotropic lattice
- tree-level clover quark
- continuum extrapolation
 - with $1/a_s = 1.0 \sim 2.9$ GeV
 - ▶ reliable extrapolation
 - needs $1/a_s \gg m_{\text{charm}}$
 - ▶ large dep. on Dirac op.

Good agreement with the experimental value

except for the hyperfine splitting

A study by QCD-TARO Collaboration

QCD-TARO Collaboration, JHEP08 (2003) 022

- Quenched QCD (without dynamical quarks)
- Isotropic lattices with large cutoff $1/a \gg m_{\text{charm}}$
- Nonperturbative improved clover quark
M.Lüscher et al., Nucl.Phys.B491(1997)323
- Continuum extrapolation

Lattice setup

Gauge field : plaquette gauge

Quark field : nonperturbative improved clover quark

(tree level clover quark, standard Wilson)

beta	$L^3 \times T$	a(fm)	La(fm)	Csw	#conf
6.0	$18^3 \times 48$	0.0931	1.68	1.769	190
6.2	$24^3 \times 72$	0.0677	1.62	1.614	90
6.4	$32^3 \times 96$	0.0513	1.64	1.526	60
6.6	$32^3 \times 96$	0.0397	1.27	1.467	130

scale set by r_0 , quark mass set by $m(J/\psi)$

finite volume effect is negligible at $La > 1.1 \text{ fm}$

Scale determination

- Lattice cutoff

Sommer scale $r_0 \simeq 0.5 \text{ fm}$

$r_0 \sim 10 \%$ uncertainty (due to dynamical quark effects)

Hadron mass (e.g. $m(1P) - m(1S)$)

for example CP-PACS results $\sim 16 \%$ uncertainty

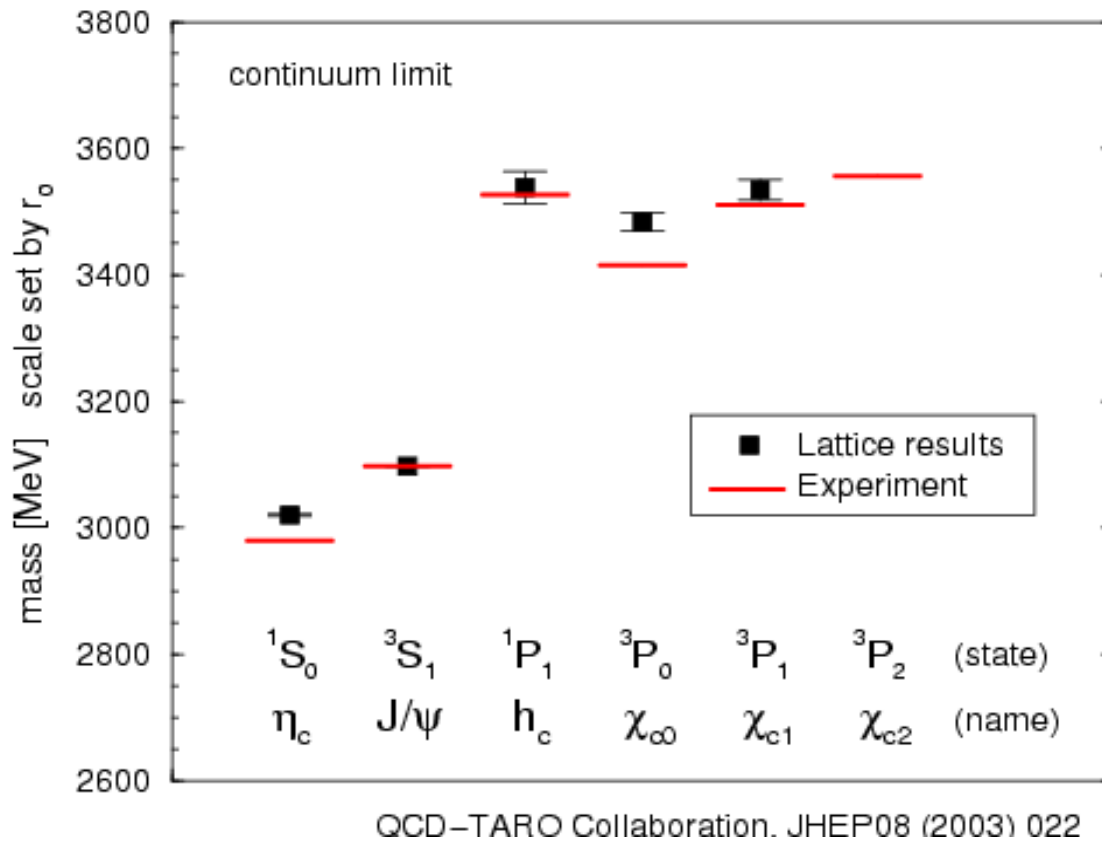
▶ discrepancy caused by
scaling violation & dynamical quark effects

- Quark mass

matching with a physical charmonium mass

$$\left. \begin{array}{l} m(J/\psi) = 3.097 \text{ GeV} \\ m(\eta_c) = 2.980 \text{ GeV} \end{array} \right\} \sim 6 \% \text{ uncertainty}$$

Charmonium spectrum

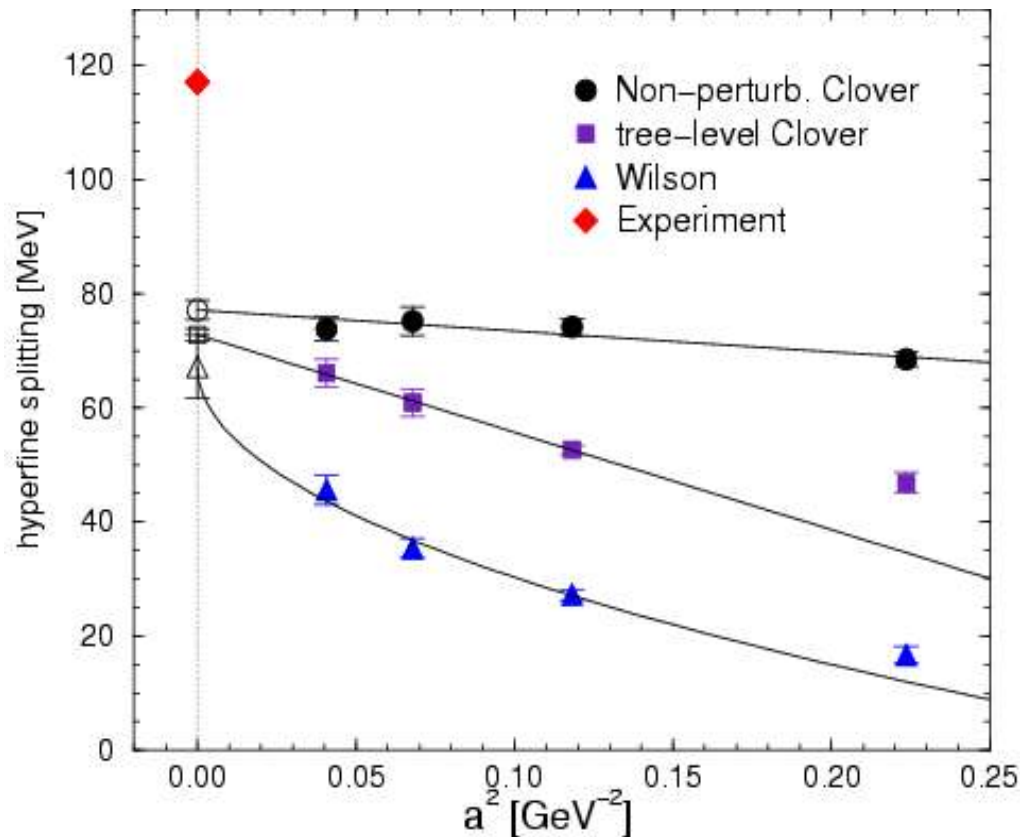


- nonperturbative clover
- continuum limit

quark mass set by $m(J/\psi)$
scale set by r_0

The results are consistent with the previous works
hyperfine splitting is smaller than the Exp. value

Hyperfine splitting in the $a=0$ limit



- nonperturbative clover $O(a^2)$
 - tree-level clover $O(g^2 a), O(a^2)$
 - standard Wilson $O(a)$
- ▶ nonperturbative clover shows small scale violation even for $m(J/\psi) a \sim 1.4$
 - ▶ other Dirac op. consistent with each other in the $a=0$ limit for $m(J/\psi) a < 1$

$m(J/\psi) - m(\eta_c) = 77 \text{ MeV}$ in the continuum limit

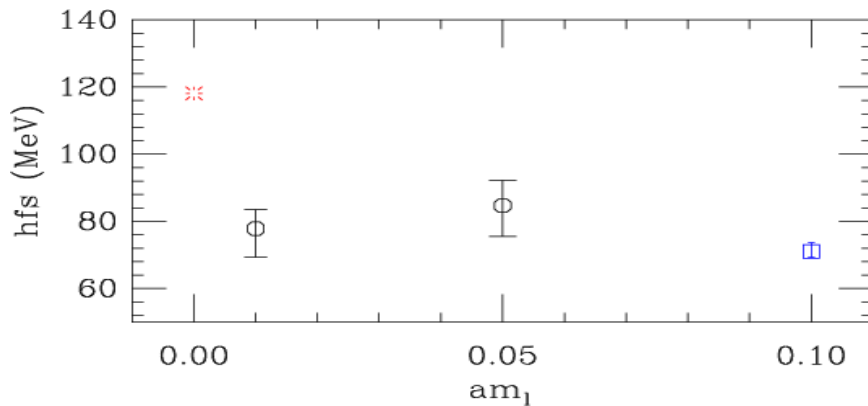
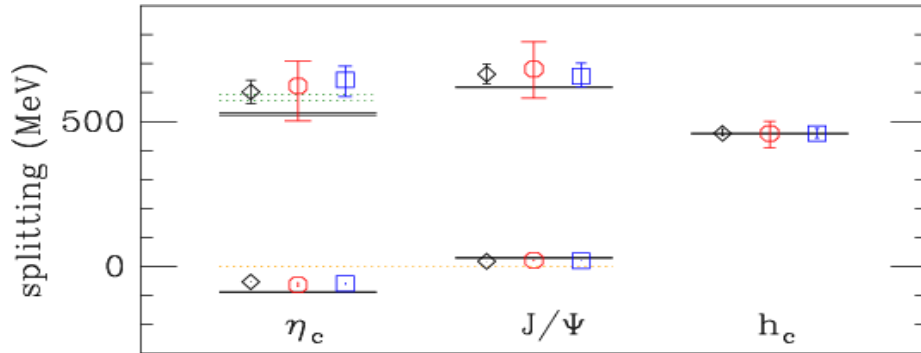
30–40% smaller than the experimental value = 117 MeV

Contained uncertainties

$$m(J/\psi) - m(\eta_c) = 77(2)(6) \text{ MeV} \quad (117\text{MeV in Experiment})$$

- statistical error ~ 2 MeV
- quark mass determination (m_{charm}) ~ 5 MeV
- continuum extrapolation ~ 3 MeV
- finite volume effects ~ checked & negligible
- scale determination ($1/a$) ~ dynamical quark effects
- choice of Dirac operators ~ well controlled
- dynamical quark effects
- disconnected diagram contributions

Dynamical quark effects



Nucl. Phys. B(PS)119(2003)586
M.di Pierro et al.

- valence quark : Fermilab action
- sea quark : KS quark
- $N_f = 0, 3, 2+1$
- no continuum limit $a = 0.13 \text{ fm}$

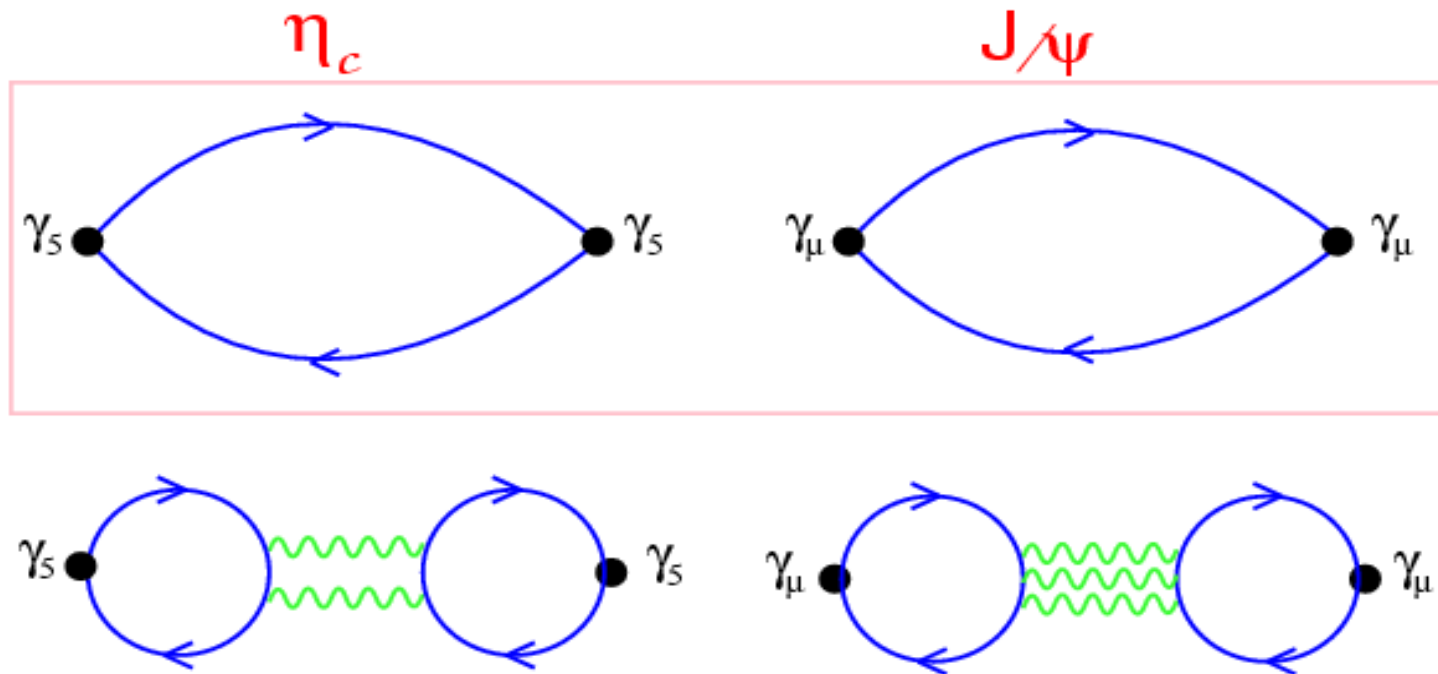


Other full QCD studies \longrightarrow **Dynamical quark effects $\sim 10\%$**

T.Manke et al., Phys. Rev. D62 (2000) 114508.

C.Stewart and R.Koniuk, Phys. Rev. D63(2001) 054503.

OZI forbidden “disconnected” diagram



Disconnected diagrams are neglected

because high cost & very small contribution

however, it may contribute to HFS $\sim O(10)$ MeV ?

Charmonium correlators

$$C(t) = \sum_{\vec{x}} \langle \text{Tr}[\Gamma D^{-1}(0, \vec{0}; t, \vec{x}) \Gamma D^{-1}(t, \vec{x}; 0, \vec{0})] \rangle$$

$$D(t) = \sum_{\vec{x}} \langle \text{Tr}[\Gamma D^{-1}(0, \vec{0}; 0, \vec{0})] \text{Tr}[\Gamma D^{-1}(t, \vec{x}; t, \vec{x})] \rangle$$

$D^{-1}(t, \vec{x}; t', \vec{x}')$: quark propagator

- $\Gamma = \gamma_5, \gamma_\mu$ (Pseudoscalar, Vector)
- source & sink operators are extended with $\phi(\vec{x}) \propto \exp(a|\vec{x}|^p)$
- disconnected diagrams are evaluated with
the Z2-noise method

Lattice setup

- Sea quark : **Nf=2 KS quark** : $a m_q = 0.1$

plaquette gauge : $\beta = 5.50$

lattice size : $12^3 \times 24$

lattice spacing : $a=0.16\text{fm}$ ($1/a=1.2\text{GeV}$) set by r_0

16,000 traj. (measurement at every 5 traj.)

☆ cutoff is not so sufficient for m_{charm} , this is an exploratory study

- Valence quark : **Fermilab action**

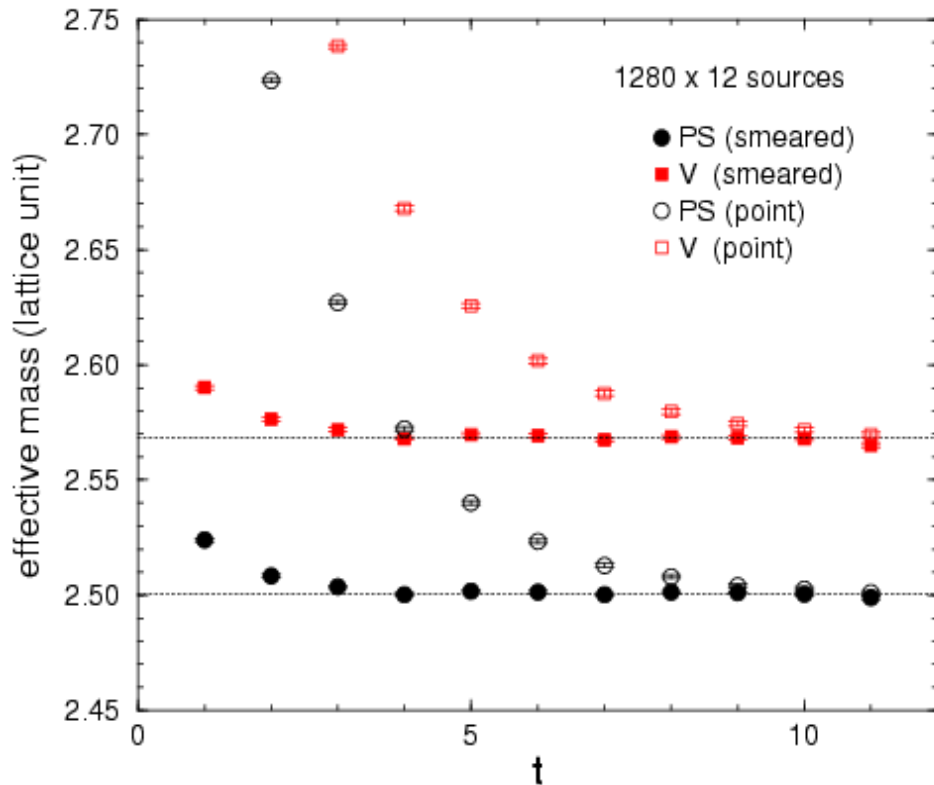
Csw : tadpole improved tree-level (u_0 in Landau gauge)

quark mass set by $m(J/\psi)$

- Z2-noise method

$$N_{\text{noise}}=600$$

Smeared operators



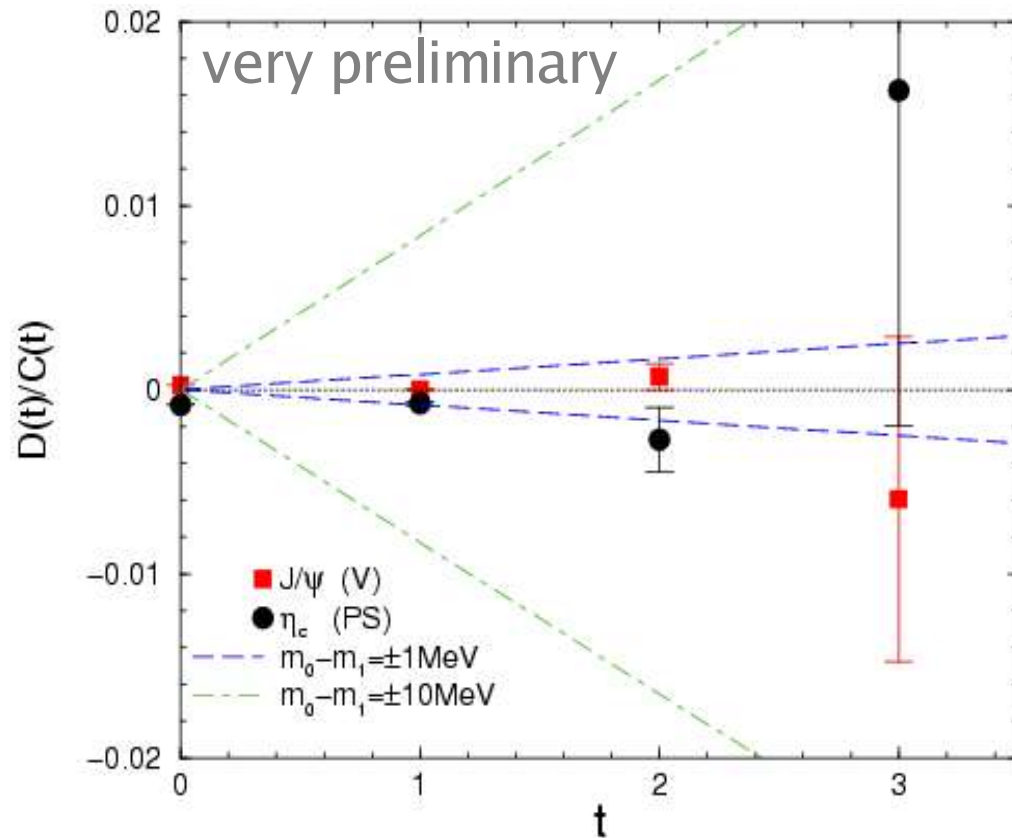
Effective mass plot
of connected diagram $C(t)$

Smearing functions $\phi(\vec{X})$ are
determined from wavefunction

Ground states dominate at $t \sim 2$

$$a(m(J/\psi) - m(\eta_c)) = 0.0676 \text{ (} \sim 81 \text{ MeV)}$$

Disconnected diagram contributions



When ground state dominates,
We have

$$C(t) = c \exp(-m_1 t)$$

$$C(t) + D(t) = d \exp(-m_0 t)$$



$$\frac{D(t)}{C(t)} = \frac{d}{c} \exp[(m_1 - m_0) t] - 1$$

- Too noisy at large t
- Contributions of the disconnected diagram may be very small ?

Summary & Outlook

We study the problem of charmonium HFS
and consider a possibility of disconnected diagram contributions

- **Isotropic lattice + Nonperturbative improved clover quark**

is good choice for the charmonium spectrum

In quenched QCD, $m(J/\psi) - m(\eta_c)$ is

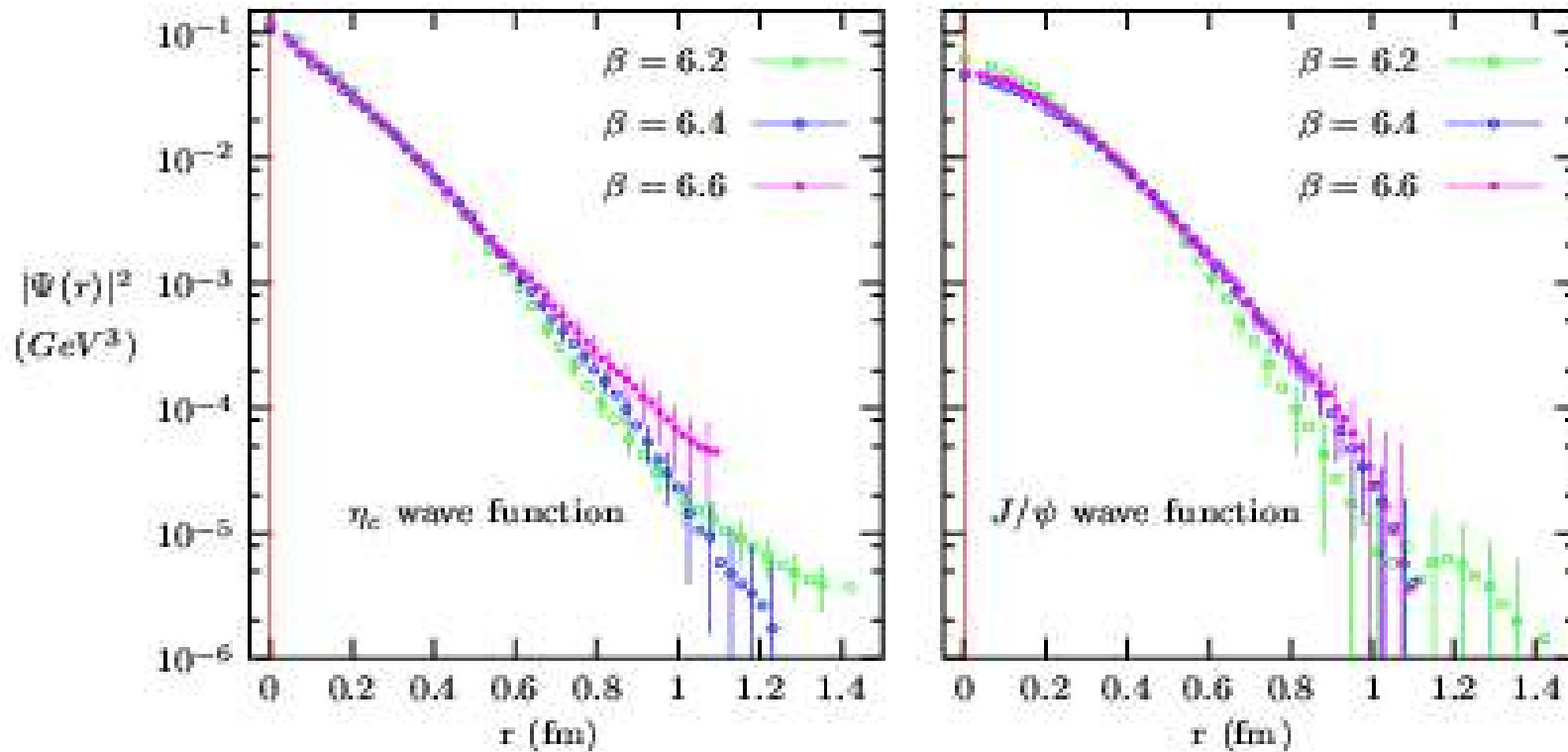
30 – 40% smaller than the experimental value

- Disconnected diagram contributions are very small
or hidden by large error

- ▶ Same calculations with smaller valence quark mass

-  **extrapolation from light quark mass region**

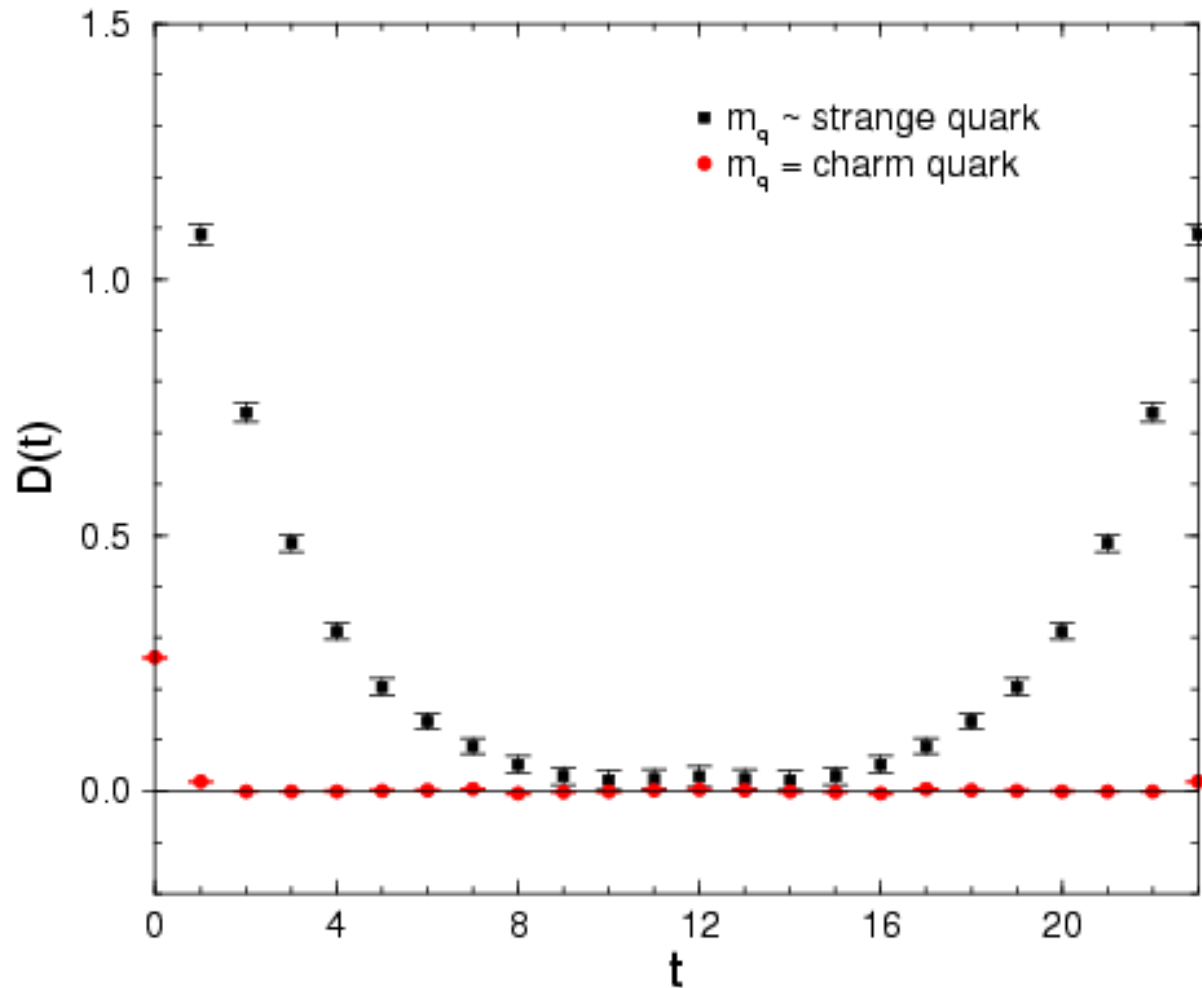
Appendix: Wave function



- ▶ scaling violations are very small
- ▶ $\Psi(\vec{0})$'s agree with phenomenological expectations qualitatively

In a heavy quark model : $\Psi_{\eta_c}(\vec{0}) > \Psi_{J/\psi}(\vec{0})$

Appendix: lighter valence quarks



Appendix: Z2-noise method

