

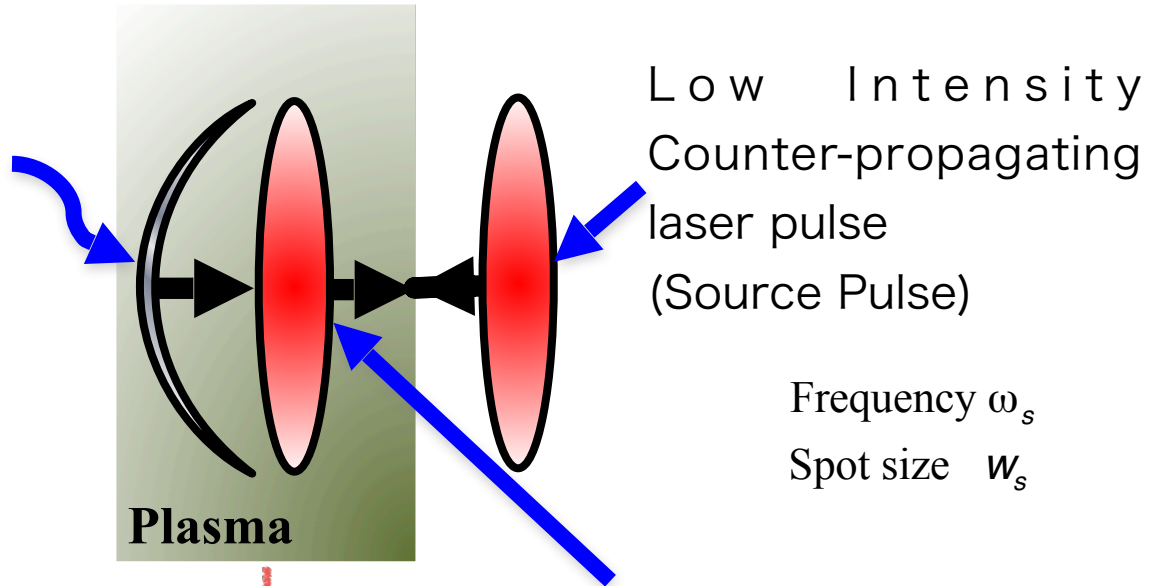
Reflection of Strong Long Wavelength Laser Pulses from Relativistic Mirrors

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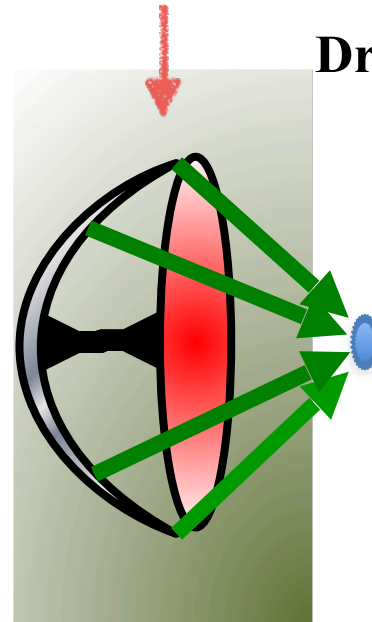
Introduction

Relativistic Flying Mirror
 Breaking plasma wave



Driver laser pulse

- S. V. Bulanov, et al., Phys. Rev. Lett. 91, 085001 (2003)
- M. Kando et al., Phys. Rev. Lett. 99, 135001 (2007)
- M. Kando, et al., Phys. Rev. Lett. 103, 235003 (2009)
- S. V. Bulanov et al., Plasma Sources Sci. Technol. 25, 053001 (2016)



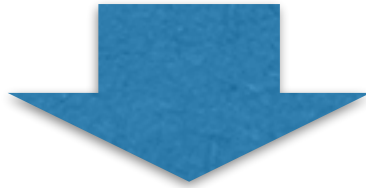
Phase velocity $\beta_{ph}c$

Up-shifted $\omega'' = \frac{1 + \beta_{ph}}{1 - \beta_{ph}} \omega_s \approx 4\gamma_{ph}^2 \omega_s$

Sharply focused $w' = \sqrt{\frac{1 - \beta_{ph}}{1 + \beta_{ph}}} w_s \approx \frac{w_s}{2\gamma_{ph}}$

Source pulse

- Same wavelength as Driver
- Low intensity



- Long wavelength
- Lower critical density → Better reflection
- Finite existence time of wake → Few cycles
- Nonlinear regime
- Harmonics

Parameters

• Plasma

$$L = 50\mu m$$

$$n_p = 5 \times 10^{19} cm^{-3}$$

• Driver laser

$$\lambda_D = 1\mu m$$

$$I_D = 2 \times 10^{19} W/cm^2$$

$$a_D = 3.8$$

$$l_D \approx 3.3\mu m$$

• Source laser

$$\lambda_S = 3\mu m$$

$$I_S = 10^{12} \sim 10^{18} W/cm^2$$

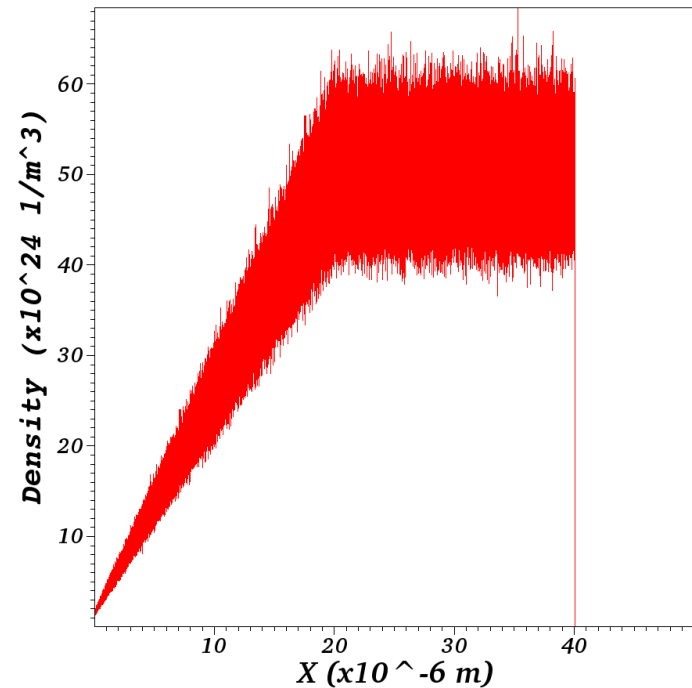
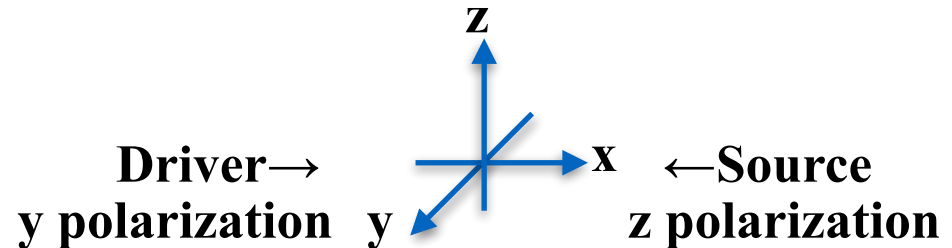
$$a_S = 2.5 \times 10^{-3} \sim 2.55$$

$$l_S \approx 10\mu m$$

• Simulation

$$\cdot 1D \quad \Delta_x = 1/9000\mu m$$

$$\cdot 10 \text{ particles/cell}$$



EPOCH code

Wave breaking

$$a^{3/2} \geq (\omega/\omega_{pe}) \quad a = 2.81 \quad a_D > a$$

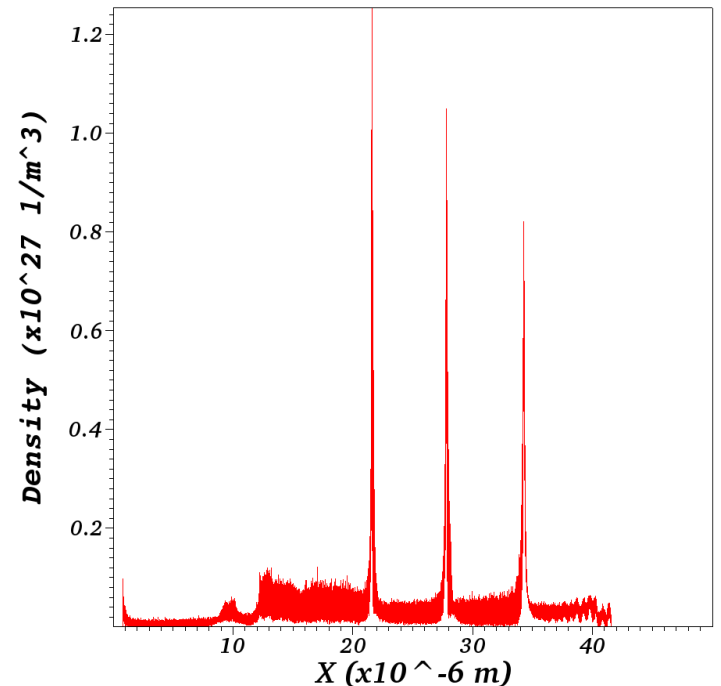
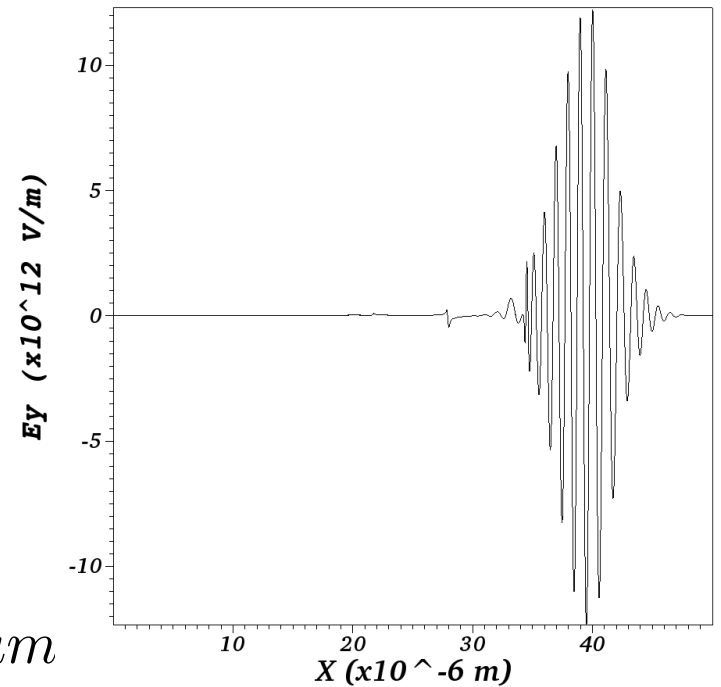
Optimal laser length

$$a^2 \gg 1$$

$$l_{opt} = 2|a| + \left(\frac{1}{2} + 2 \ln 2 - \ln |a| \right) \frac{1}{|a|} + O \left[\frac{1}{|a|} \right]^3 \quad l_{opt} \approx 7.9 \mu m$$

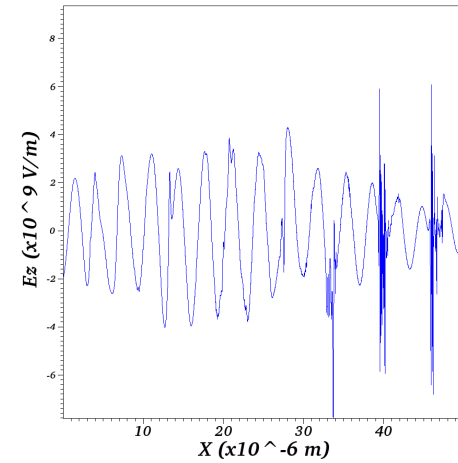
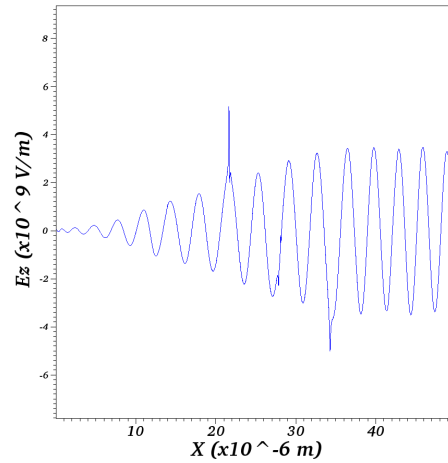
Wake wavelength

$$\lambda_w = 4ca/\omega_{pe} \quad \lambda_w \approx 11 \mu m$$

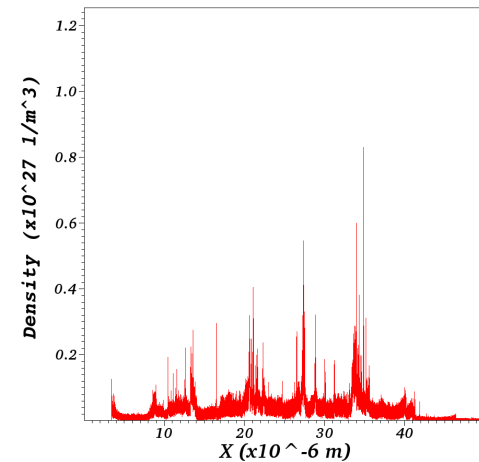
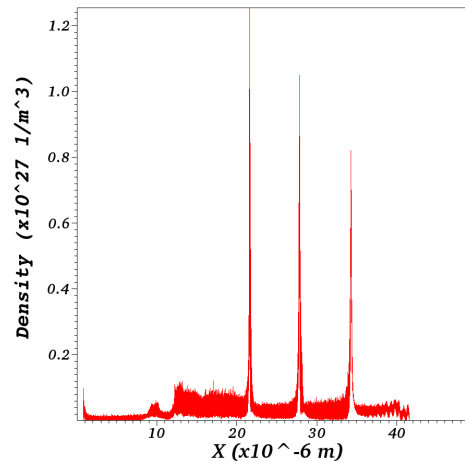


Weak Source pulse

10^{12} W/cm^2

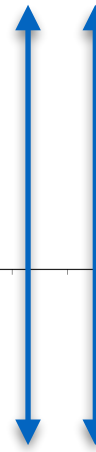
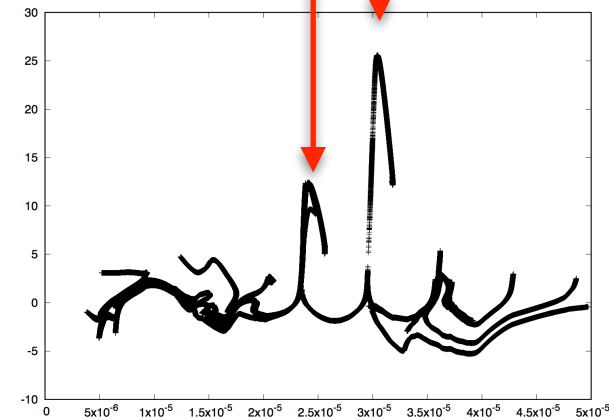
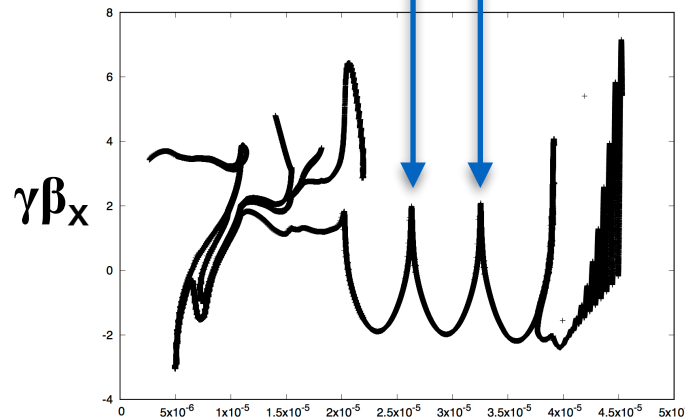
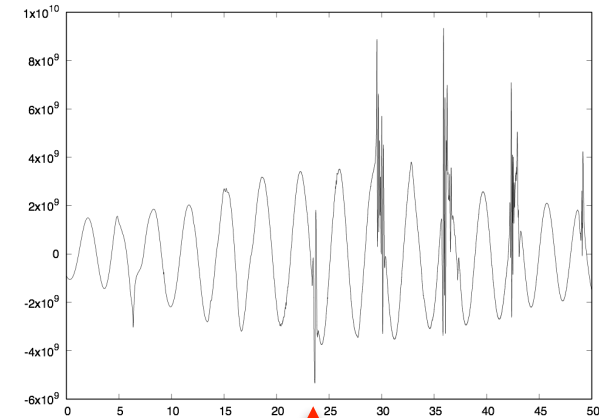
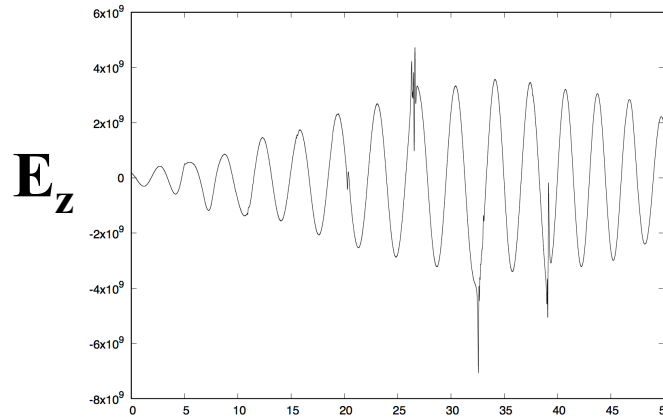


- Reflections can be seen
- Occur at various points



Phase Space

- Source pulse reflection from waves of various momenta
- non-breaking and breaking



Spectrum

- Fourier transform of source pulse (Hanning filter)
- Broad spectrum
- Circularly polarized pulse

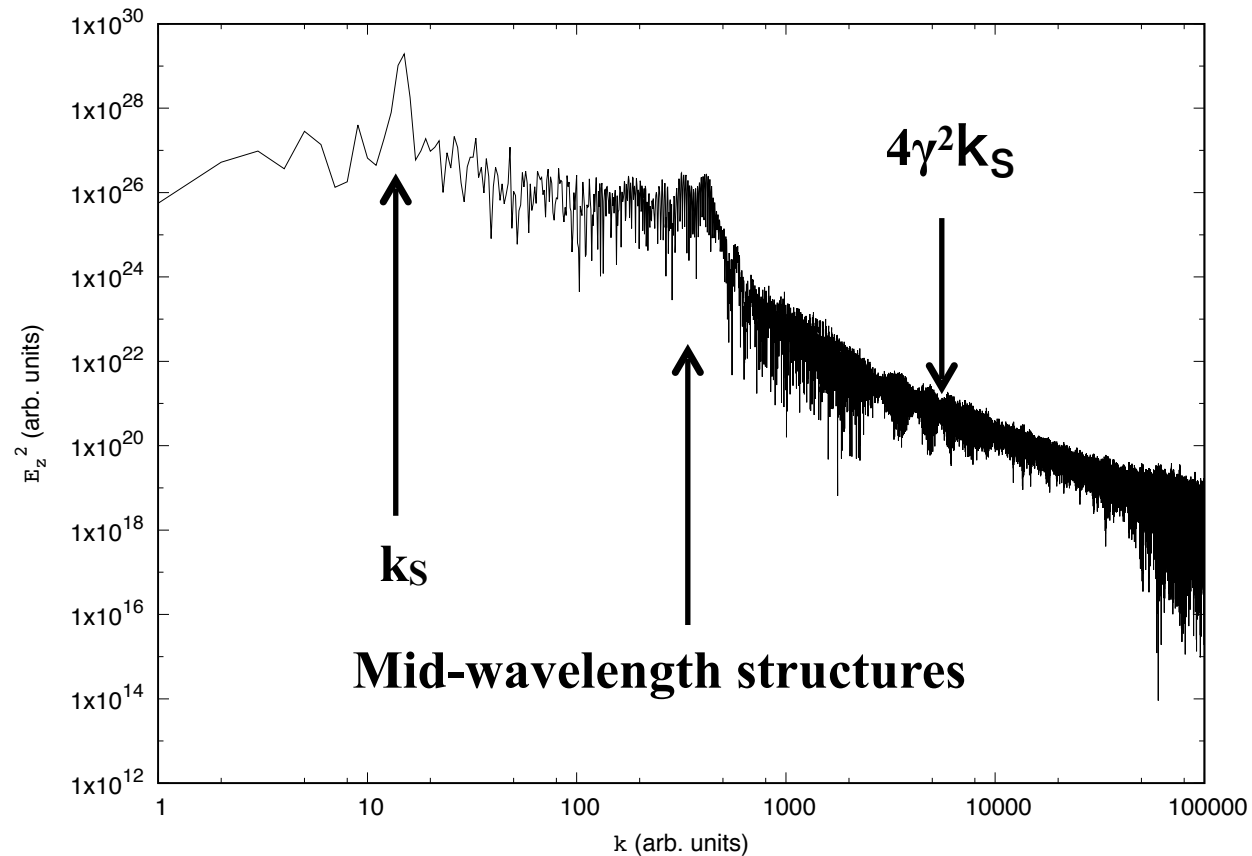
$$\gamma = \left(\frac{\omega}{\omega_{pe}}\right)(1 + a^2)^{1/4}$$

$$a = 3.8$$

$$\frac{\omega}{\omega_{pe}} = 4.72$$

$$\gamma = 9.36$$

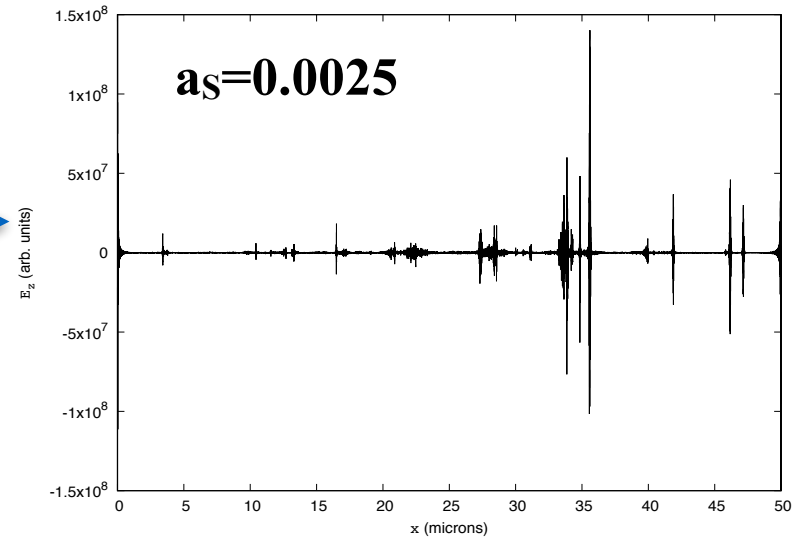
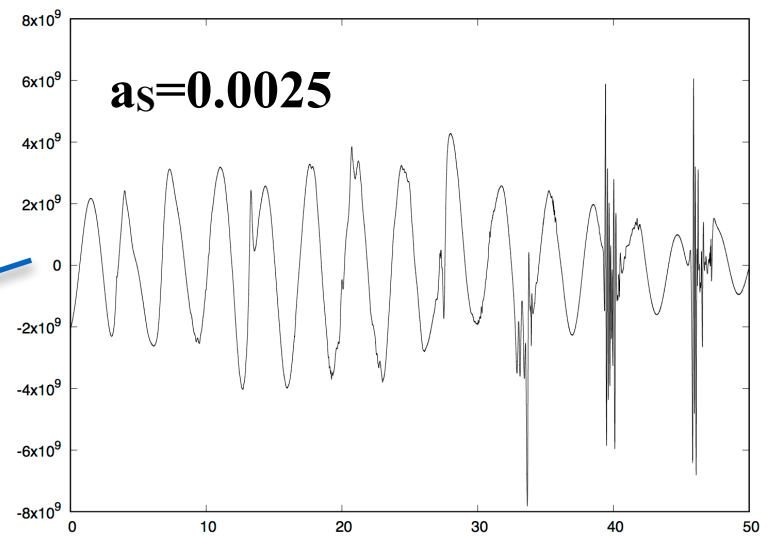
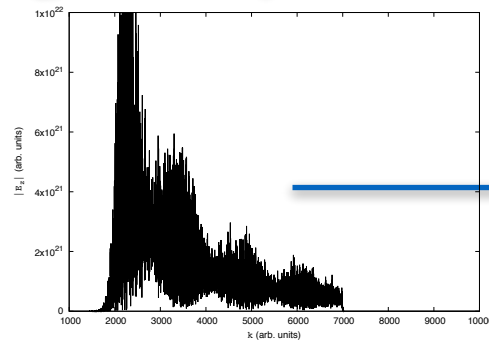
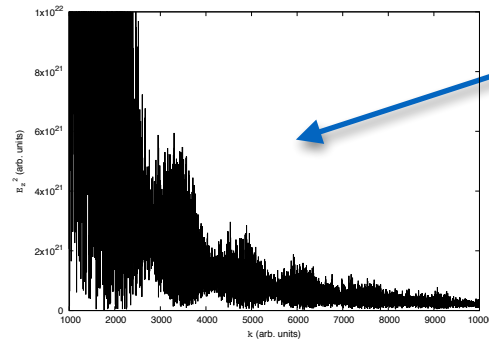
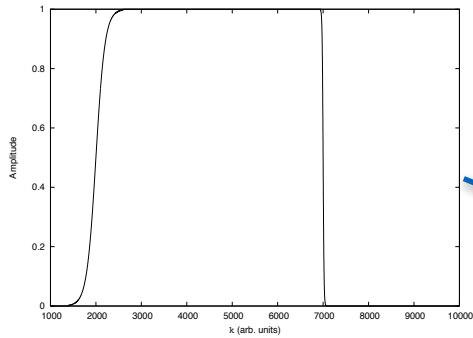
$$4\gamma^2 = 350.57$$



A. I. Akhiezer and R. V. Polovin, *Sov. Phys. JETP* 3, 696 (1956)

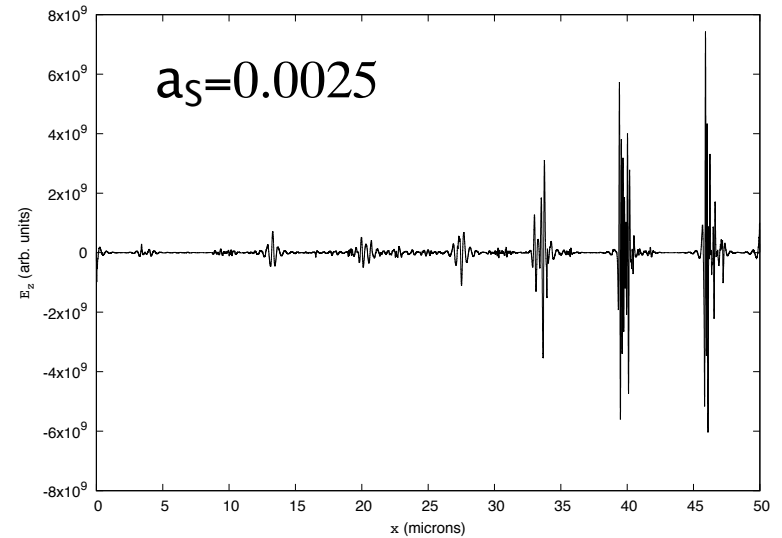
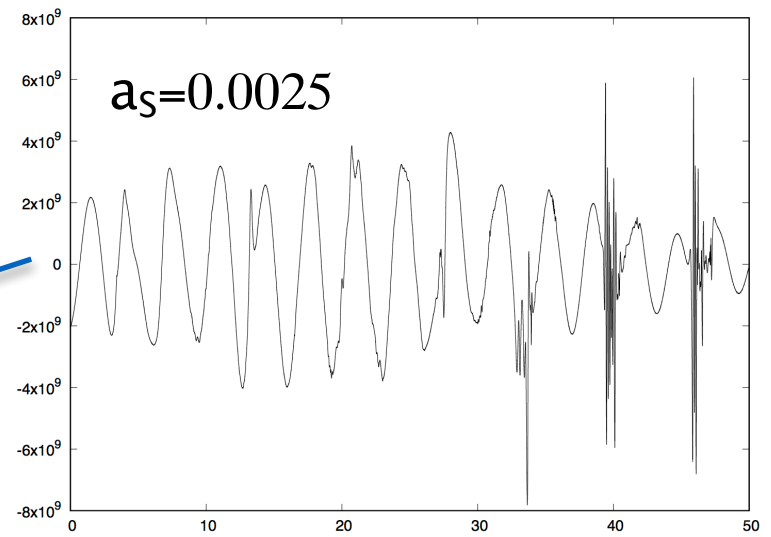
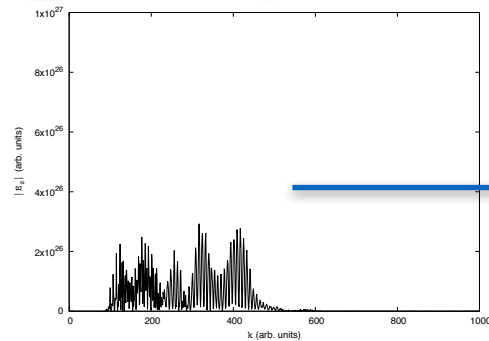
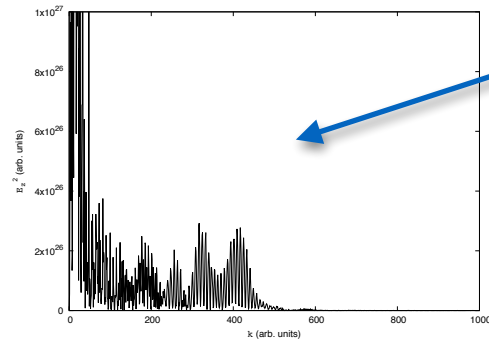
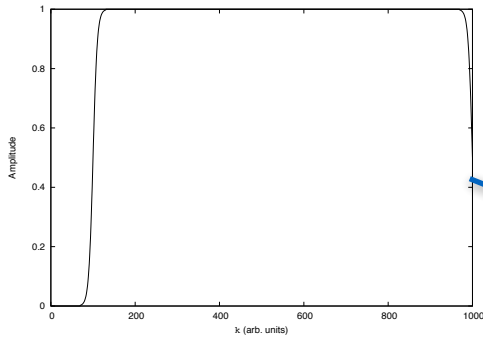
Weak Source pulse

- High bandpass filter to locate source signal
- Tanh function at edges



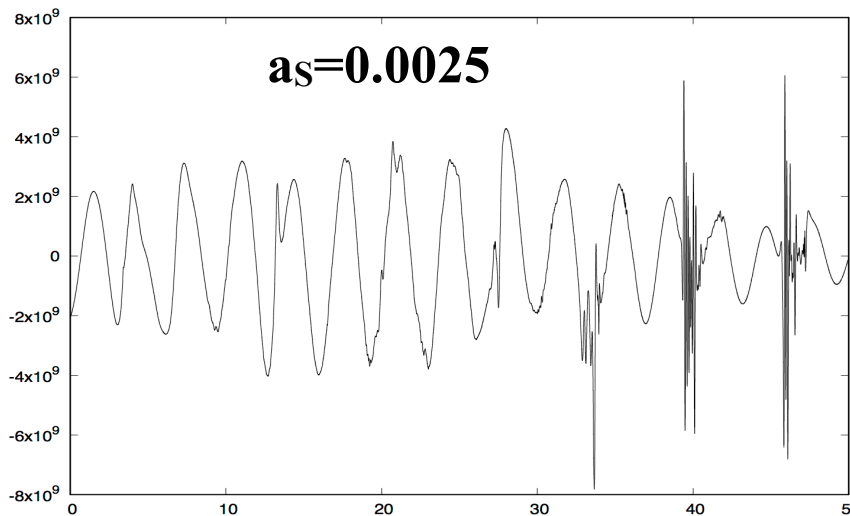
Weak Source pulse

- Low bandpass filter
- Tanh function at edges

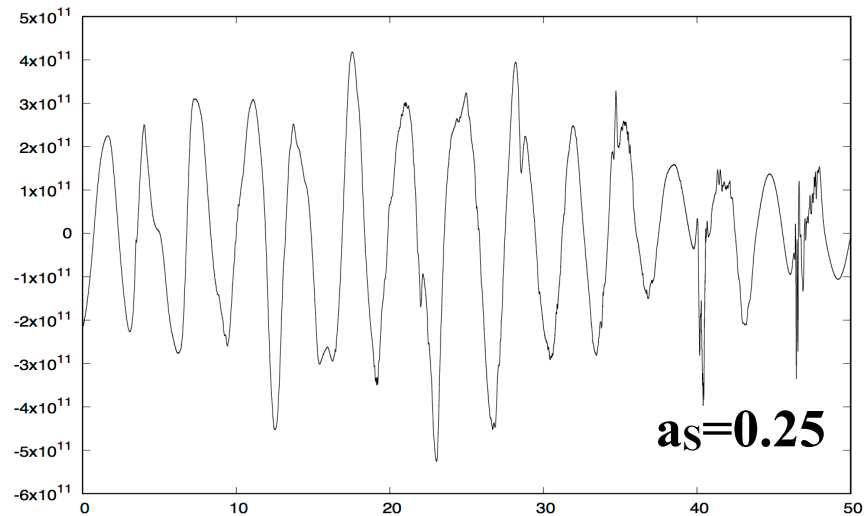


Higher Source Intensity

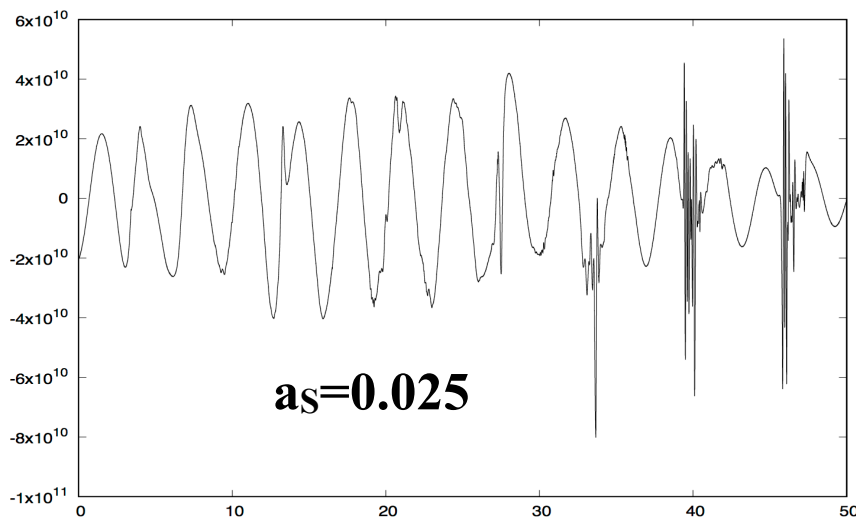
10^{12} W/cm^2



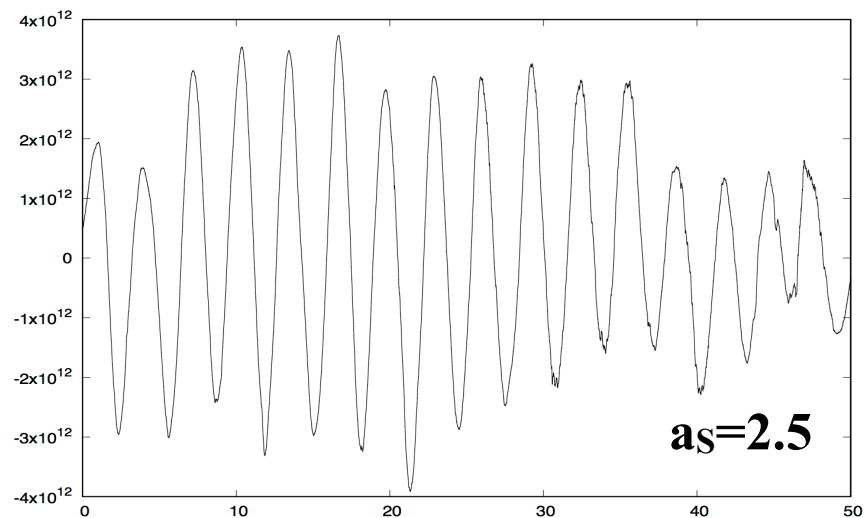
10^{16} W/cm^2



10^{14} W/cm^2



10^{18} W/cm^2



Spectra

10^{18} W/cm^2

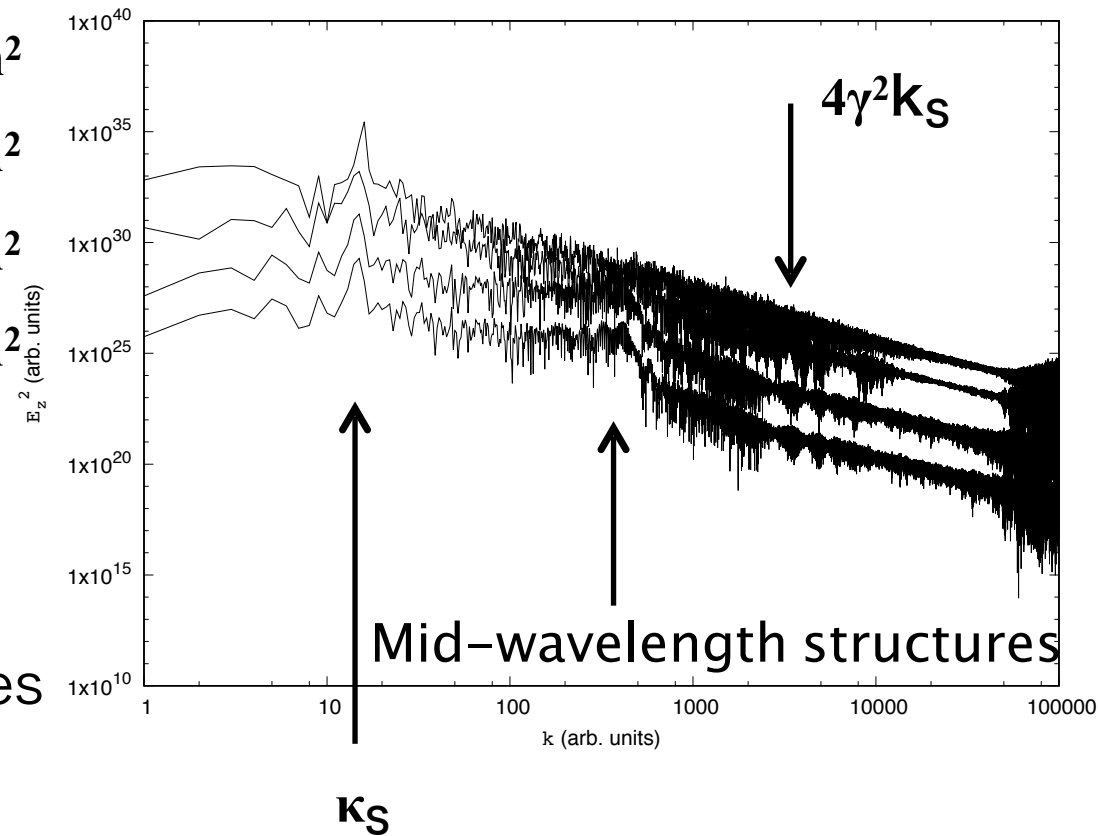
10^{16} W/cm^2

10^{14} W/cm^2

10^{12} W/cm^2

- Harmonics $a_s=0.25$
- Mid-wavelength structures not present $a_s>0.25$

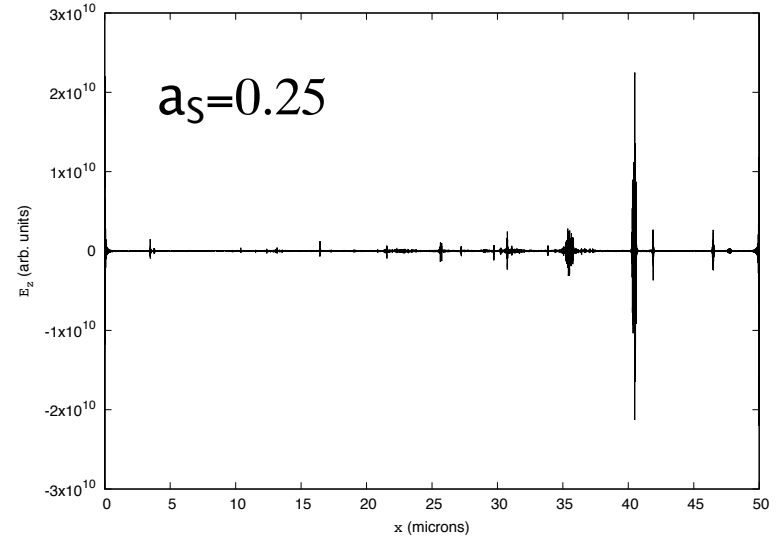
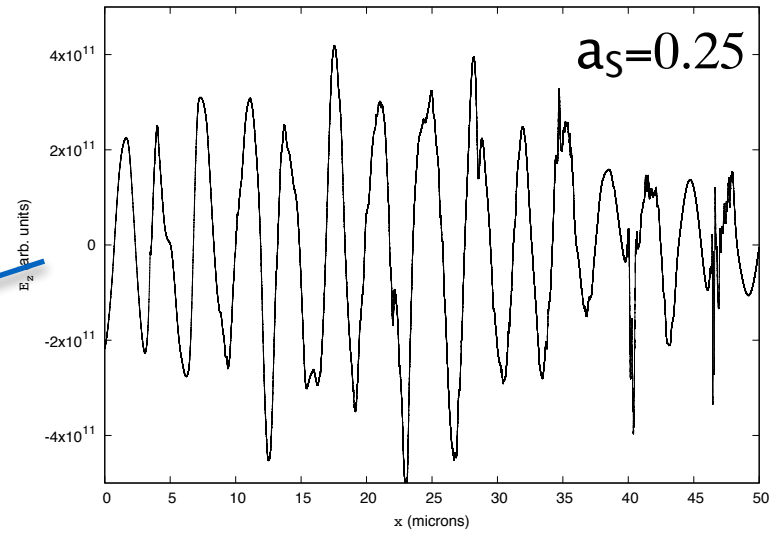
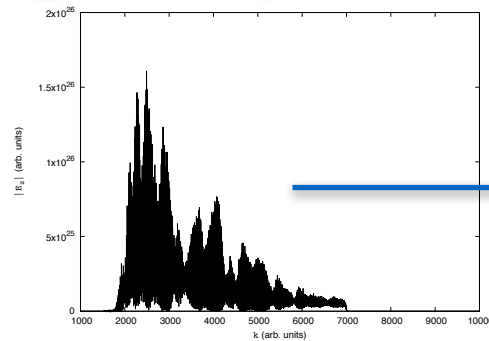
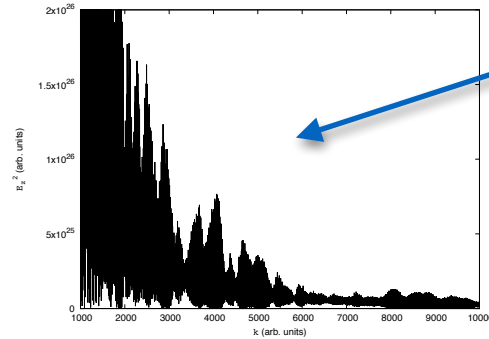
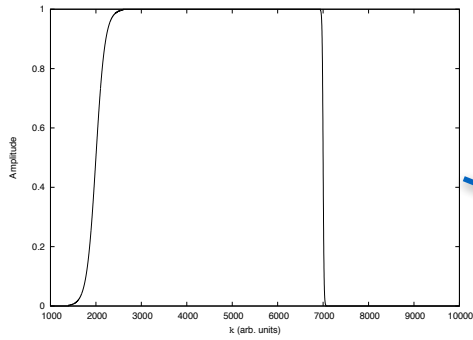
$|E_z(\mathbf{k})|$ (arb. Units)



k (arb. Units)

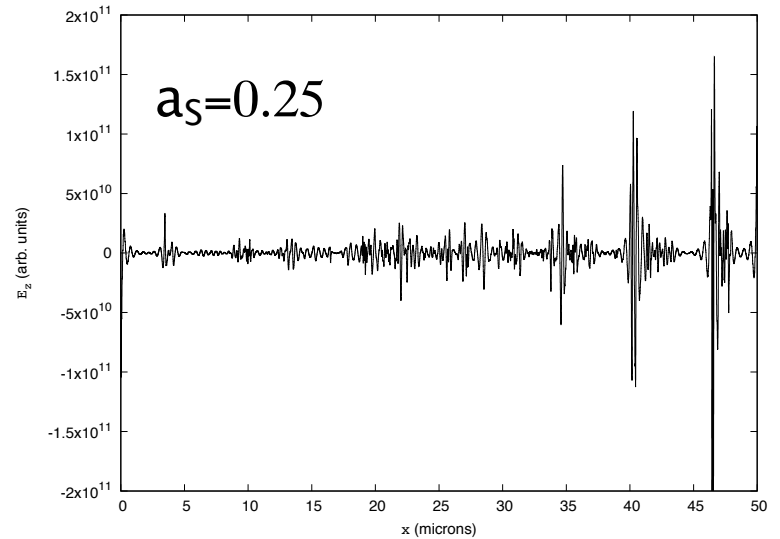
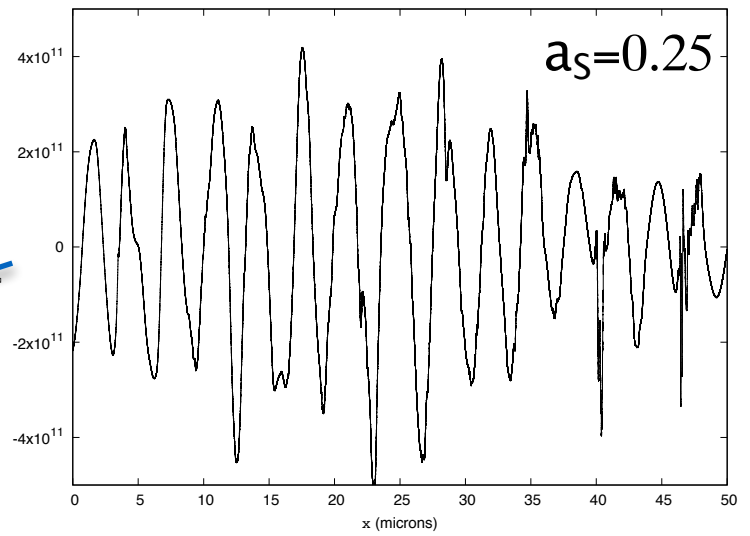
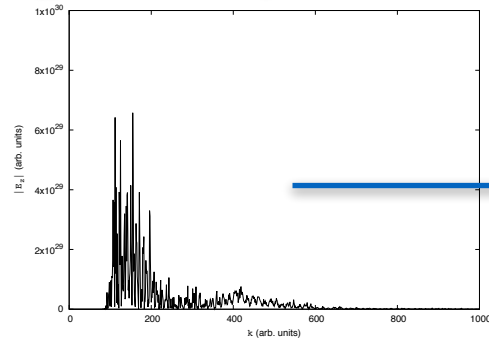
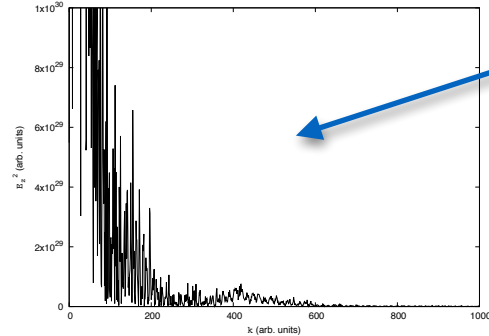
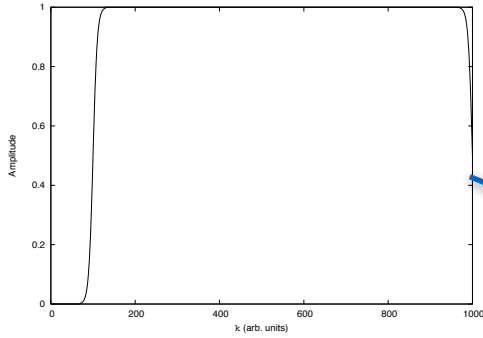
Moderate Source pulse

- High bandpass filter
- Tanh function at edges



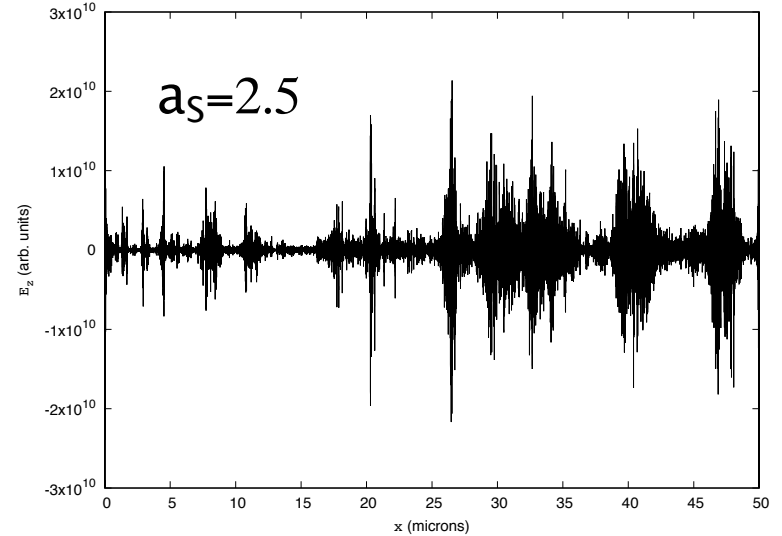
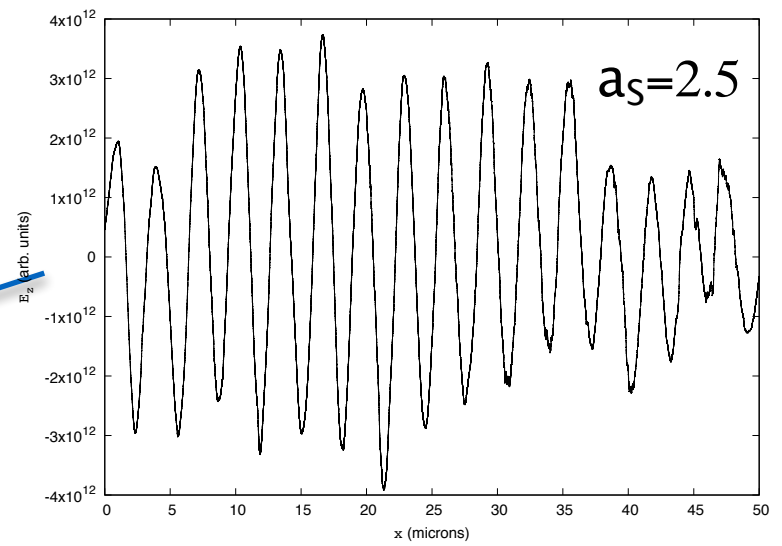
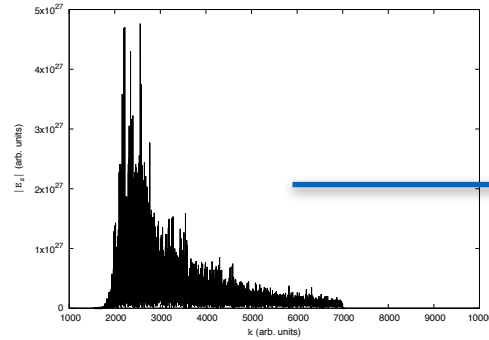
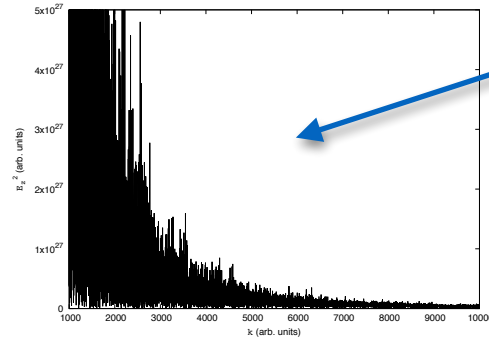
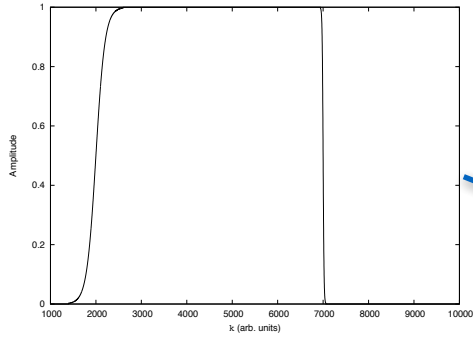
Moderate Source pulse

- Low bandpass filter
- Tanh function at edges



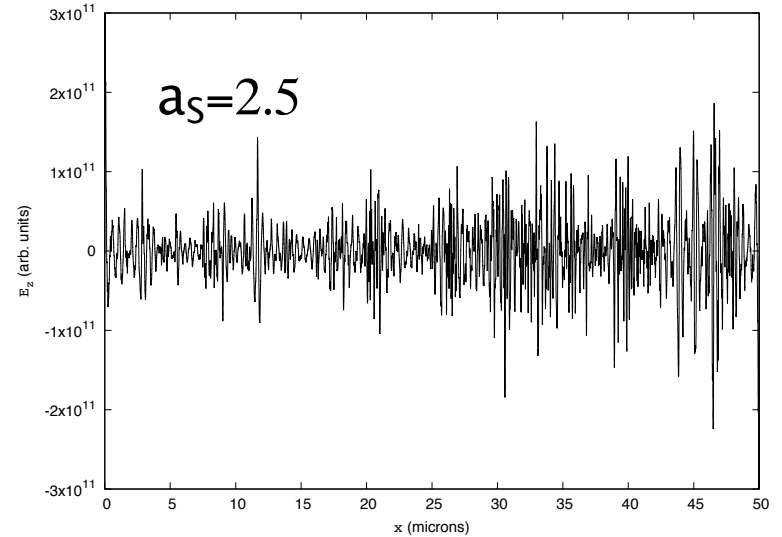
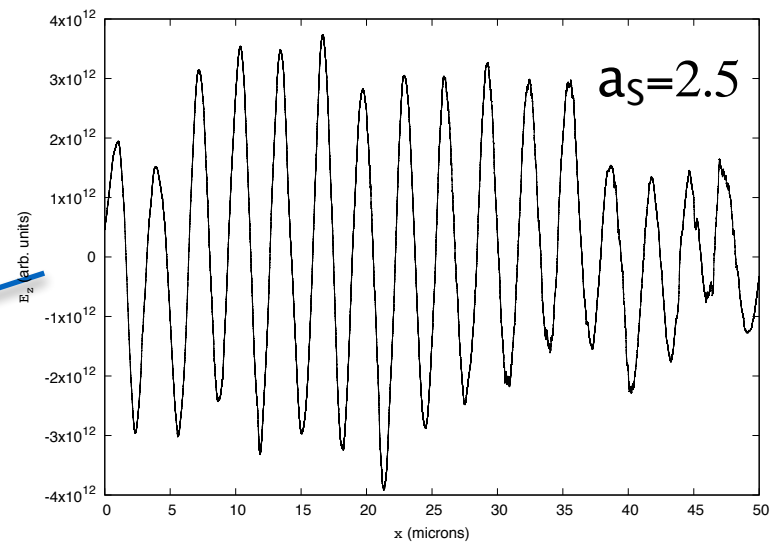
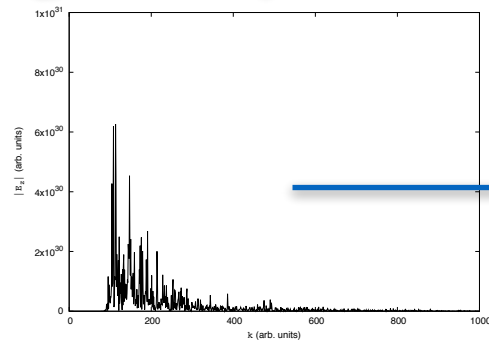
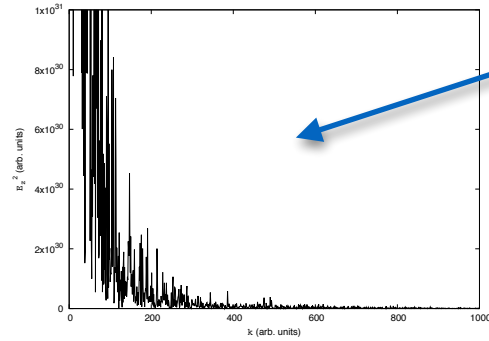
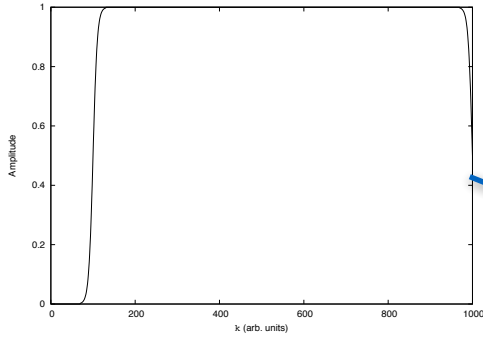
Strong Source pulse

- High bandpass filter
- Tanh function at edges



Strong Source pulse

- Low bandpass filter
- Tanh function at edges



Conclusions

- Reflection from relativistic flying mirror
- Long wavelength/nonlinear regime
- Broad spectrum observed
- Reflection from non-breaking and breaking waves
- Harmonics appear at $a_s=0.25$
- High bandpass filter position of reflected pulses
- Low bandpass filter position of mid-wavelength structures
- At higher source intensities mid-wavelength structures not present
- Source pulse is generating its own wake at high intensities
- Next
- More detailed analysis of 1D results
- 2D Simulations
- After that Radiation Reaction effects

Simulations performed on “SGI ICE X” super computer at Japan Atomic Energy Agency